

Search For New And Unusual Strangeonia In Photoproduction Using CLAS[†]

Mukesh S. Saini

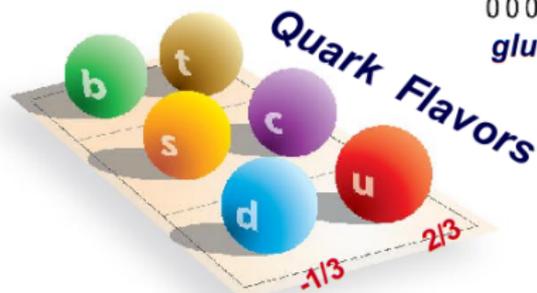
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[†] Research funded by the US Department of Energy (DOE)

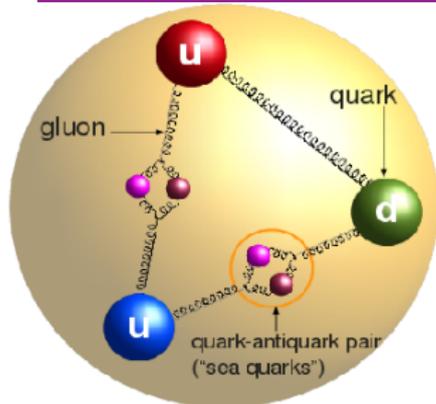
- ▶ Meson Spectroscopy
- ▶ Experiment
- ▶ Data Selection
- ▶ Strangeonium Analysis
- ▶ Summary



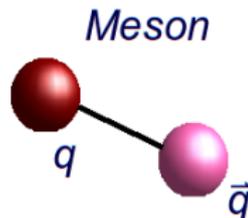
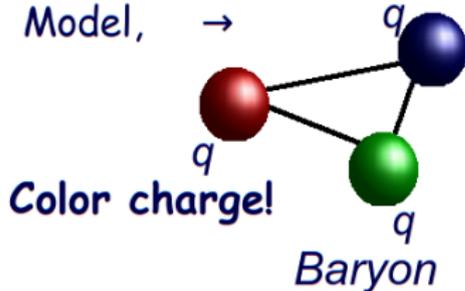
Quantum chromodynamics (QCD) is theory of the strong interaction (color force).

It describes the interactions of the quarks and gluons making up the hadron

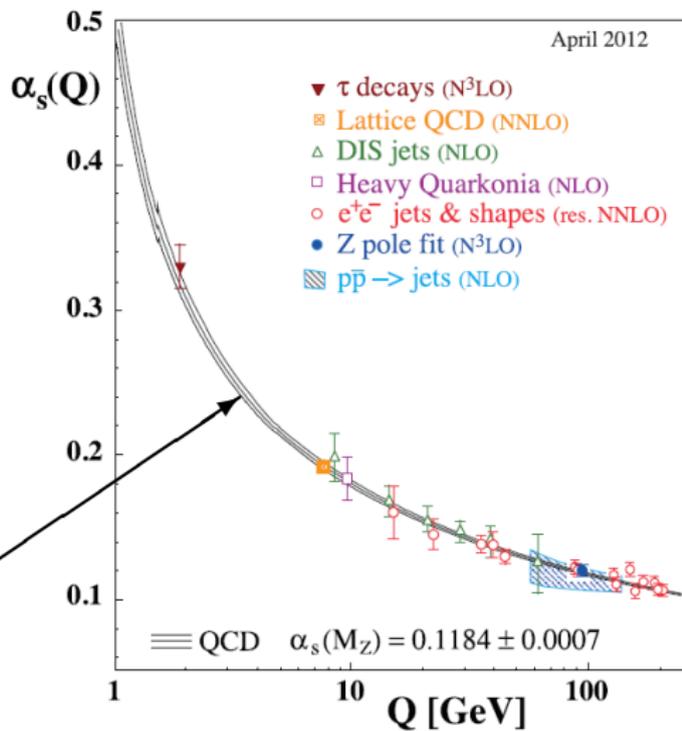
QCD Picture



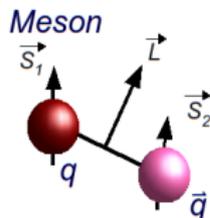
In constituent quark Model, \rightarrow



- ▶ Free quarks and gluons have not been observed in nature due to confinement.
- ▶ QCD predicts exotic hadrons beyond the naive quark model [hybrids, glueballs and multi-quark states]
- ▶ Mapping of the meson spectra will help us identify exotic unconventional mesons and decays, to further our insight into soft (Non-perturbative) QCD



PDG, 2012



$$\vec{J} = \vec{L} + \vec{S}$$

$$P = (-1)^{L+1}$$

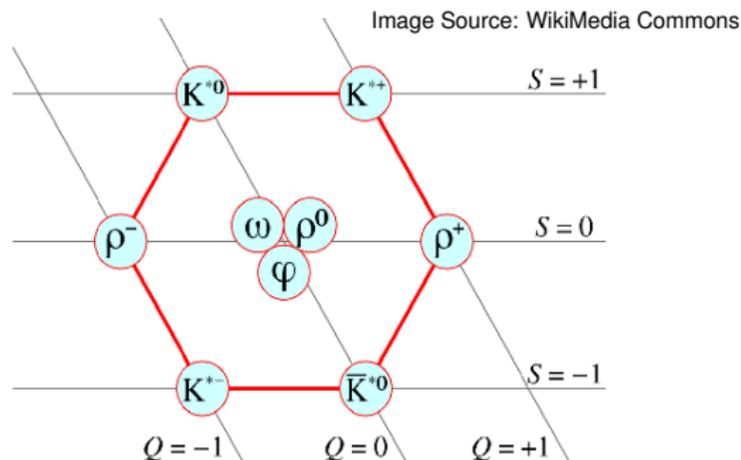
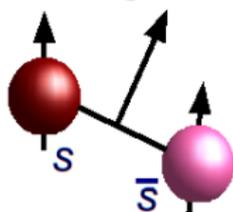
$$C = (-1)^{L+S}$$

$$J^{PC} \Big|_{\text{allowed}} = 0^{-+}, 0^{++}, 1^{--}, 1^{+-}, 1^{++}, 2^{--}, \dots$$

$$J^{PC} \Big|_{\text{exotic}} = 0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, 3^{-+}, \dots$$

Usual Mesons	EXOTICA			
	Gluonic Hybrids	Tetra-quarks	Glue-Balls	...
$q\bar{q}$	$q\bar{q}g$	$q\bar{q}q\bar{q}$	gg	...

Strangeonia



- ✓ Of the **22** expected resonances from 3P_0 model predictions, only **7** candidates identified

$\eta - \eta'$	ϕ (1020)	h_1 (1380) [†]
f_1 (1420)	f_2' (1525)	ϕ (1680)
ϕ_3 (1850)	?	?

[†] Not Included in the PDG Summary Tables.

			J^{PC}	Name	Mass (MeV)
n=2	L=0	S=0	0^{-+}	η_s	1415
		S=1	1^{--}	ϕ	1680
	L=1	S=0	1^{+-}	h_1	1850
		S=1	0^{++}	f_0	2000
			1^{++}	f_1	1950
		2^{++}	f_2	2000	
n=3	L=0	S=0	0^{-+}	η_s	1950
		S=1	1^{--}	ϕ	2050

← Radial excitations of ($l = 0, s\bar{s}$) meson.

			J^{PC}	Name	Mass (MeV)
n=1	L=0	S=0	0^{-+}	η, η'	548, 958
		S=1	1^{--}	ϕ	1020
	L=1	S=0	1^{+-}	h_1'	1380
		S=1	0^{++}	f_0'	1500
			1^{++}	f_1'	1530
		2^{++}	f_2'	1525	
	L=2	S=0	2^{-+}	η_2	1850
		S=1	1^{--}	ϕ_1	1850
			2^{--}	ϕ_2	1850
		3^{--}	ϕ_3	1854	

← Orbital excitations of ($l = 0, s\bar{s}$) meson.

Tables from reference:
T. Barnes, N. Black and P. R. Page,
Phys. Rev. D 68, 054014 (2003)

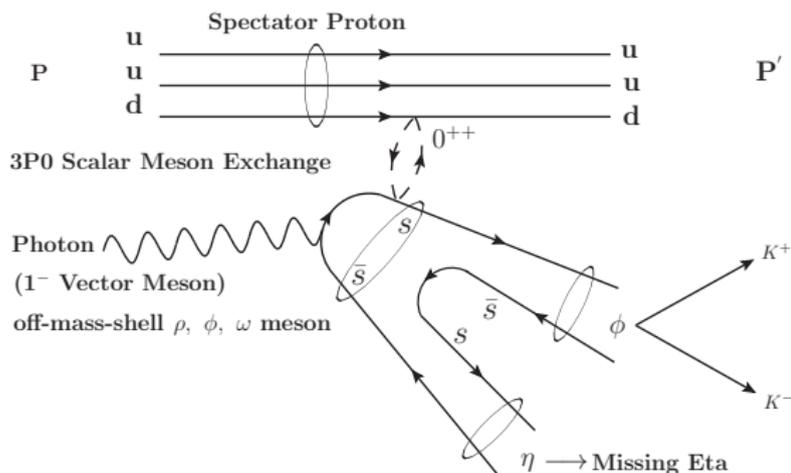
Why Study Strangeonia?

- ▶ Due to the intermediate mass of the strange quarks, study of the strangeonium states will serve as a bridge between short and large distance behavior of QCD confinement potential, a study of the transition from light quark sector to the HQET

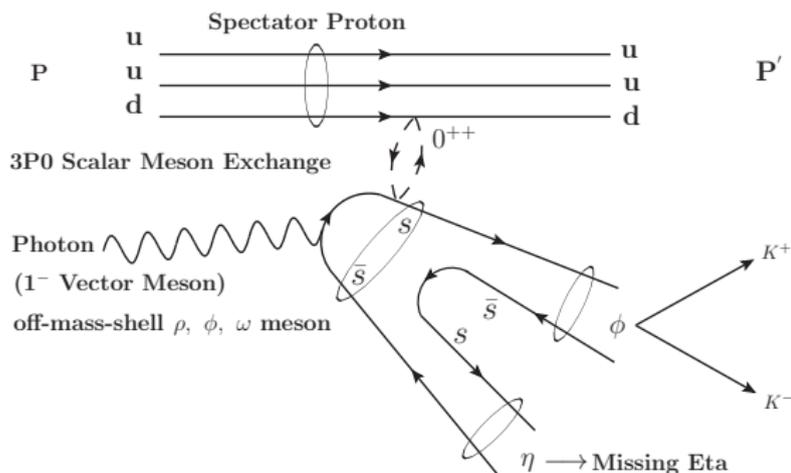
How To Study Strangeonia

- ▶ "Due to the *OZI* rule, the observation of a state with a large branching fraction to $\eta\phi$, $\eta'\phi$ or $\phi\phi$ and small branches to nonstrange final states can serve as a "smoking gun" for an initial $s\bar{s}$ state." - Barnes, Black & Page (Strong decays of Strange Quarkonia)
- ▶ Open strangeness decay modes (KK , KK^*) for initial $s\bar{s}$ are susceptible to being confused with $n\bar{n}$, where $n \in \{u, d\}$, which can also decay via KK , KK^*
- ▶ $\phi\eta$ and $\phi\eta'$ decay modes are most likely to originate from an initial $s\bar{s}$ state, even though η has $n\bar{n}$ component

