



Beam Asymmetry Measurements from g8b for $\gamma p \rightarrow p \pi^+ \pi^-$

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On behalf of the CLAS Collaboration
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Tallahassee, FL



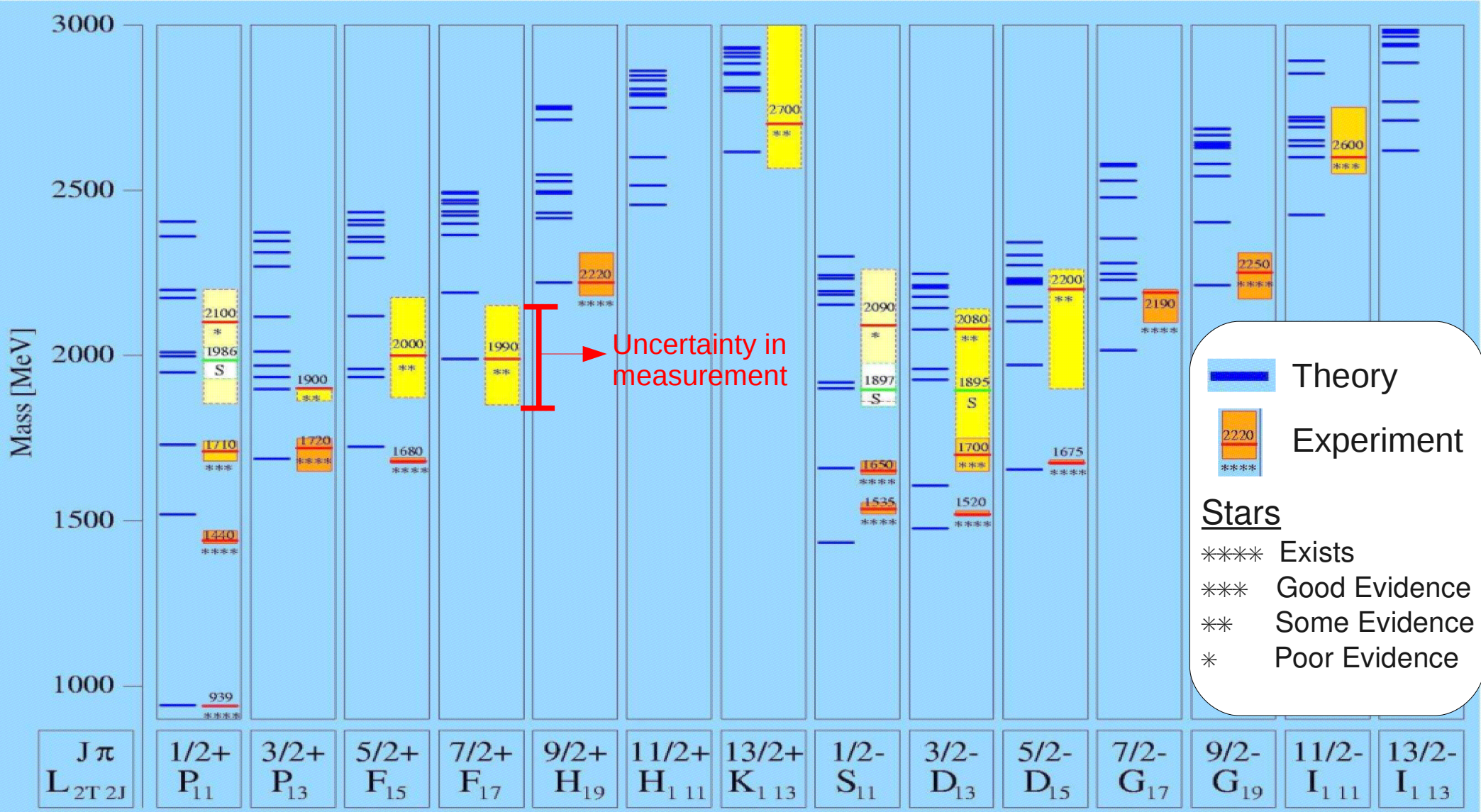
Outline

- Introduction:
 - Problems in Hadron Spectroscopy
- Experimental Facility
- Analysis
 - Final State Equations
 - Coordinate system
 - Preliminary Results from g8b (first measurements of I^S and I^C for $\gamma p \rightarrow p \pi^+ \pi^-$)

Constituent Quark Model

U. Loring

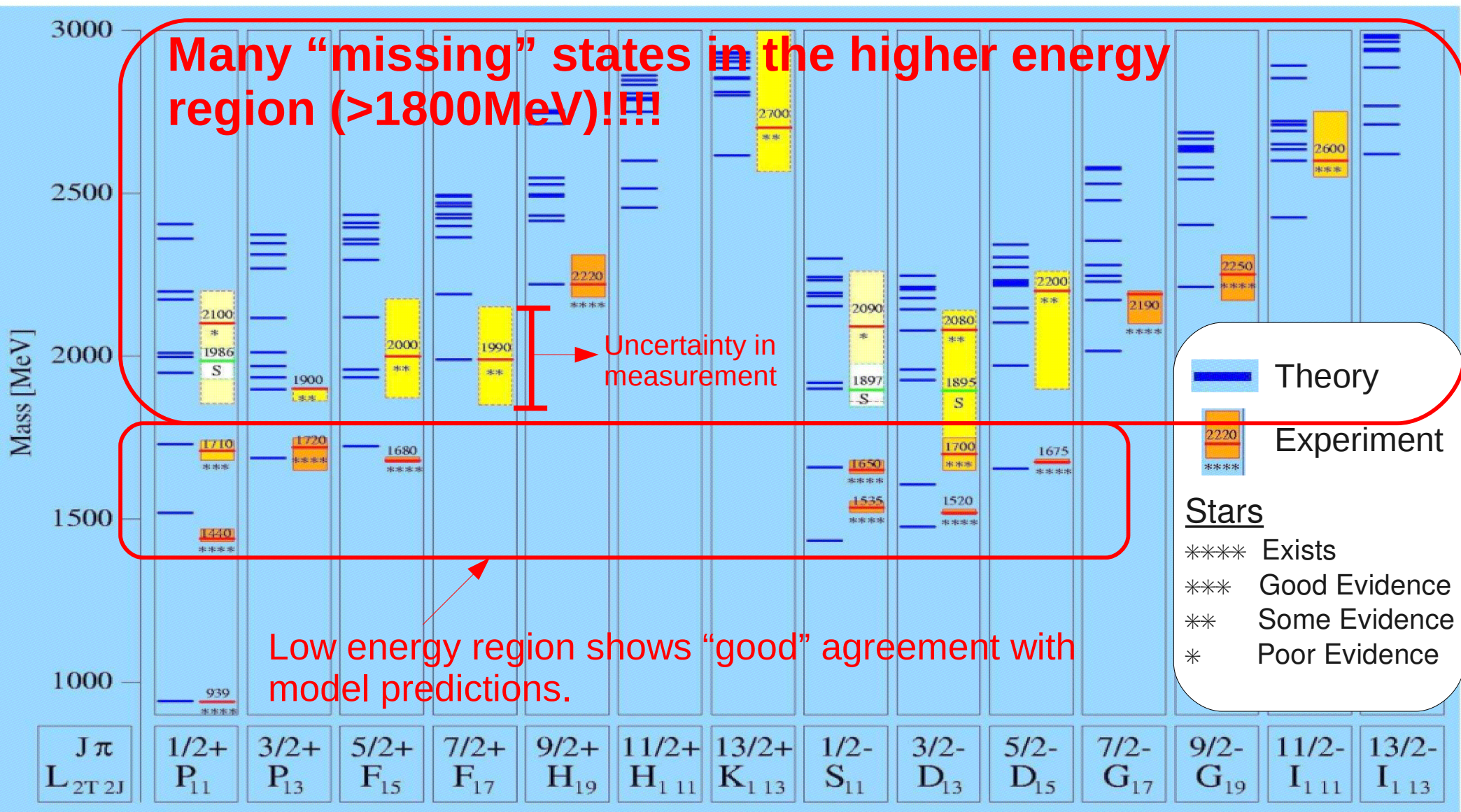
N^* Resonances (Isospin = 1/2 states)