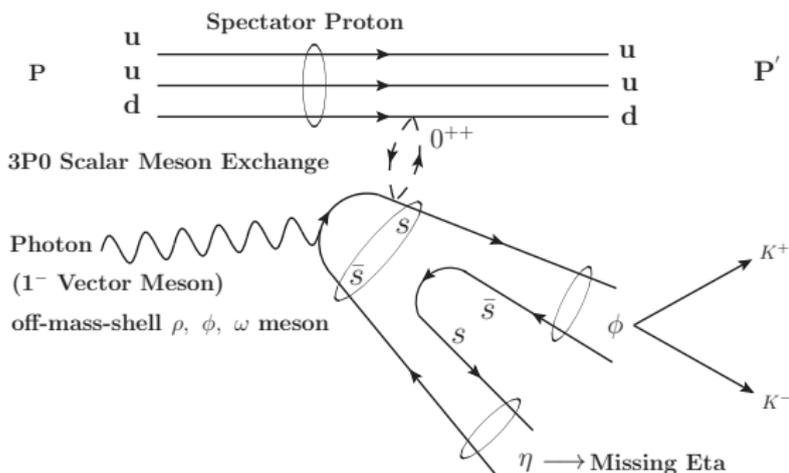
- ▶ In hadronic interactions, $[\gamma \iff \{\rho, \omega, \phi\}]$ - with an important $s\bar{s}$ component - Vector Meson Dominance (VMD)
- ▶ Study of diffractive photoproduction reaction $\gamma p \rightarrow p \phi \eta$, should lead to observation of many $C=-1$ $s\bar{s}$ states
- ▶ If the decay products of a meson have a ϕ meson with another mostly strange meson like η , OZI rule dictates that initial state be $s\bar{s}$



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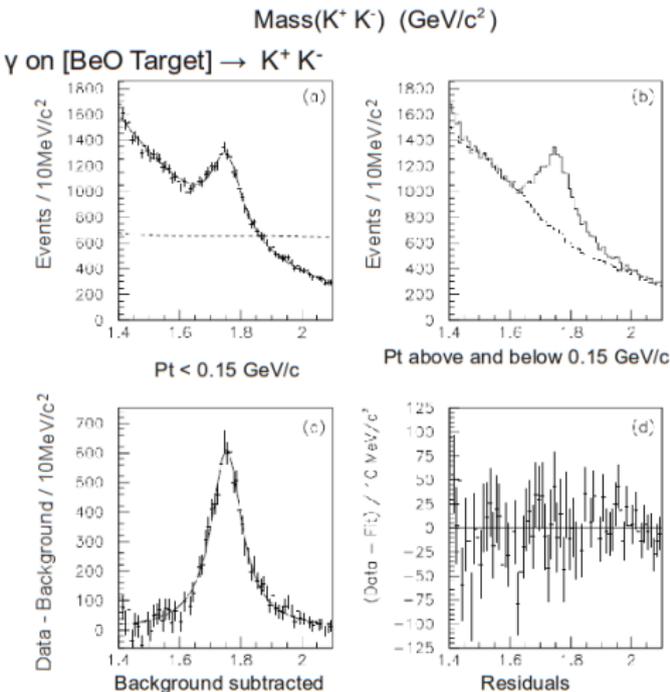


Image from reference:
 J. M. Link et al., [FOCUS Collaboration],
 Phys. Lett. B 545, 50 (2002)

FOCUS Experiment

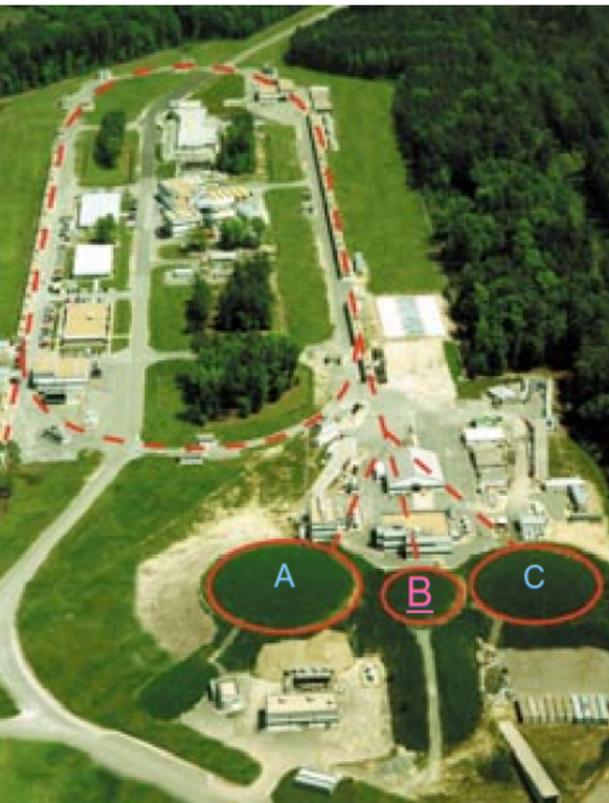
- ▶ e^+e^- annihilation observe $\phi(1680)$ in KK^* decay
- ▶ Photoproduction experiments observe a resonance at 1750 MeV/c² in $X(1750) \rightarrow K^+K^-$
- ▶ Earlier photoproduction results had low statistics \sim few 100 events
- ▶ FOCUS@Fermilab has 11,700 $X(1750) \rightarrow K^+K^-$
- ▶ FOCUS studied $KK : KK^*$ branching ratio and claimed $X(1750) \neq \phi(1680)$
- ▶ If $X(1750) \in \{s\bar{s}\}$, cleanest way to observe it will be $X \rightarrow \phi\eta$

Strangeonium State		Predicted Decay Width Theory MeV/c ²
$\phi(1680)$	2S	$\Gamma_{theory} = 378$ $\Gamma_{\phi\eta} = 44$
$\phi(2050)$	3S	$\Gamma_{theory} = 378$ $\Gamma_{\phi\eta} = 21$
$h_1(1850)$	2P	$\Gamma_{theory} = 193$ $\Gamma_{\phi\eta} = 33$
$\phi_3(1854)$	1D	$\Gamma_{theory} = 104$ $\Gamma_{\phi\eta} = 3$
$\phi_2(1850)$	1D	$\Gamma_{theory} = 214$ $\Gamma_{\phi\eta} = 53$
$\phi(1850)$	1D	$\Gamma_{theory} = 652$ $\Gamma_{\phi\eta} = 29$
$h_3(2200)$	1F	$\Gamma_{theory} = 249$ $\Gamma_{\phi\eta} = 5$

Experimentally Observed States

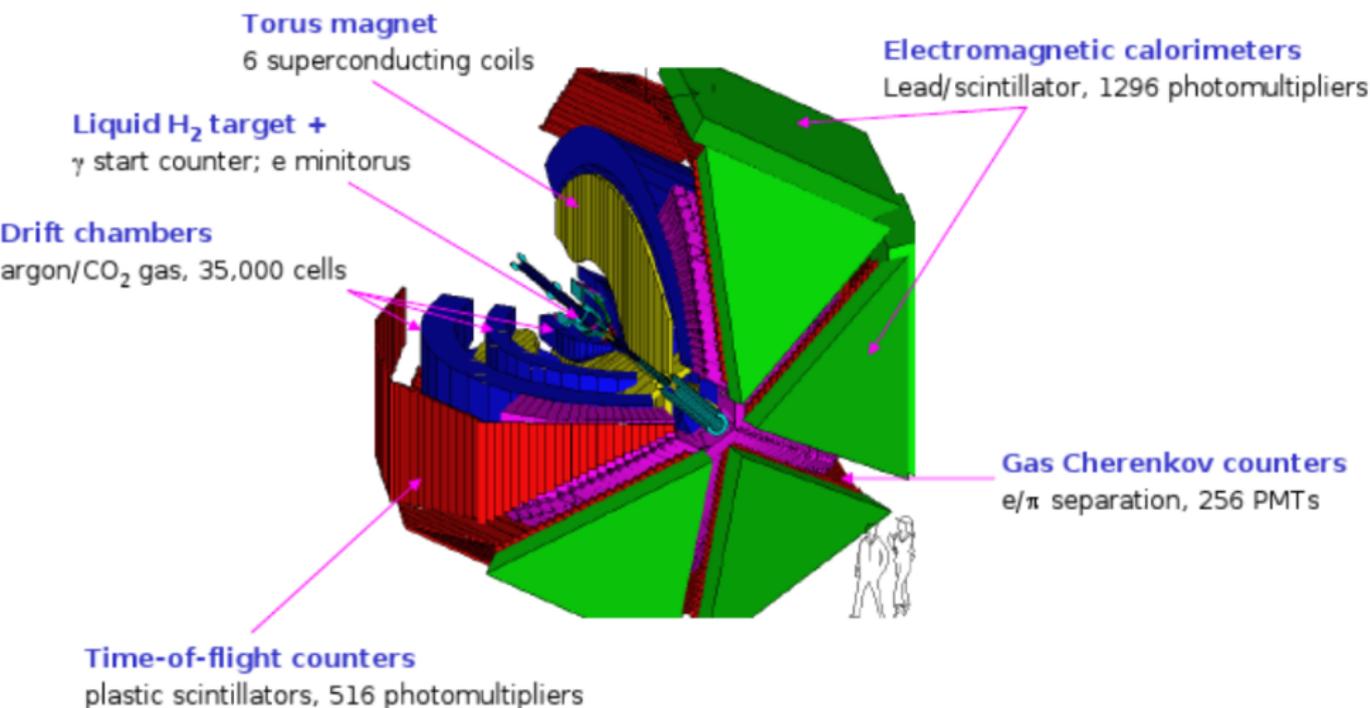
- ▶ Only 2 of these claimed $s\bar{s}$ observed
 - In e^+e^- exp. - $\phi(1680) \rightarrow KK^*$
 $\Gamma_{measured} = 150 \pm 50 \text{ MeV}/c^2$
 - In pK^- exp. - $\phi_3(1850) \rightarrow K^+K^-$
 $\Gamma_{measured} = 87^{+28}_{-23} \text{ MeV}/c^2$
- ▶ All existing $s\bar{s}$ observations in open strangeness decay modes
- ▶ Lack of observed $s\bar{s}$ states implies lots of potential for physics discovery

- Theoretical predictions for $s\bar{s}$ expected to be observed in the $\phi\eta$ invariant mass distribution from the reference - T. Barnes, N. Black and P. R. Page, Phys. Rev. D 68, 054014 (2003)



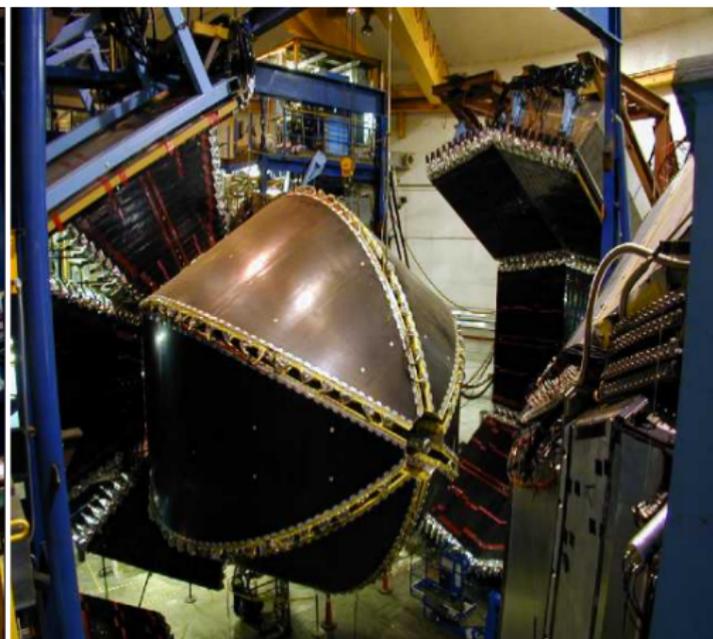
- ▶ **CEBAF:**
Continuous **E**lectron **B**eam **A**ccelerator **F**acility,
hosted at
Thomas Jefferson National Accelerator Facility,
Newport News, Virginia
- ▶ CEBAF delivers e^- beams to the 3 Halls,
polarised upon request in 5 passes with e^-
Energies up-to 6 GeV (1.2×5)
- ▶ Hall-B is the smallest experimental Hall with the
largest detector “CLAS”
- ▶ Major upgrades at CEBAF and the Halls for the 12
GeV upgrade as well as addition of a new Hall-D
which will house GLUEX created with meson
spectroscopy as the primary purpose

CEBAF Large Acceptance Spectrometer





- ◆ Skeletal superconducting Toroidal Magnets for CLAS.

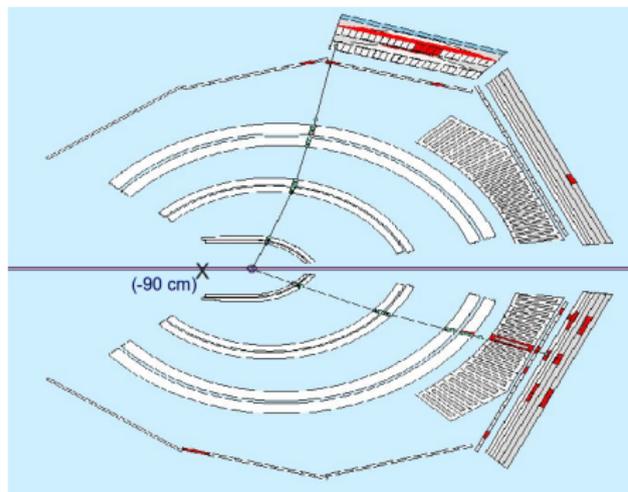


- ◆ CLAS detector during assembly.

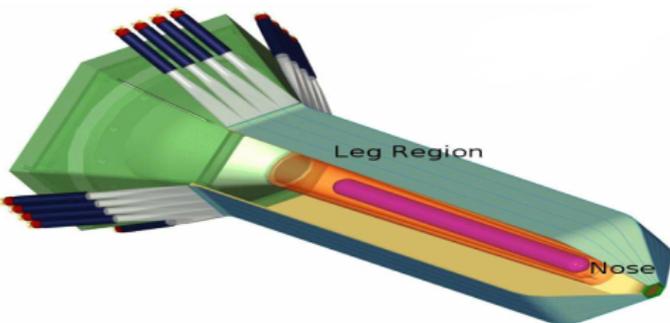
- ▶ 44.2 Days of beam-time over 70 days, 1st April to 9th June, 2008
- ▶ Beam current \rightarrow 60-65 nA ; $E_e \rightarrow$ 5.71 GeV ; DAQ Rate \rightarrow 8 KHz
- ▶ 26.2 billion triggers; Main Trigger \rightarrow 2 prong or more with $E_\gamma \geq 4.4$ GeV, 3 prong with no MOR ...
- ▶ 126 TB of raw data
- ▶ 250 TB of reconstructed data
- ▶ 68 pb^{-1} of photoproduction data

Detector Calibrations

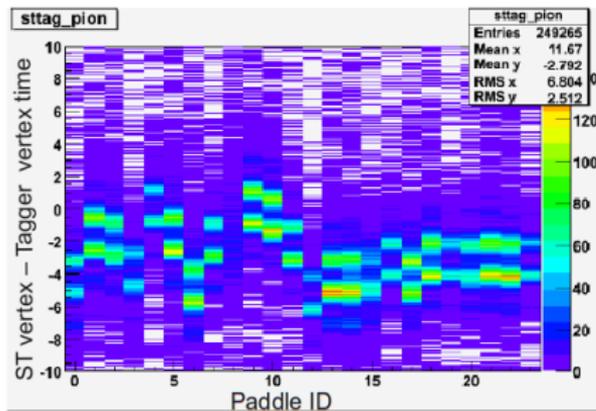
- ▶ Start Counter - Mukesh Saini
- ▶ Tagger - Mukesh Saini
- ▶ Time Of Flight - Craig Bookwalter
- ▶ Drift Chamber - Diane Schott
- ▶ Reconstruction - Johann Goetz



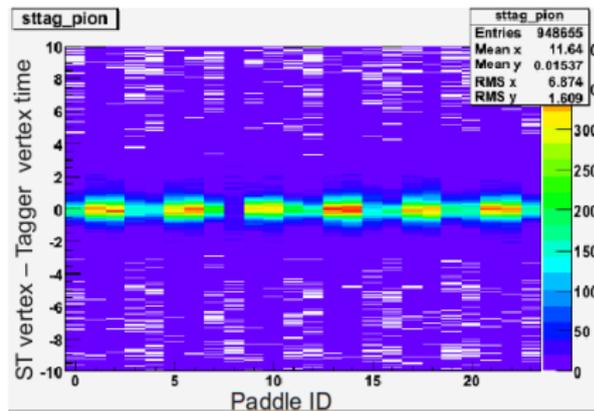
Start Counter – Tags the start time for a track



- ▶ Incorporates the independent sector based tracking of CLAS
- ▶ g12 pulled ST 90cm back from the center of CLAS to increase acceptance for low t , forward going particles
- ▶ ST crucial for picking the right photon and Particle ID



(a) Before



(b) After

Start Counter calibrated by Mukesh Saini

Tagger - Tags the Energy and the Timing for the incoming photon

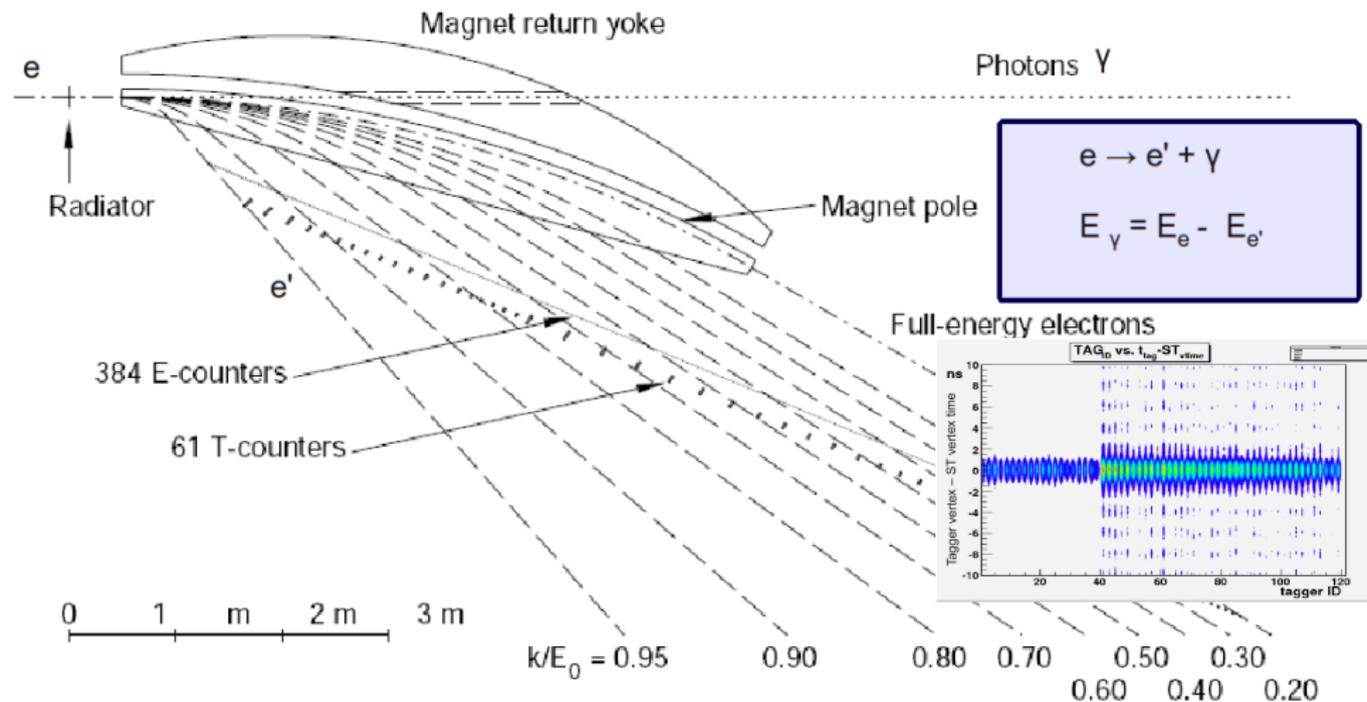


Fig. 23. Hall B photon-tagging system.

Tagging system calibrated by Mukesh Saini

Particle Selection

- ▶ 3 charged particle tracks detected in CLAS - Proton, K^+ , K^-
- ▶ Apply energy-momentum conservation - $\gamma_\mu + P_\mu = P'_\mu + K_\mu^+ + K_\mu^- + MM_\mu$
- ▶ Select - η - in the missing mass distribution
- ▶ Add four-vectors for K^+K^- to get their invariant mass and hence identify the - ϕ - meson
- ▶ Reconstruct the invariant mass for ($\phi\eta$) using the above selected ϕ & η mesons