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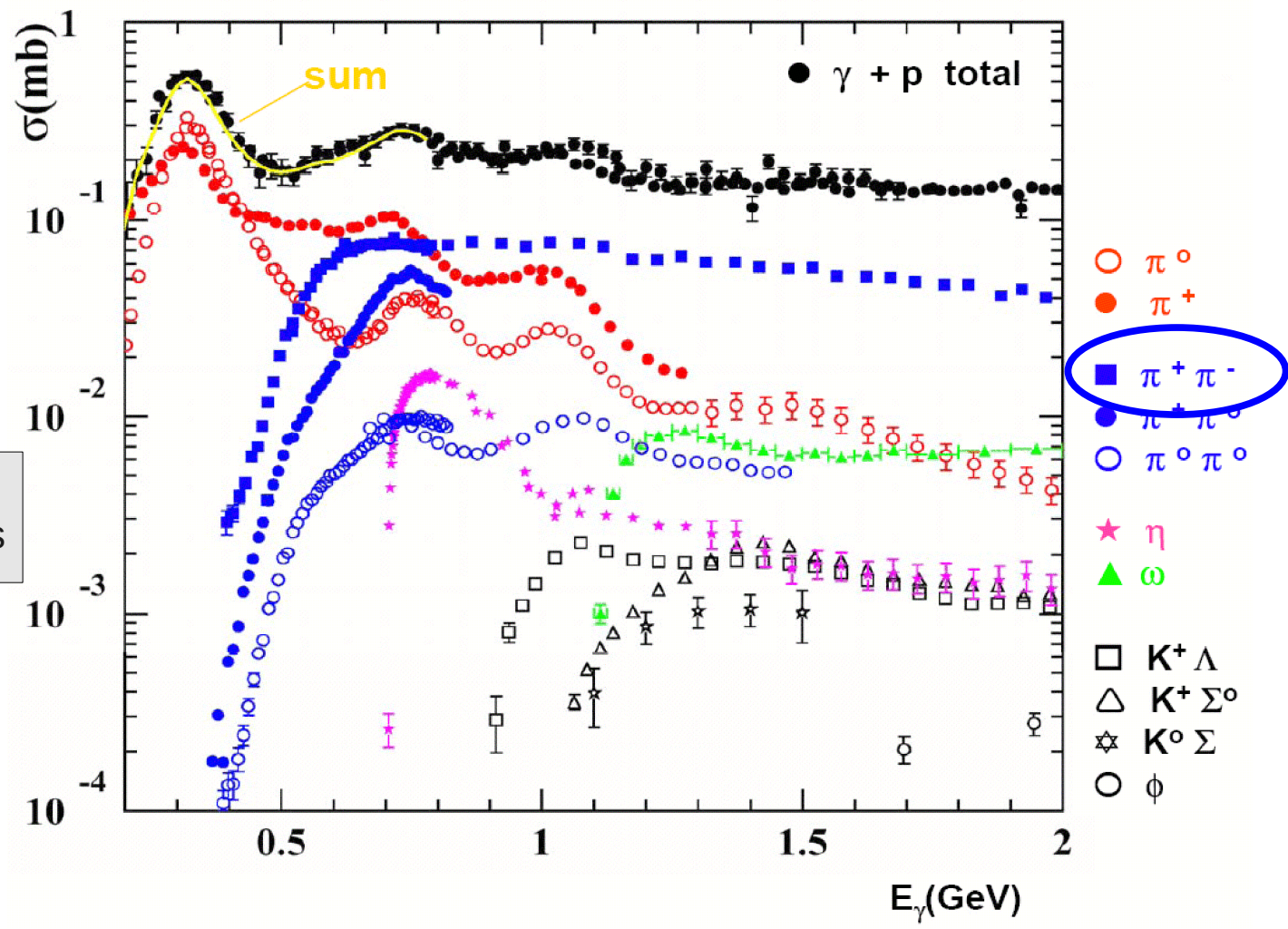




Why the lack of evidence? What can we do?

- Most of the existing data regarding N^* resonances involves πN elastic scattering.
 - Photoproduction (γp) is predicted to be a promising method for producing these states.
- Existing polarized photoproduction data mainly covers a mass range up to about 1800 MeV with analyses involving a single meson.
 - Go higher in energy; analyze a channel with two or more mesons as it accounts for most of the cross section for $W \geq 2$ GeV
- Analysis of unpolarized data leads to ambiguous results.
 - The inclusion of polarization avoids such ambiguities.

$\gamma+p$ Cross Section



Experimental Cross Sections

- $\pi^+\pi^-$ cross section is large but devoid of structure. \rightarrow resonances are broad and overlapping.
- Cannot go “peak hunting”.
- Constraints provided by polarization are needed to disentangle and isolate resonance contributions.

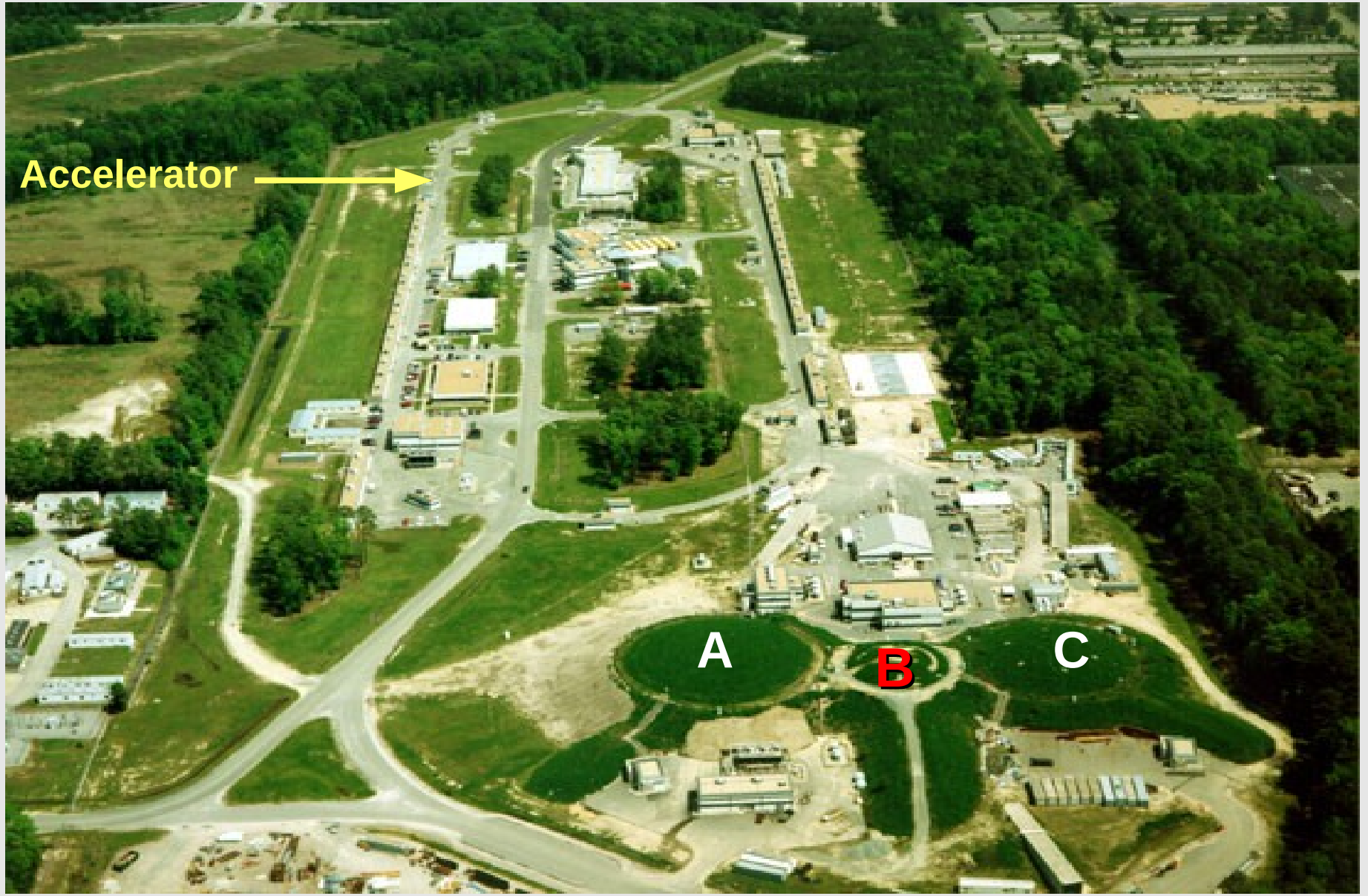


The Facility: Jefferson Lab in Newport News, VA



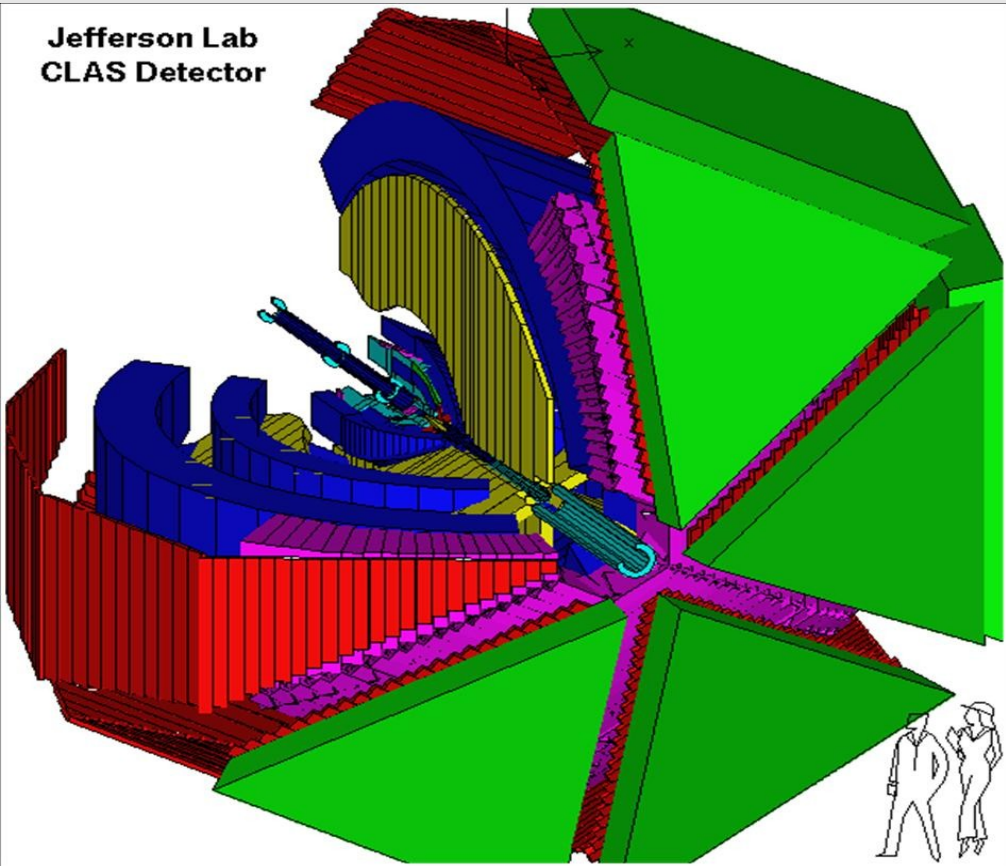


The Facility: Jefferson Lab in Newport News, VA



The Hall B Detector : CLAS

Jefferson Lab
CLAS Detector

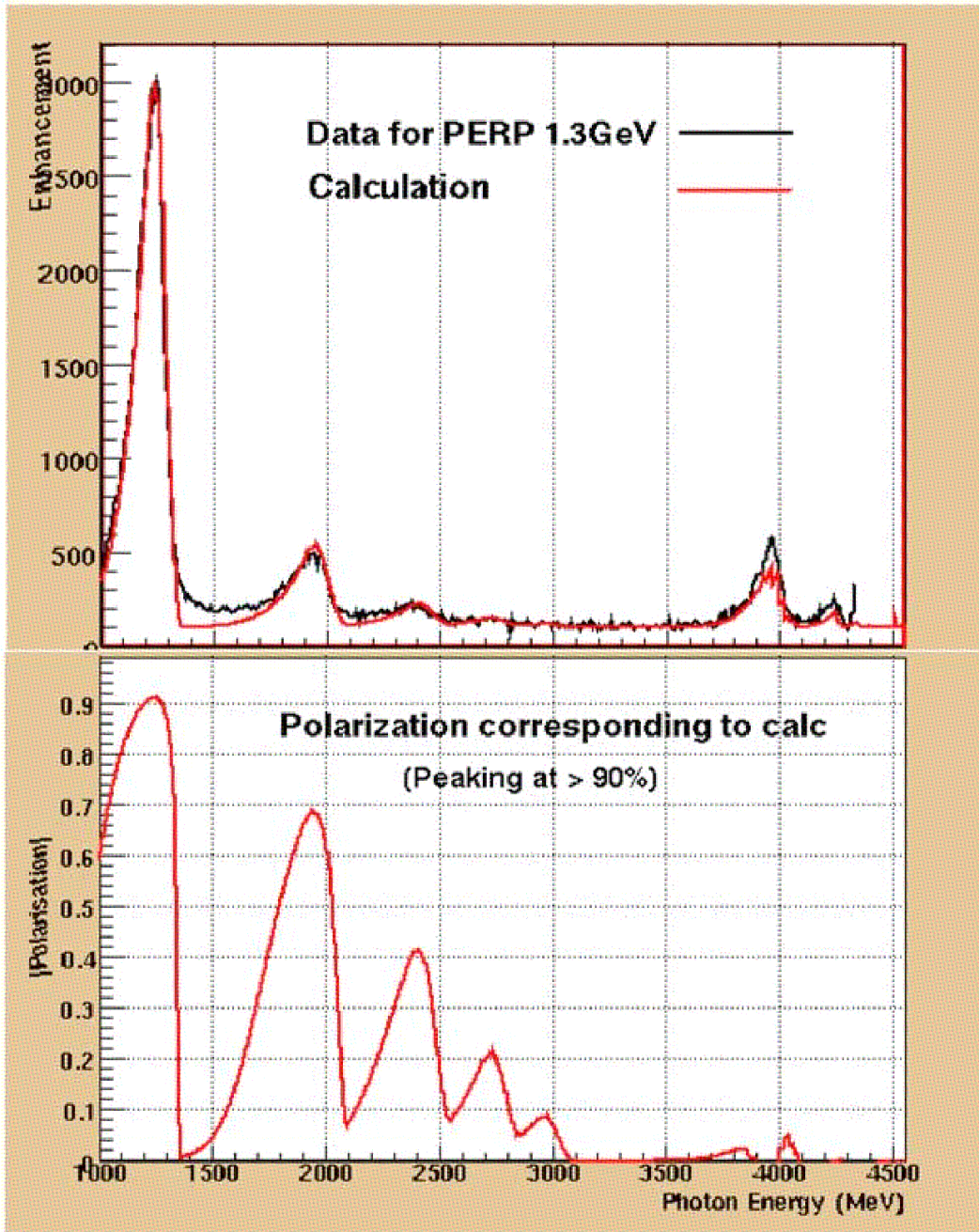


- Yellow : Torus Magnet
- Blue : Drift Chambers
- Purple : Cerenkov Counters

- Red : Time of Flight Scintillators
- Green : Electromagnetic Calorimeters

Polarized and Tagged Photon Beam

- Hall B has the ability of producing a beam of tagged photons.
- Using coherent bremsstrahlung radiation, these tagged photons can also be polarized.
- linearly polarized photon beam = unpolarized electron beam + oriented diamond radiator
- Can obtain 90% polarization



Final State Equation