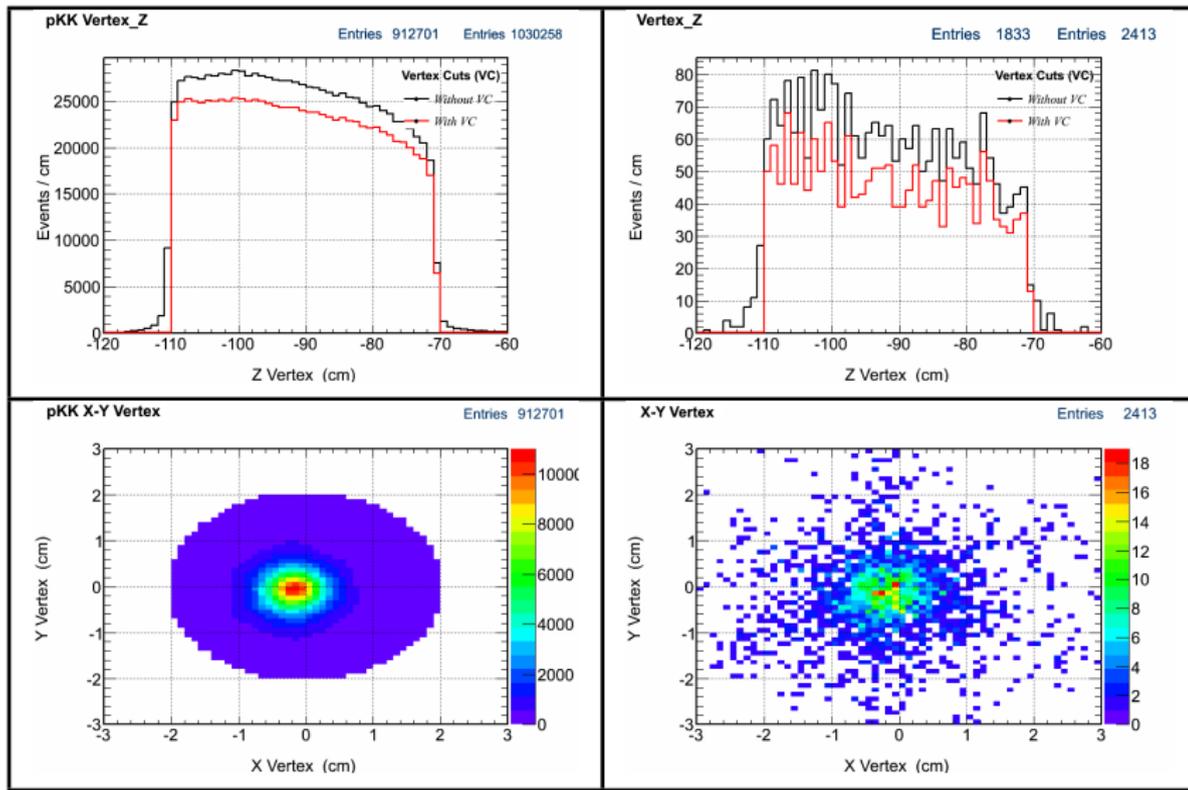
Meson ID

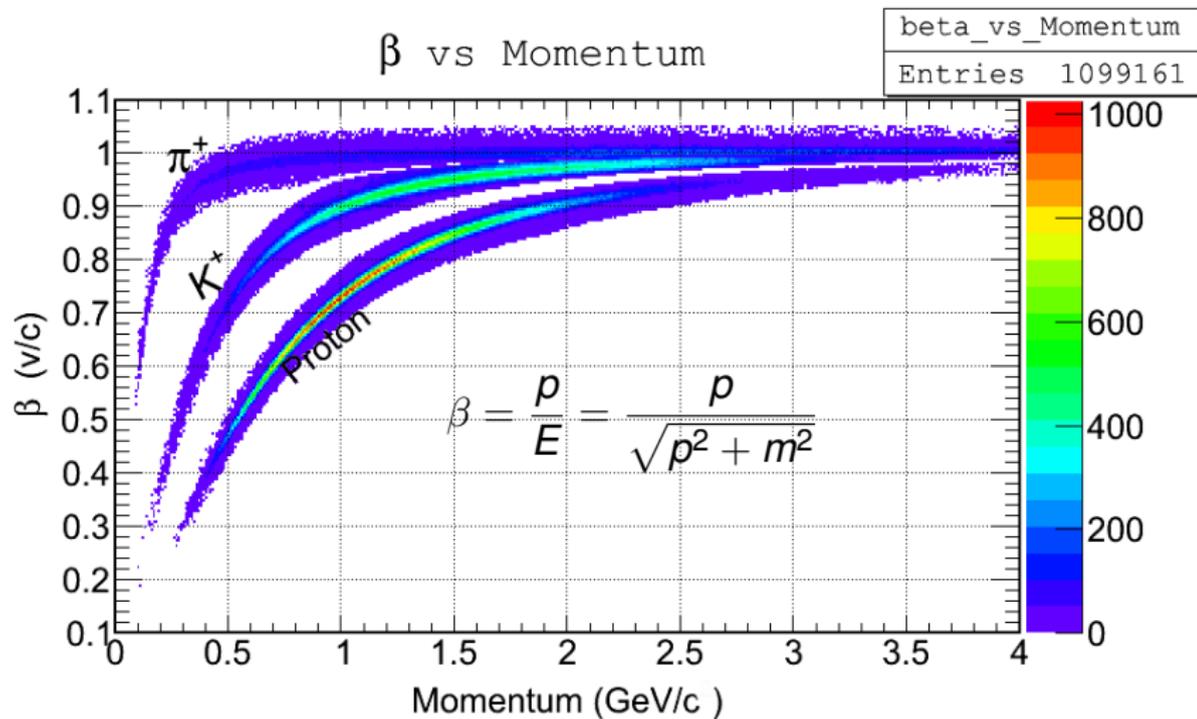
- ▶ To Identify ϕ , Select $1010 \text{ MeV} \leq \text{IM}(K^+ K^-) \leq 1030 \text{ MeV}$
- ▶ To identify η , Select $510 \text{ MeV} \leq \text{MM} \leq 580 \text{ MeV}$

Cuts

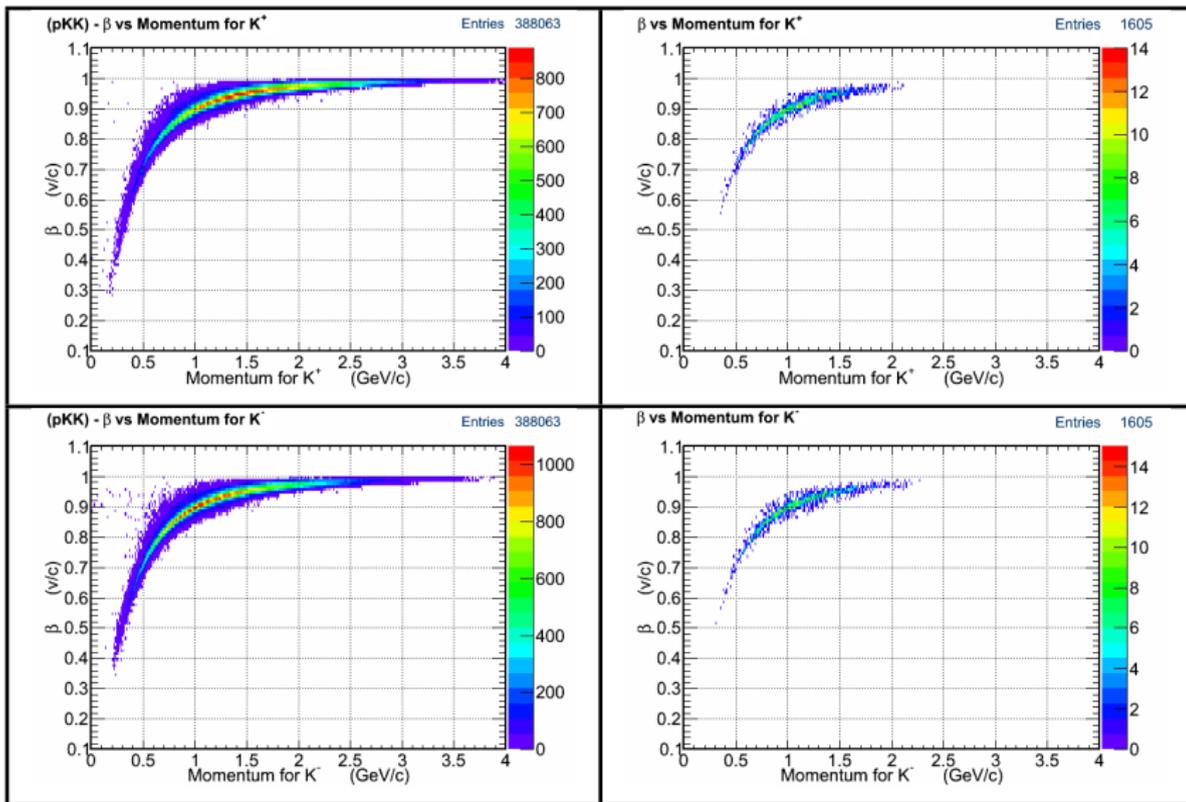
- ▶ Event vertex time within 1.002 ns of the Start Counter vertex time
- ▶ All particles have the difference between their measured and calculated ' β ' less than 0.05
- ▶ Event vertex is required to be within the Target