- To isolate resonance contributions we aim to measure polarization observables.
- The final state equation for two mesons in the final state has a total of 15 observables.

$$I = I_0 \{ (1 + \Lambda_i \cdot \mathbf{P}) + \delta_{\odot} (\mathbf{I}^{\odot} + \Lambda_i \cdot \mathbf{P}^{\odot}) + \delta_l [\sin 2\beta (\mathbf{I}^s + \Lambda_i \cdot \mathbf{P}^s) + \cos 2\beta (\mathbf{I}^c + \Lambda_i \cdot \mathbf{P}^c)] \}$$

I_0 = unpolarized reaction rate
 Λ_i = degree of polarization of target
 $\delta_{\odot, l}$ = degree of polarization of photon beam
 \mathbf{P} = observables arising from target polarization
 $\mathbf{I}^{\odot, s, c}$ = observables arising from use of polarized photons
 β = orientation of polarization w.r.t. a final state particle

- Through the use of experimental conditions/setup, we can reduce the number of observables making a measurement possible.



Experimental Setup : g8b

- The g8b experiment ran from July - Sep 1st, 2005.
- Produced polarized photons at 5 coherent edge energies: 1.3, 1.5, 1.7, 1.9, 2.1 GeV
- Used linearly polarized photons incident on an unpolarized LH₂ target.