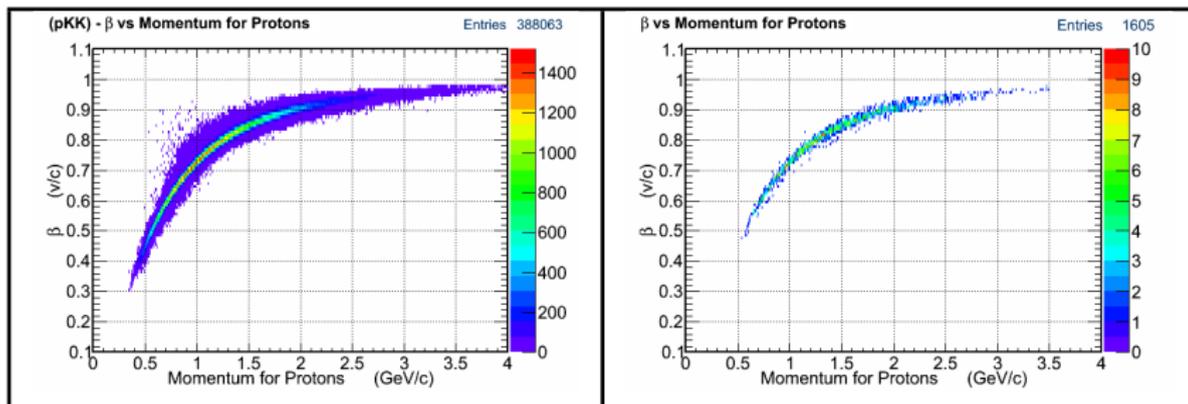
Vertex Distribution for Z and X-Y plane





β Vs Momentum Distributions



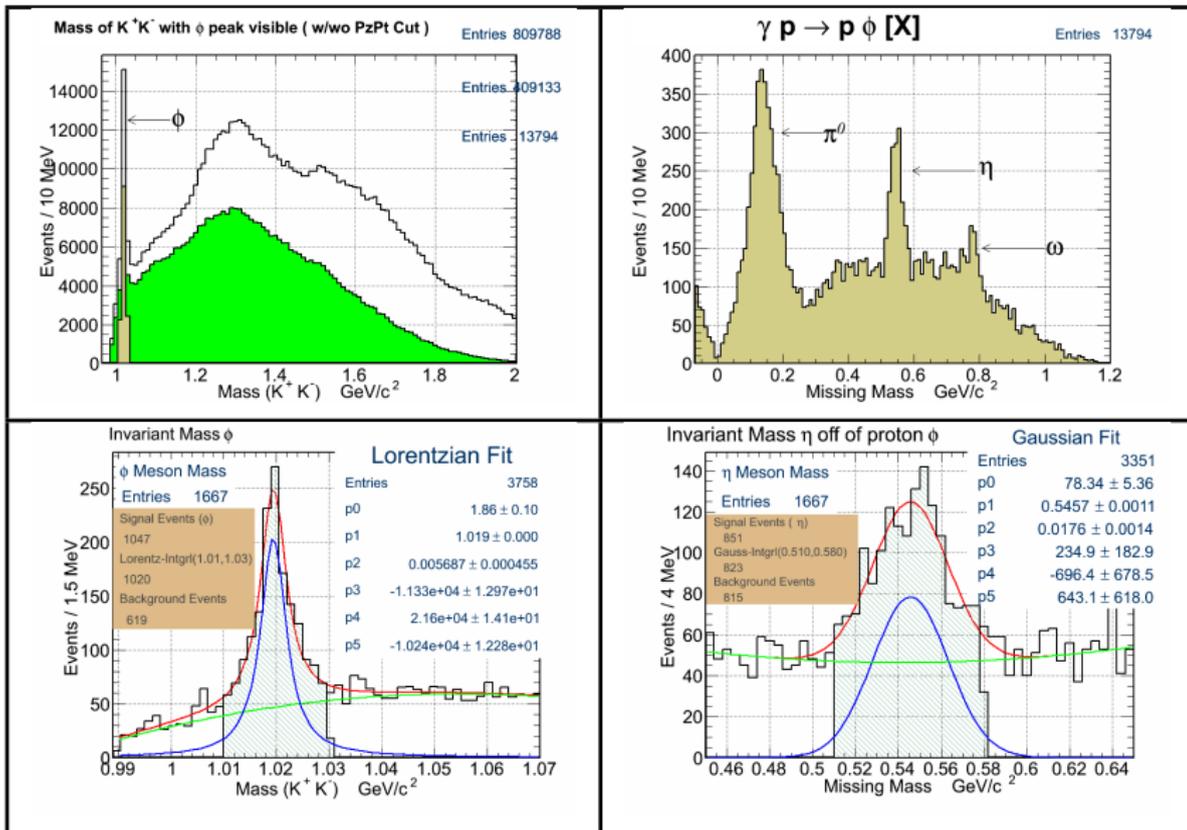


Observations

- ▶ Vertex and Timing Distributions are acceptable
- ▶ β for *Proton*, K^+ & K^- have no cross-contamination bands
- ▶ PID and Cuts employed work reasonably well



Selection Plots - ϕ η

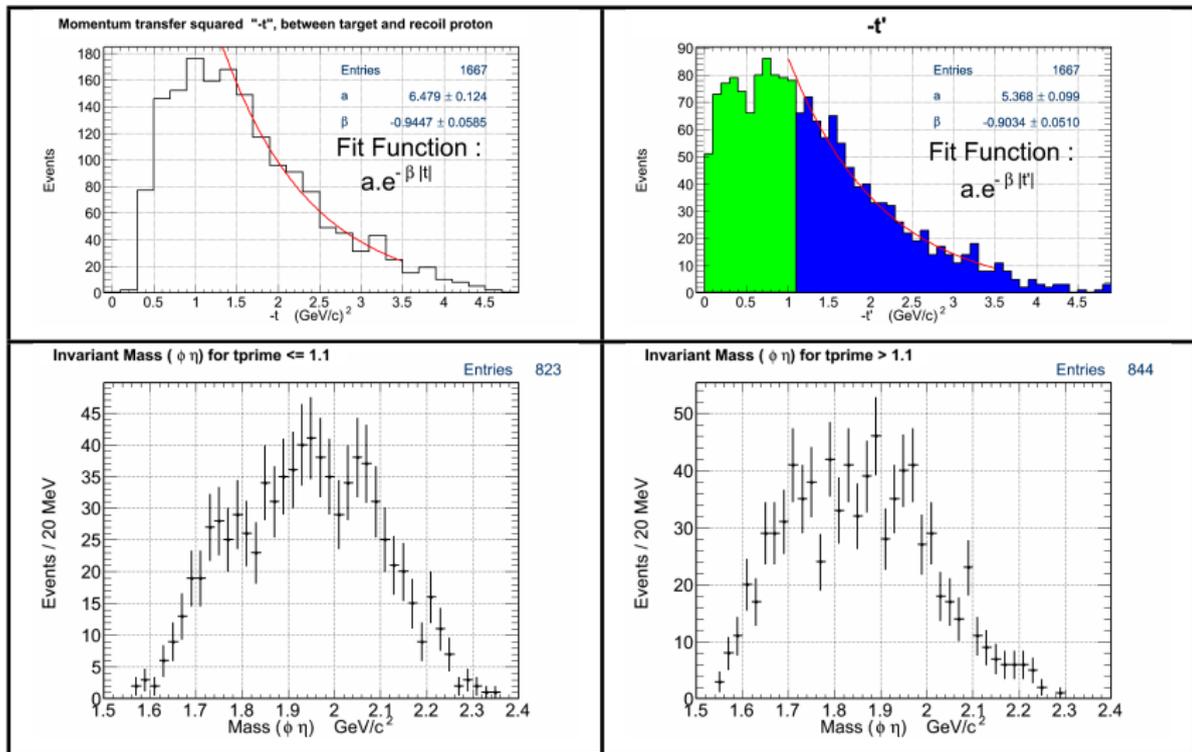


Selection Based On Kinematics Using t'

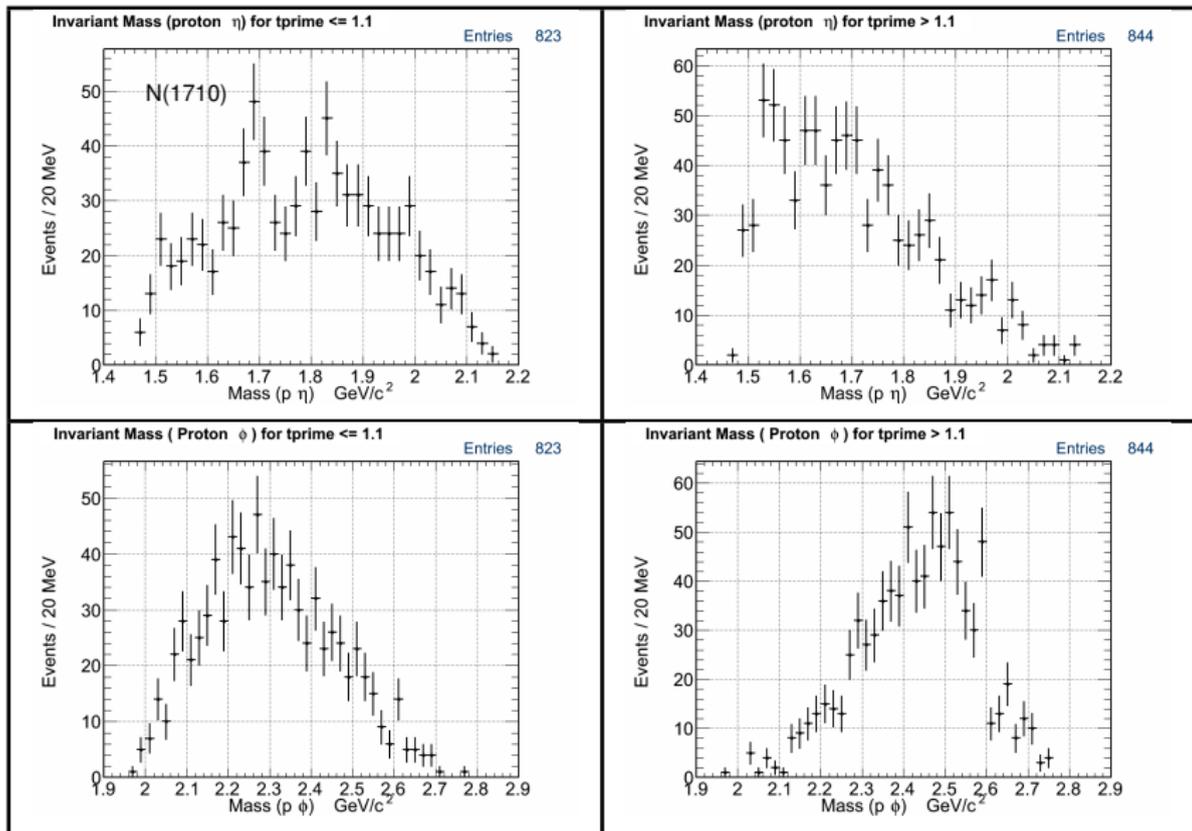
$$\text{Mandelstam's } t = |\mathbf{P}_\gamma^\mu - \mathbf{P}_X^\mu|^2 = |\mathbf{P}_{\text{target}}^\mu - \mathbf{P}_{\text{recoil}}^\mu|^2$$

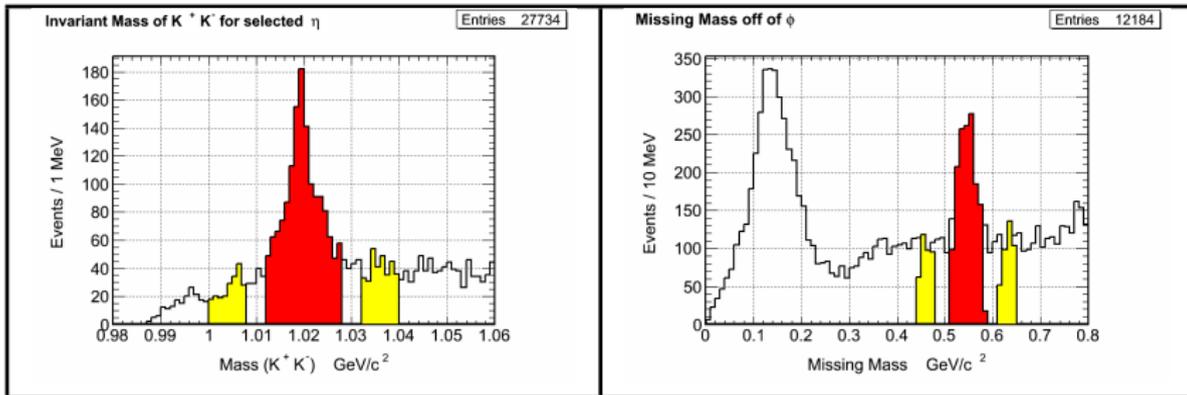
$$t' = t - t_{\text{min}}, \text{ where } t_{\text{min}} = \text{Minimum Four-Momentum Transfer Squared}$$