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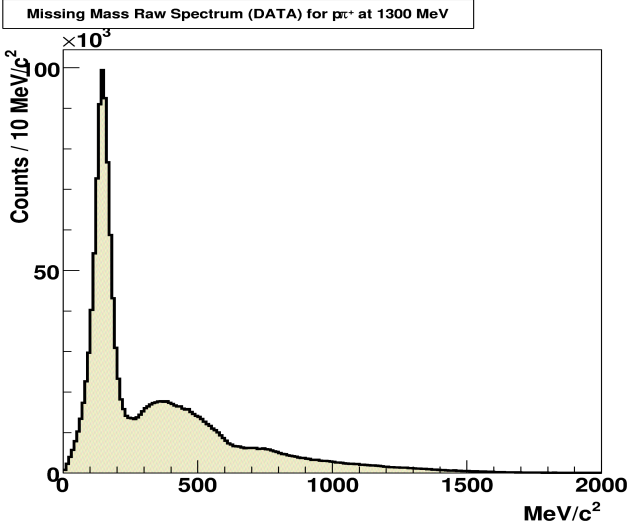


$$I = I_0 \{ 1 + \delta_l [\mathbf{I}^s \sin 2\beta + \mathbf{I}^c \cos 2\beta] \}$$

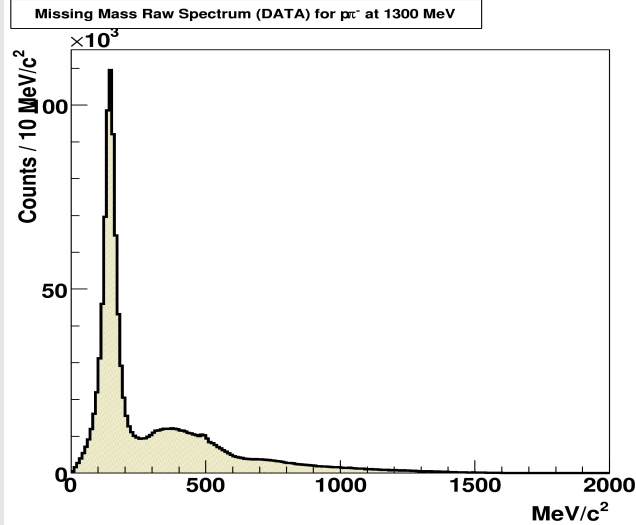
- I^c (also known as Σ in the single meson equation)
- I^s
- Measuring both for $\gamma p \rightarrow p \pi^+ \pi^-$

The g8b Data Set

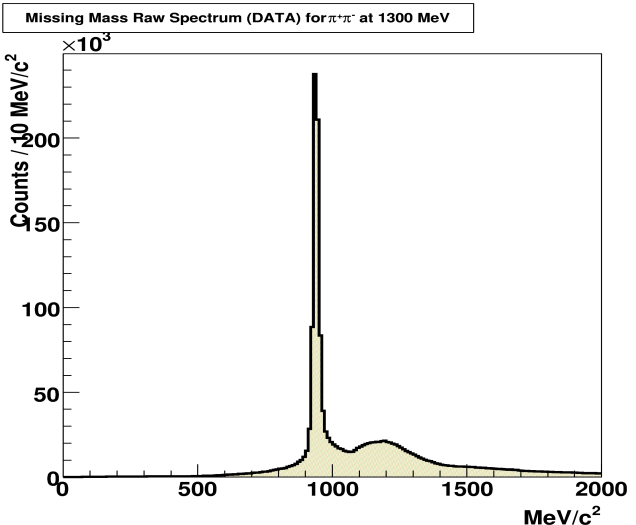
- $\gamma p \rightarrow p \pi^+ \pi^-$ from the g8b data set
- Kinematically fitting four topologies:



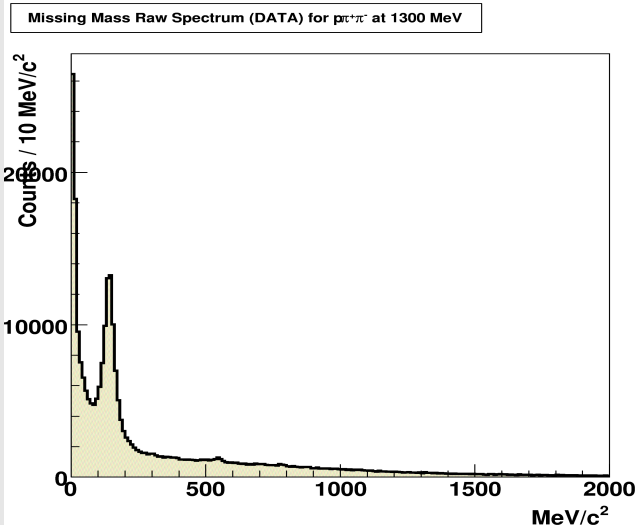
$p \pi^+ (\pi^-)$



$p \pi^- (\pi^+)$



$\pi^+ \pi^- (p)$



$p \pi^+ \pi^- ()$

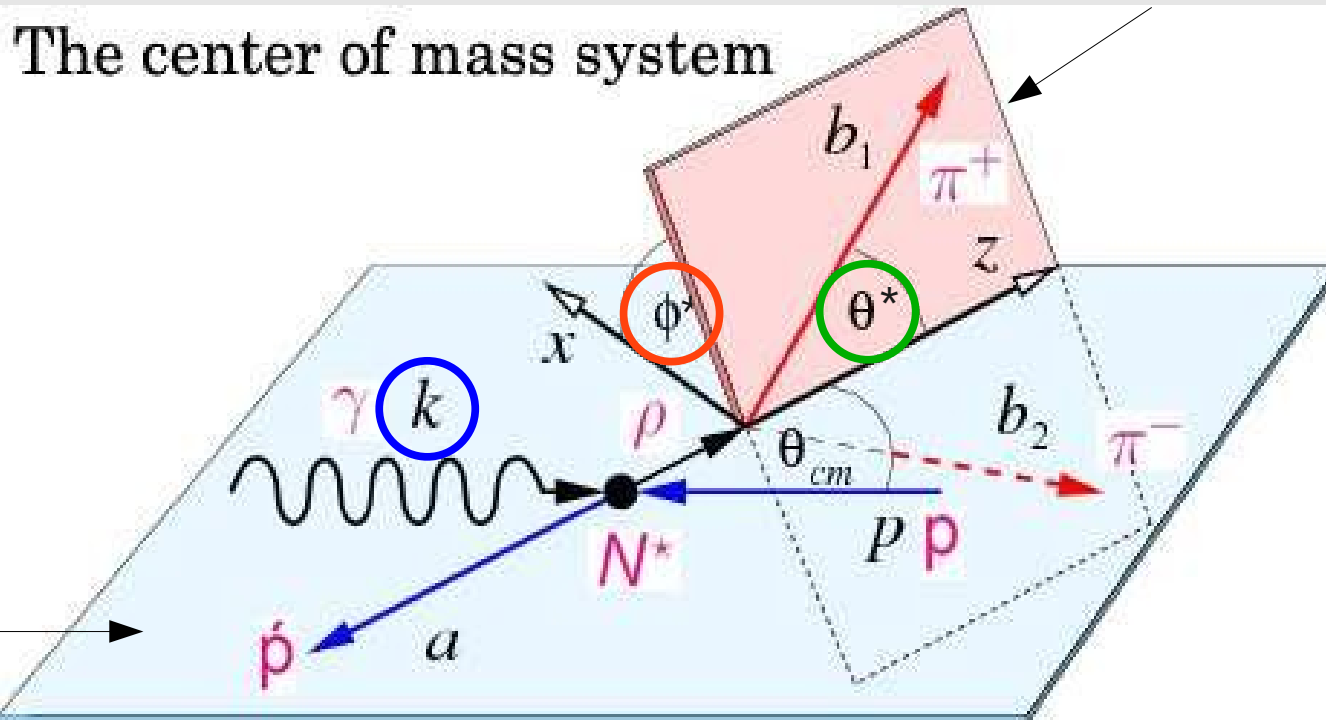
Correcting the g8b data set

- Momentum corrections were determined using the pulls generated from fitting to a $p \pi^+ \pi^-()$ final state.
- The errors applied to the covariance matrix (used for kin fitting) were determined by starting with the g11a errors. Errors were altered to obtain pull widths of one.
- A study into the sag of the tagger E-counters produced an energy dependent photon correction.
- Momentum corrections were refined after the application of the photon corrections.
- Covariance matrix errors were also refined/finalized.

CM Coordinate System for $\gamma p \rightarrow p \pi^+ \pi^-$

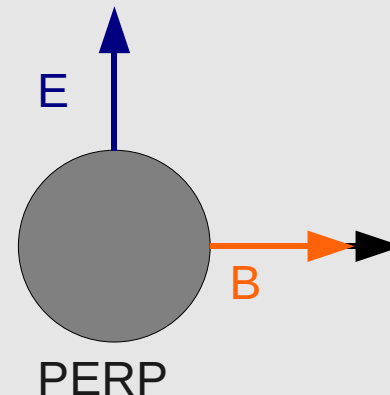
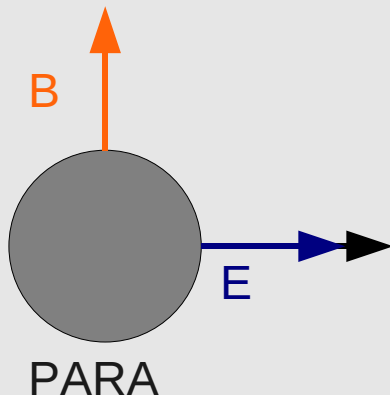
- Analysis of the $\gamma p \rightarrow p \pi^+ \pi^-$ channel requires the employ of 5 independent variables, for example : $\cos\theta_{cm}$, m , k , ϕ^* , θ^*
- $\gamma p \rightarrow N^* \rightarrow N \rho \rightarrow p \pi^+ \pi^-$

Event plane formed by 2 final state particles



Preliminary Results: Phi Distributions

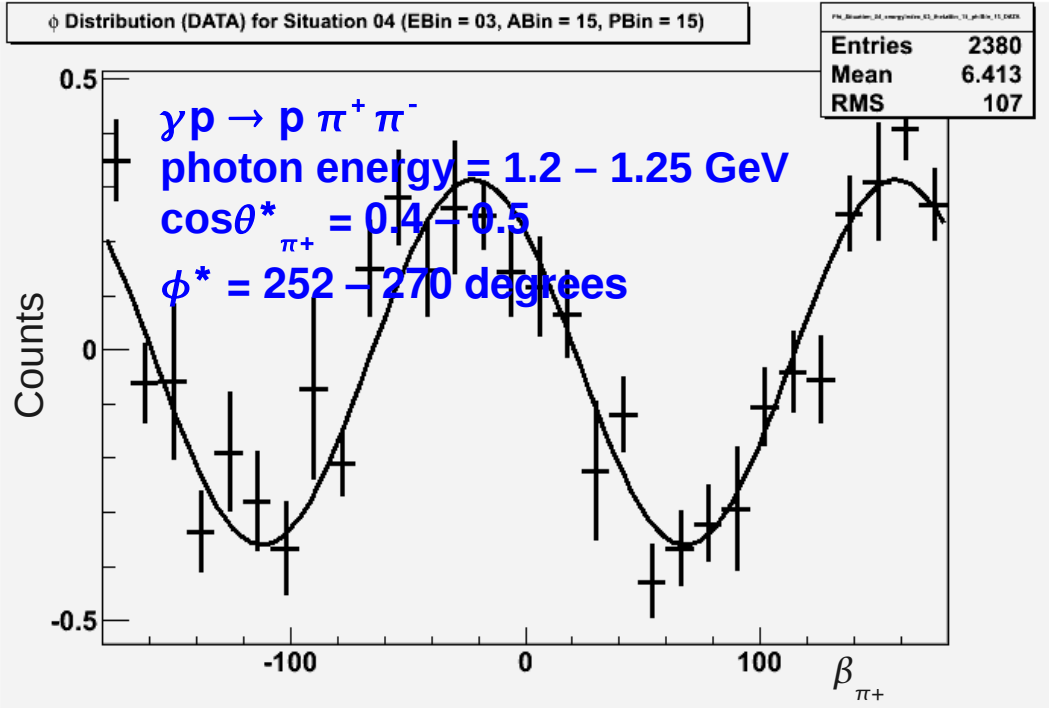
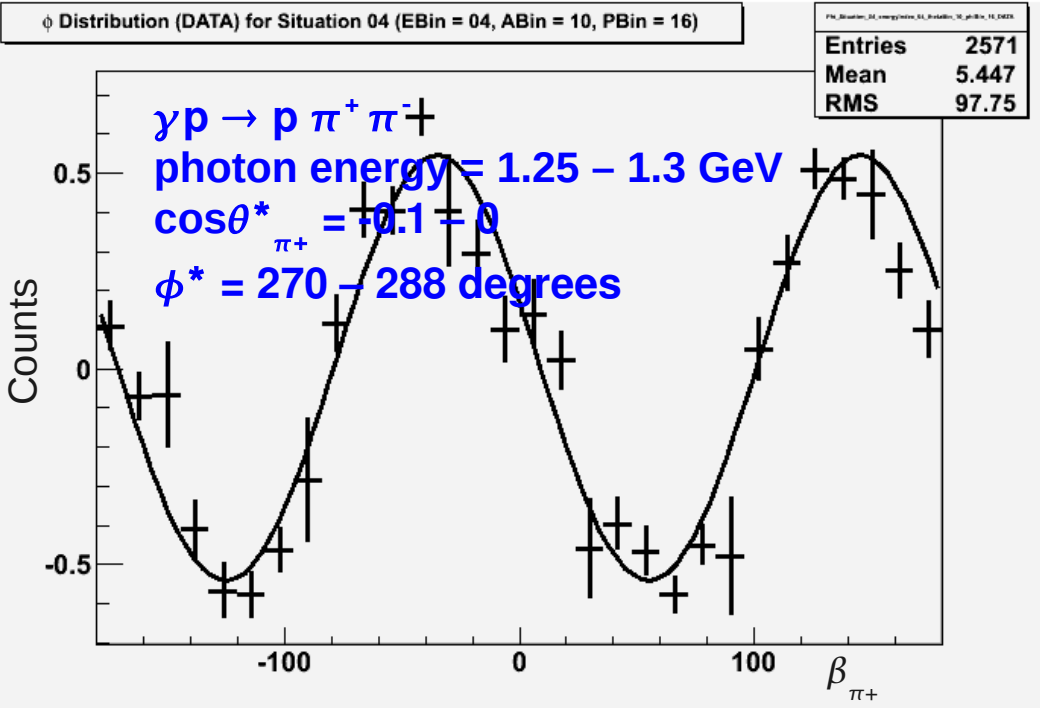
- Distribution of $p \pi^+ \pi^-$ events is normally independent of the lab angle ϕ but the polarized photons break that symmetry.
- CLAS Language: Linear Polarization
 - PARA = E field parallel to the floor
 - PERP = E field perpendicular to the floor
 - AMO = no polarization