Required For Resonance Production



Selection Based On Kinematics



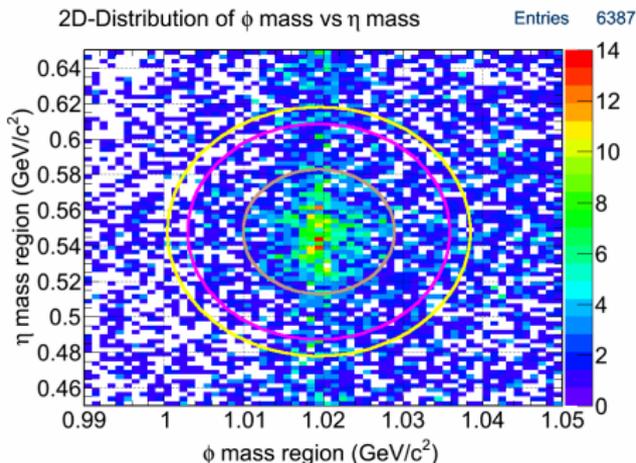


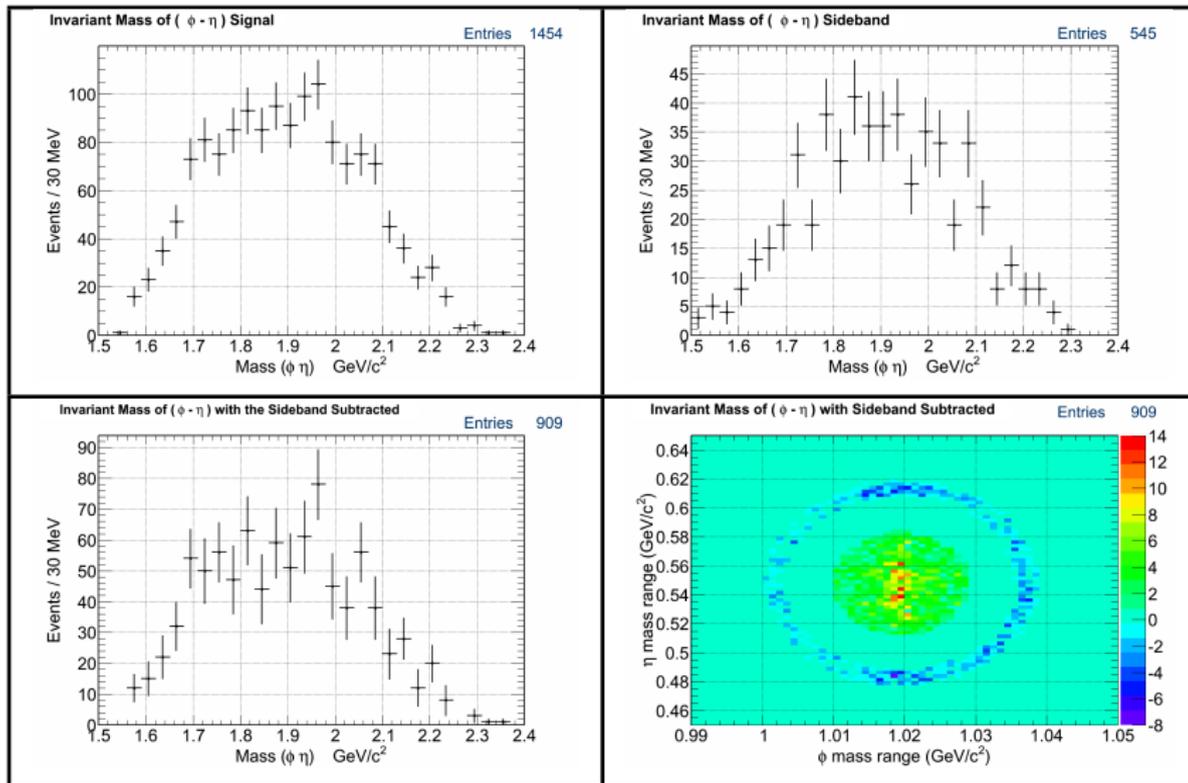
Side-Band Observations:

- ▶ ϕ width is 4 MeV and the peak is at 1020 MeV
- ▶ η width is 17 MeV and the peak is at 547 MeV
- ▶ Signal region is chosen to be peak $\pm 2\sigma$
- ▶ $\phi \rightarrow 1012 - 1028 \text{ MeV}$, $\eta \rightarrow 513 - 581 \text{ MeV}$
- ▶ Gap of 1σ is used between signal and sideband to minimize loss of ϕ 's
- ▶ Gap of 2σ is used between signal and sideband to minimize loss of η 's

$$\left(\frac{x - \phi_{\text{mass}}}{\frac{\phi \text{ mass range}}{2}} \right)^2 + \left(\frac{y - \eta_{\text{mass}}}{\frac{\eta \text{ mass range}}{2}} \right)^2 = 1$$

ϕ mass	1.0195 GeV/c ²
ϕ mass range selected	0.0019 GeV/c ²
ϕ sideband mass gap selected	$\sqrt{3} \times 0.0019$ GeV/c ²
ϕ sideband mass range selected	2×0.0019 GeV/c ²
η mass	0.5478 GeV/c ²
η mass range selected	0.070 GeV/c ²
η sideband mass gap selected	$\sqrt{3} \times 0.070$ GeV/c ²
η sideband mass range selected	2×0.070 GeV/c ²





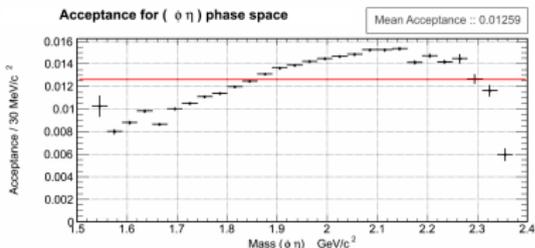
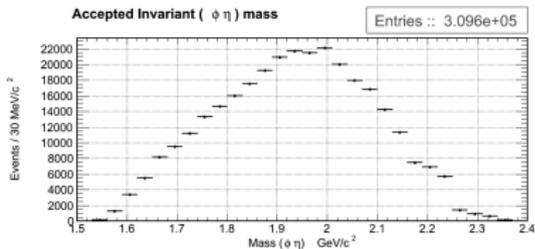
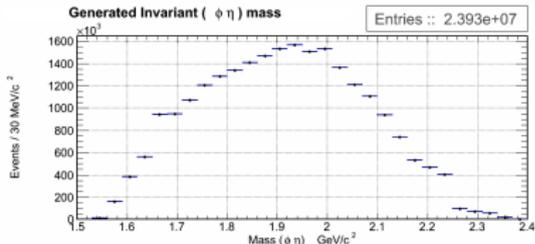
$$\frac{N_{target}}{V} = 2 \times N_{avogadro} \times \frac{\rho_{target}}{A_{H_2}}$$

$$\frac{N_{observed}}{Acceptance} = \sigma \times N_{incident} \times N_{target}$$

$$\sigma = \frac{N_{observed} \times A_{H_2}}{2 \times N_{incident} \times N_{avogadro} \times \rho_{target} \times L_{target} \times Acceptance}$$

where,

- N_{target} is the number of target protons that on average lie in the path of the incoming beam photons,
- $N_{observed}$ is the number of observed events in the experiment, aka the yield,
- ρ_{target} is the density of the LH₂ used in the experiment,
- $N_{incident}$ is the integrated flux (total number) of the incoming beam photons that were incident on the target to achieve the observed yield,
- L_{target} is the length of the target cell,
- Acceptance is derived from the simulations of the reaction phase space for the experiment and represents the corrections due to the finite acceptance of the detector.

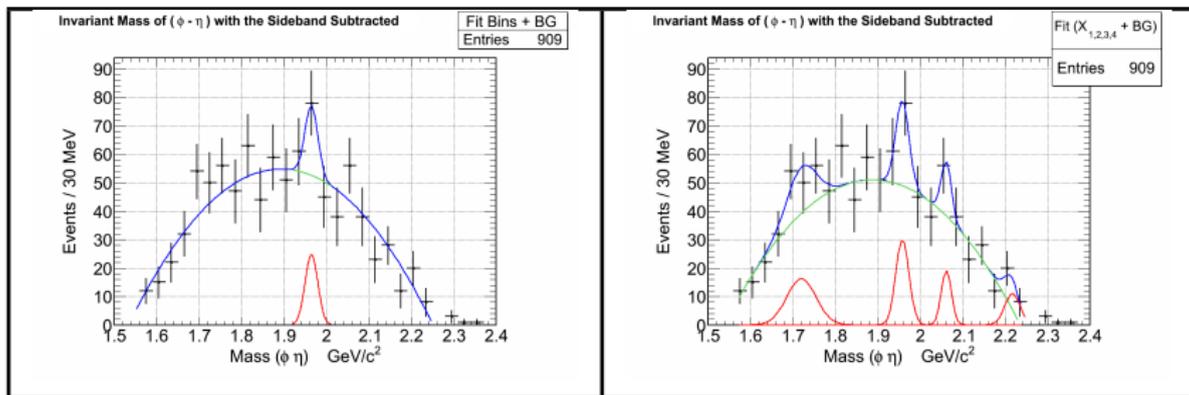


Measured Yield of $\phi \eta$ events	909 ± 45
Reconstruction inefficiency at 60 nA	16%
ST inefficiency per charged track	6%
Branching ratio $\phi \rightarrow K^+ K^-$	$48.5\% \pm 0.5\%$
Corrected Yield of $\phi \eta$ events	2686 ± 133
E_γ Flux (10^{13})	2.18 ± 0.17
Overall phase space $\phi \eta$ MC Acceptance	1.26%
Systematic error - $\phi \eta$ simulations	7.9%
Systematic error - Sideband Subtraction	0.8%
Systematic error - PID	1.3%
Systematic error - Photon Flux Normalisation	7.9%
Systematic error - Miscellaneous	0.26%
Conservative total systematic error	12.5%

Cross section - $\sigma_{\chi \rightarrow \phi \eta}$ for $E_\gamma \in \{4.40, 5.45\}$

► 5.8 ± 0.3 (stat) ± 0.72 (sys) nb

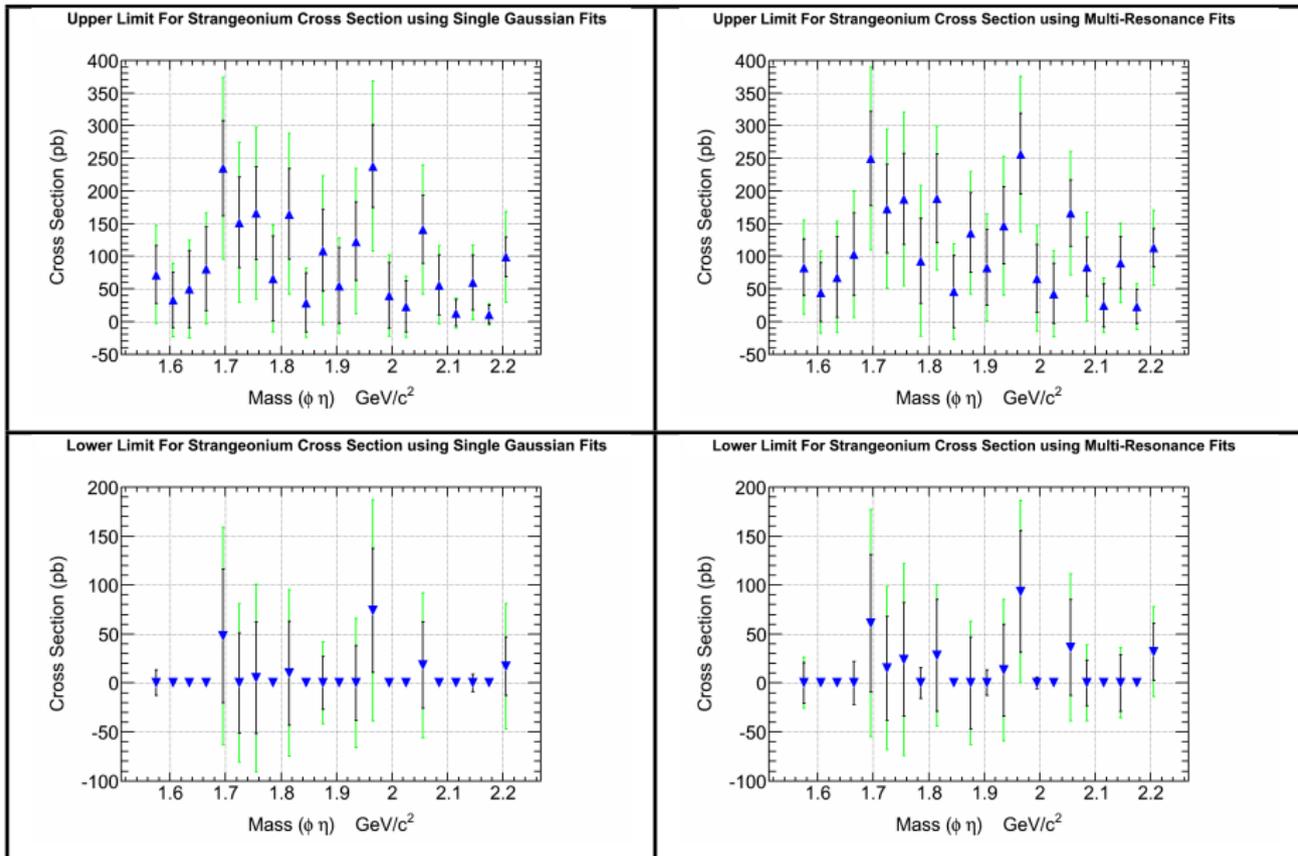
Background Estimate For Feldman-Cousins Method