Preliminary Results: Phi Distributions

- To remove systematics while maintaining statistics, the phi distributions for PARA are normalized to the PERP phi distributions.

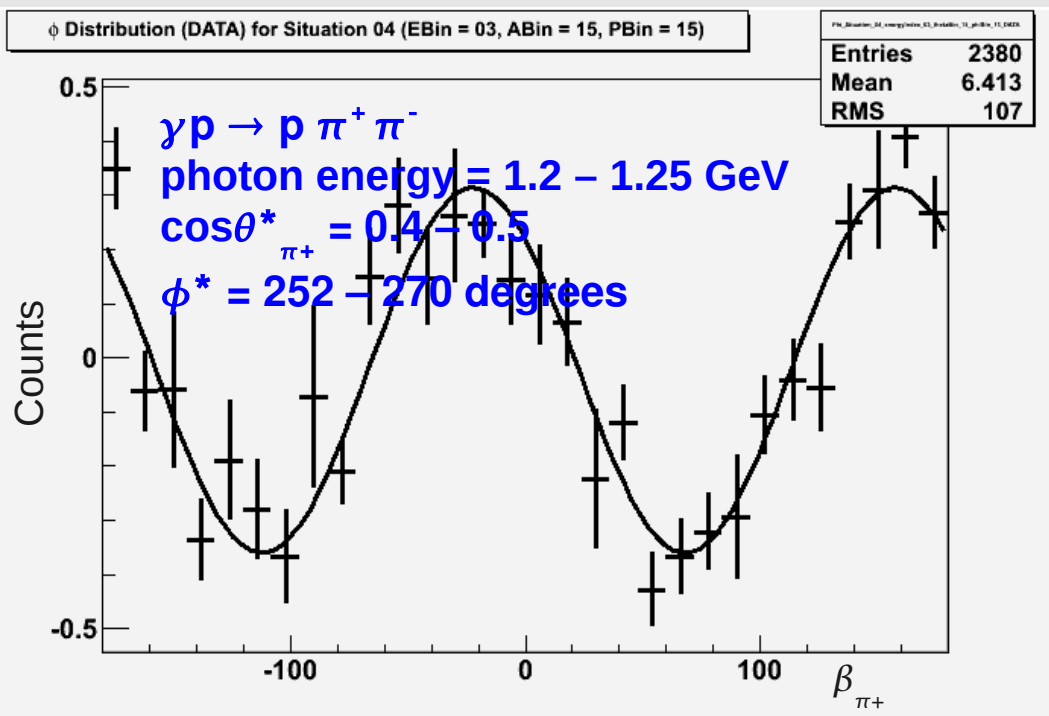
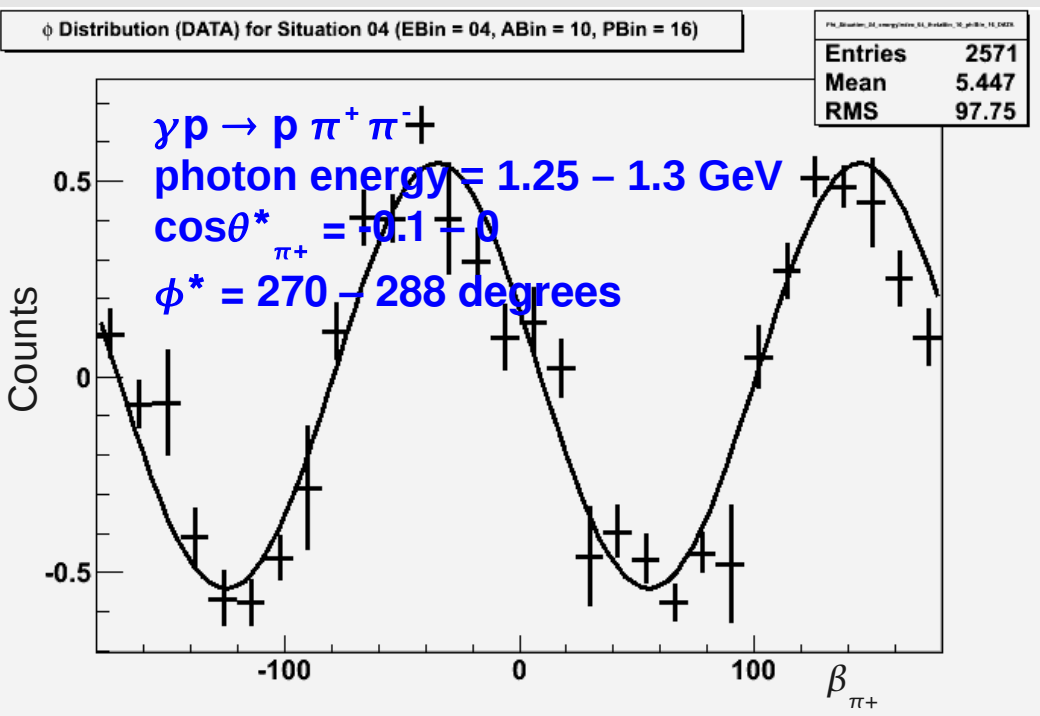
- Fit to :
$$\frac{PARA - PERP}{PARA + PERP} = \frac{I_{PARA} - I_{PERP}}{I_{PARA} + I_{PERP}}$$