Feldman-Cousins Method

- ▶ Establishing discovery requires rejecting the background only hypothesis
- ▶ Particle physics experiments - 3σ signal significance \rightarrow evidence, 5σ signal significance \rightarrow discovery
- ▶ Feldman-Cousins method - a frequentist approach, uses relative frequency of an event
 - Perform multiple experiments and number of positive results determines probability of that result
- ▶ It uses only the observed number of events and the estimated background count to calculate the confidence limits
- ▶ The fixed-unknown-true value will fall within the confidence interval in 90% of the repeats of the same experiment

Strangeonium Cross Section Confidence Limits



- ▶ The two known and expected resonances - $\phi(1680)$ and $\phi_3(1850)$ - were not observed in the $\phi \eta$ invariant mass spectrum
- ▶ Largest cross section upper limit for strangeonium decay - ' $s\bar{s} \rightarrow \phi \eta$ '
 - 1695 MeV/c² - 250 pb
 - 1965 MeV/c² - 250 pb
- ▶ Calculated lower limits for $s\bar{s} \rightarrow \phi \eta$ consistent with zero

Comment on X(1750)

- ▶ Omega spectrometer measured X(1750) $\rightarrow K^+K^-$ photoproduction cross section at $E_\gamma = 45 \text{ GeV}/c^2$ to be $8 \pm 3 \text{ nb}$
- ▶ 3P_0 model predicts the branching ratio of $K^+K^- : \phi \eta$ to be 2 : 1
- ▶ Extrapolating using froissart bound, the state is expected at least at a level of 1 nb
- ▶ If photoproduced X(1750) $\rightarrow K^+K^- \in \{s\bar{s}\}$, Our analysis would have observed the state in $\phi \eta$ at a level well above 250 pb
- ▶ The analysis thus supports FOCUS's claim that X(1750) $\neq \phi(1680)$

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- ▶ If photoproduced X(1750) $\rightarrow K^+K^- \in \{s\bar{s}\}$, Our analysis would have observed the state in $\phi \eta$ at a level well above 250 pb
- ▶ The analysis thus supports FOCUS's claim that X(1750) $\neq \phi(1680)$

- ▶ The two known and expected resonances - $\phi(1680)$ and $\phi_3(1850)$ - were not observed in the $\phi \eta$ invariant mass spectrum
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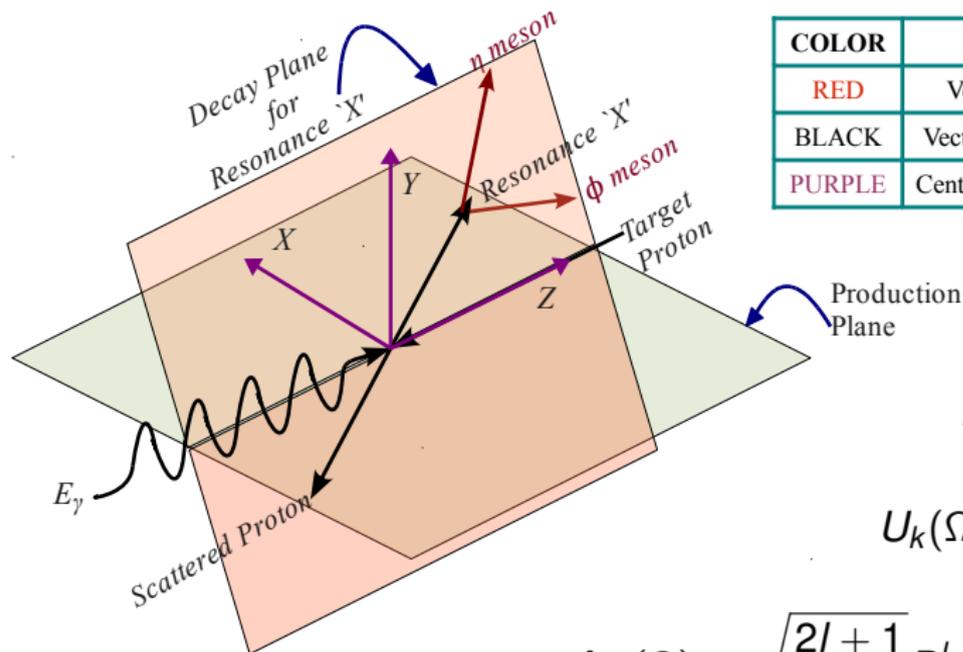
Comment on X(1750)

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COLOR	ATTRIBUTE
RED	Vectors In Decay Plane
BLACK	Vectors In Production Plane
PURPLE	Center of Mass Frame's Axes

$$I(\Omega) = \sum_k |U_k(\Omega)|^2$$

$$U_k(\Omega) = \sum_{lm} V_{lmk} A_{lm}(\Omega)$$

$$A_{lm}(\Omega) = \sqrt{\frac{2l+1}{4\pi}} D_{m0}^l(\phi, \theta, 0) = Y_l^m(\Omega)$$

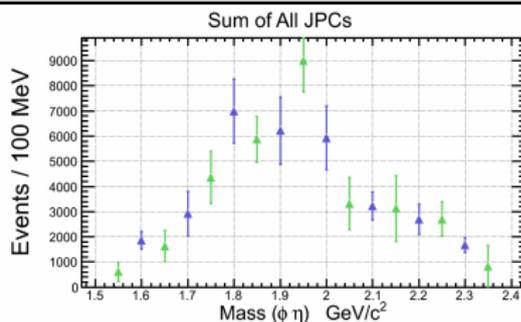
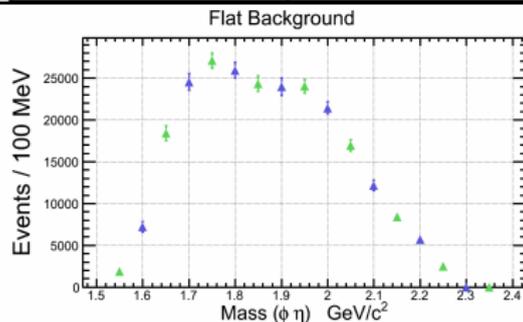
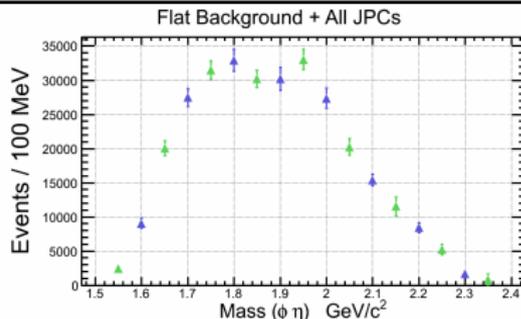
$$I(\tau) \propto \sum_k \left| \sum_{\alpha} V_{k\alpha} A_{\alpha}(\tau) \right|^2$$

$$\alpha \in \{J, P, C, M, L, I, (w, \Gamma)\}$$

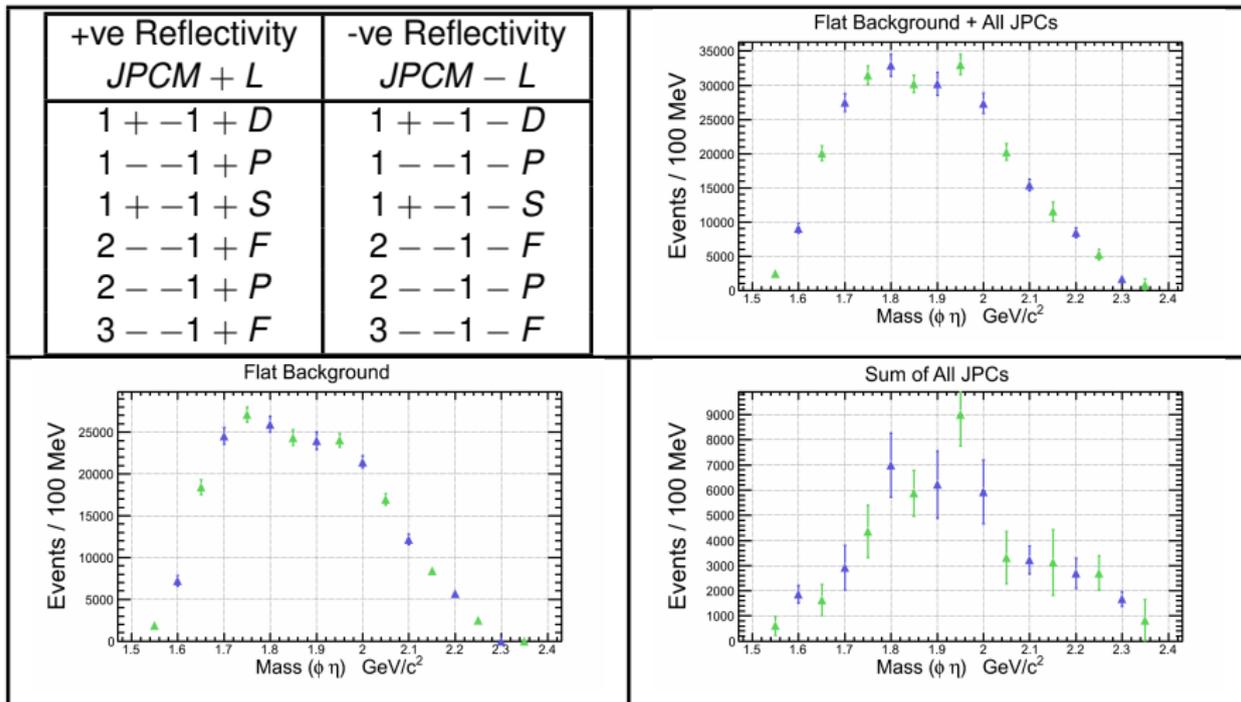
Fourier Expansion :

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} (a_n \sin(nx) + b_n \cos(nx))$$

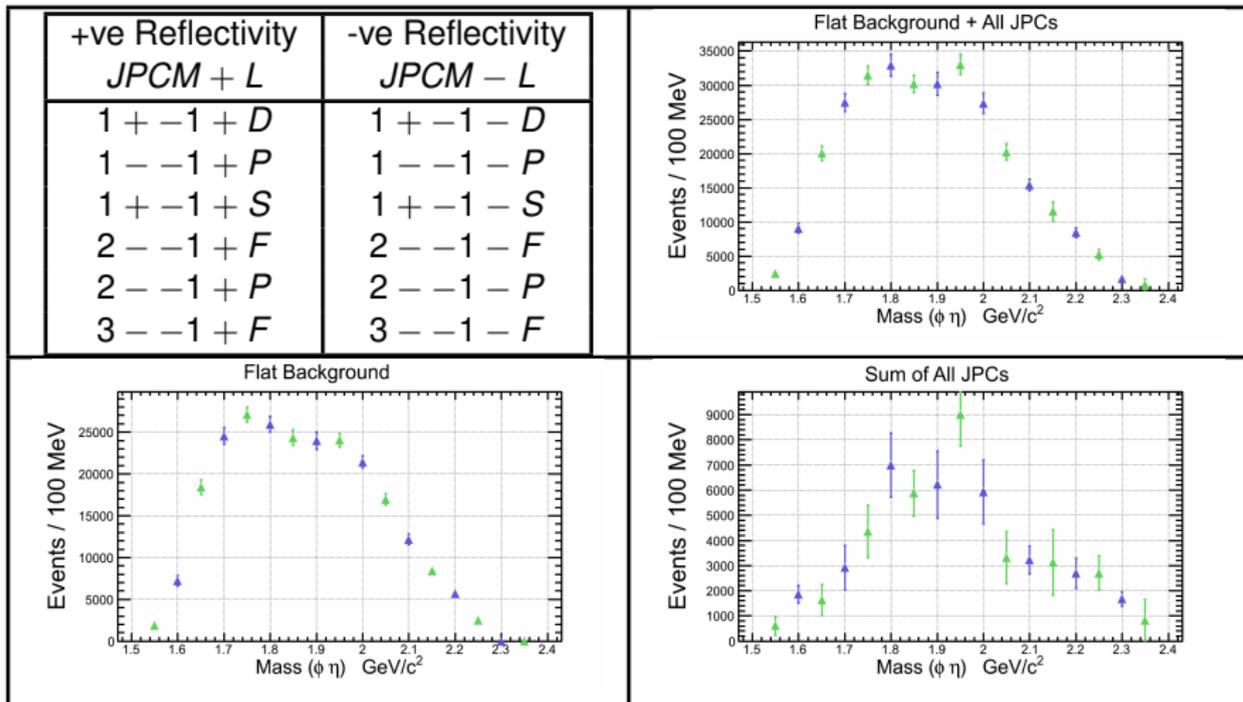
+ve Reflectivity <i>JPCM + L</i>	-ve Reflectivity <i>JPCM - L</i>
1 + - 1 + <i>D</i>	1 + - 1 - <i>D</i>
1 - - 1 + <i>P</i>	1 - - 1 - <i>P</i>
1 + - 1 + <i>S</i>	1 + - 1 - <i>S</i>
2 - - 1 + <i>F</i>	2 - - 1 - <i>F</i>
2 - - 1 + <i>P</i>	2 - - 1 - <i>P</i>
3 - - 1 + <i>F</i>	3 - - 1 - <i>F</i>



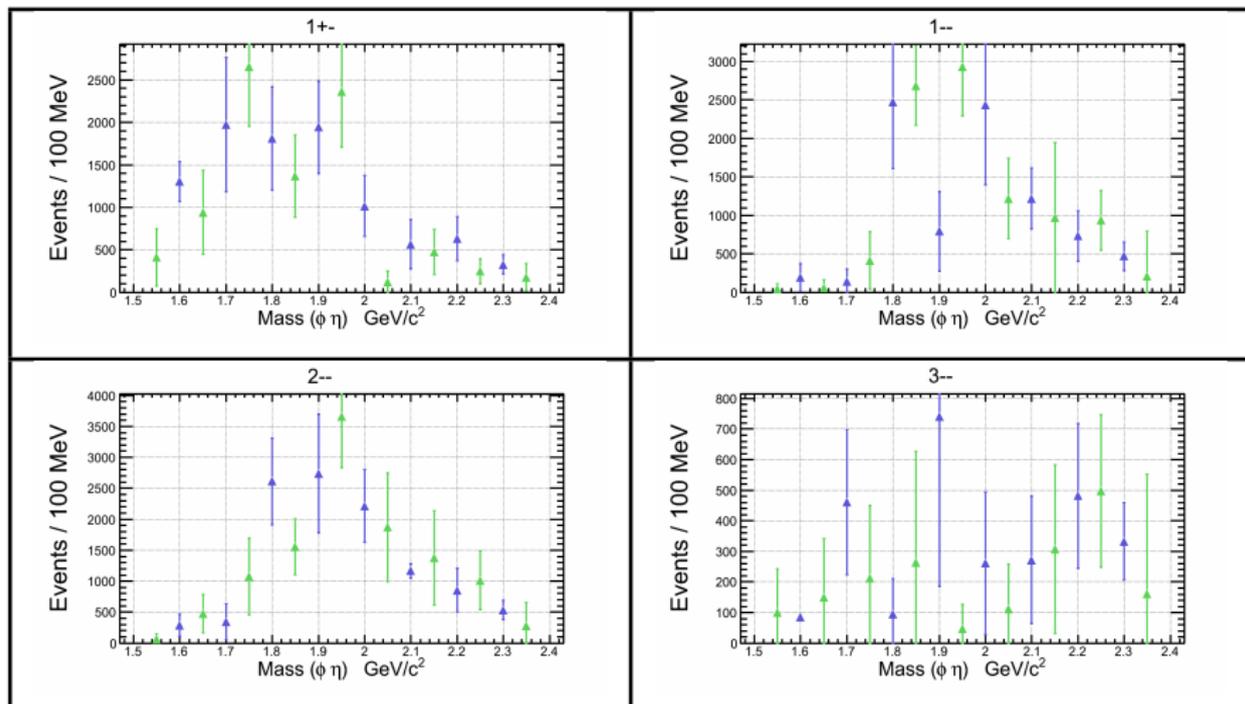
- Displayed fits are for $\phi \eta$ data binned in 100 MeV bins - then rebinned again in 100 MeV bins with an offset of 50 MeV - color coded by green and blue
- Data four-vectors used for PWA were kinematically fitted to a missing η
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- ▶ 2 — — wave is stable; but not much stronger than any other JPC
- ▶ No JPC is singularly preferred in the PWA analysis; PWA inconclusive about the makeup of the observed spectra

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 - Upper limit for $\phi(1680)$ - 140 pb
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 - Largest cross section upper limit is observed to be 250 pb
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Thank You!

$$|\bar{p}\epsilon PJM\rangle = \theta(m) \{ |\bar{p}PJM\rangle + \epsilon P(-1)^{J-M} |\bar{p}PJ - M\rangle \}$$

where,

$$\begin{aligned} \theta(m) &= \frac{1}{\sqrt{2}} \text{ for } m > 0; \quad \theta(m) = \frac{1}{2} \text{ for } m = 0; \\ \theta(m) &= 0 \text{ for } m < 0. \end{aligned}$$

Reflectivity (ϵ) / Naturality (N) for a particle, in this case, by definition is:

$$\epsilon = N = P(-1)^J$$

hence the following holds true,

$$\begin{array}{ll} \text{Natural parity exchange } \epsilon = 1; & J^P = 0^+, 1^-, 2^+, \dots \\ \text{Unnatural parity exchange } \epsilon = -1; & J^P = 0^-, 1^+, 2^-, \dots \end{array}$$

As an added advantage, states of different reflectivities / naturalities do not interfere.

$${}^\epsilon U_k(\Omega) = \sum_{lm} {}^\epsilon V_{lmk} \sqrt{\frac{2l+1}{4\pi}} {}^\epsilon D_{m0}^{l*}(\phi, \theta, 0)$$

$$I(\Omega) = \sum_{\epsilon k} |{}^\epsilon U_k(\Omega)|^2 = |{}^+ U_1(\Omega)|^2 + |{}^- U_1(\Omega)|^2 + |{}^+ U_2(\Omega)|^2 + |{}^- U_2(\Omega)|^2$$

$$\mathcal{M} = \sum_{\alpha, \phi} \underbrace{\langle K^+ K^- \eta \rho | \hat{T}_d^{\phi \rightarrow K^+ K^-} | \phi \eta \rho \rangle}_{\text{Decay} - A_\alpha(\tau)} \underbrace{\langle \phi \eta \rho | \hat{T}_d^{X \rightarrow \phi \eta} | X_\alpha \rho \rangle}_{\text{Production} - V_\alpha} \langle X_\alpha \rho | \hat{T}_p | \gamma \rho \rangle$$

$$I(\tau) \propto \sum_k \left| \sum_\alpha V_{k\alpha} A_\alpha(\tau) \right|^2$$

$$\alpha \in \{J, P, C, M, L, l, (w, \Gamma)\}$$

where,

- J = Total angular momentum of resonance 'X'
- P = Parity of the resonance 'X'
- C = Charge conjugation parity for the resonance 'X'
- M = Z-projection of the total angular momentum
- L = Angular momentum between the ϕ and the η meson
- l = Angular momentum between the decay products of ϕ
- (w, Γ) = Mass and width parameter for the Breit-Wigner

$$\mathcal{L} \propto \left[\frac{\bar{n}^n}{n!} e^{-\bar{n}} \right] \prod_i^n \left[\frac{l(\tau_i)}{\int l(\tau)\eta(\tau)pq d\tau} \right]$$

$$\bar{n} \propto \int l(\tau)\eta(\tau)pq(d\tau)$$

$$\ln \mathcal{L} \propto \sum_i^n \ln(l(\tau_i)) - \int l(\tau)\eta(\tau)pq d\tau$$

$$\propto \sum_{k \in \alpha\alpha'} \ln({}^\epsilon V_{\alpha k} {}^\epsilon V_{\alpha' k}^* {}^\epsilon A_{\alpha}(\tau_i) {}^\epsilon A_{\alpha'}^*(\tau_i)) - \eta_x \left(\sum_{k \in \alpha\alpha'} {}^\epsilon V_{\alpha k} {}^\epsilon V_{\alpha' k}^* {}^\epsilon \psi_{\alpha\alpha'}^a \right)$$

$$\eta_x = \frac{\text{MC Events Accepted}}{\text{Raw MC Events Generated}}$$

$${}^\epsilon \psi_{\alpha\alpha'}^a = \frac{1}{N_a} \sum_i^{N_a} {}^\epsilon A_{\alpha}(\tau_i) {}^\epsilon A_{\alpha'}^*(\tau_i)$$

$$N = \sum_{k \in \alpha\alpha'} {}^\epsilon V_{\alpha k} {}^\epsilon V_{\alpha' k}^* {}^\epsilon \psi_{\alpha\alpha'}^r$$

where,

- \mathcal{L} is the likelihood function,
- n is the number of events observed,
- \bar{n} is the average number of events observed if the experiment was ran multiple times,
- $\eta(\tau)$ is the finite experimental acceptance as determined by the Monte Carlo simulations,
- $pq d\tau$ is the lorentz invariant phase space element for the involved kinematics,
- ${}^\epsilon \psi_{\alpha\alpha'}^a$ is the normalization integral calculated from phase space MC simulation for N_a accepted events.

$$\gamma + \textit{proton} \rightsquigarrow \textit{proton} + \phi + \pi^0$$

$$\gamma + \textit{proton} \rightsquigarrow \textit{proton} + K^+ + K^- + [\pi^0]$$