Preliminary Results: Phi Distributions

- To remove systematics while maintaining statistics, the phi distributions for PARA are normalized to the PERP phi distributions.

- Fit to :
$$\frac{[(\delta_l^{para} + \delta_l^{perp}) I^s \sin(2\beta) + (\delta_l^{para} - \delta_l^{perp}) I^c \cos(2\beta)]}{2 + (\delta_l^{para} - \delta_l^{perp}) I^s \sin(2\beta) + (\delta_l^{para} + \delta_l^{perp}) I^c \cos(2\beta)} \quad \beta = \phi_{lab} + \phi_{polarization}$$

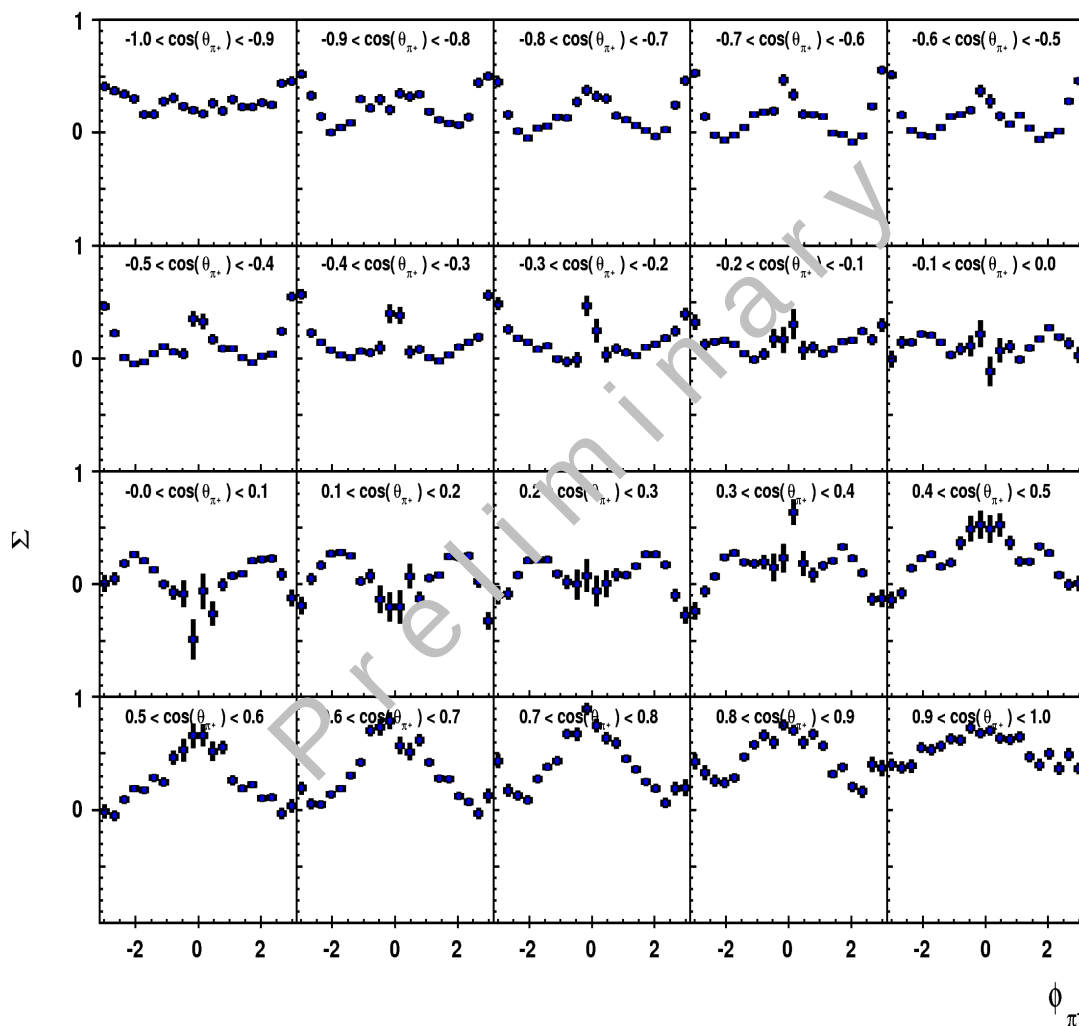


Preliminary Results : I^c

$$\frac{[(\delta_l^{para} + \delta_l^{perp}) I^s \sin(2\beta) + (\delta_l^{para} + \delta_l^{perp}) I^c \cos(2\beta)]}{2 + (\delta_l^{para} - \delta_l^{perp}) I^s \sin(2\beta) + (\delta_l^{para} - \delta_l^{perp}) I^c \cos(2\beta)}$$

$\gamma p \rightarrow p \pi^+ \pi^- (\Sigma)$

- Photon energy of 1100 – 1150 MeV
- Each square is an bin in $\cos\theta^*$ of the π^+ .
- Y-axis is the value of the observable.
- X-axis is the ϕ of the π^+ .
- We see a non-zero value for I^c (Σ).
- I^c is symmetric around the origin.

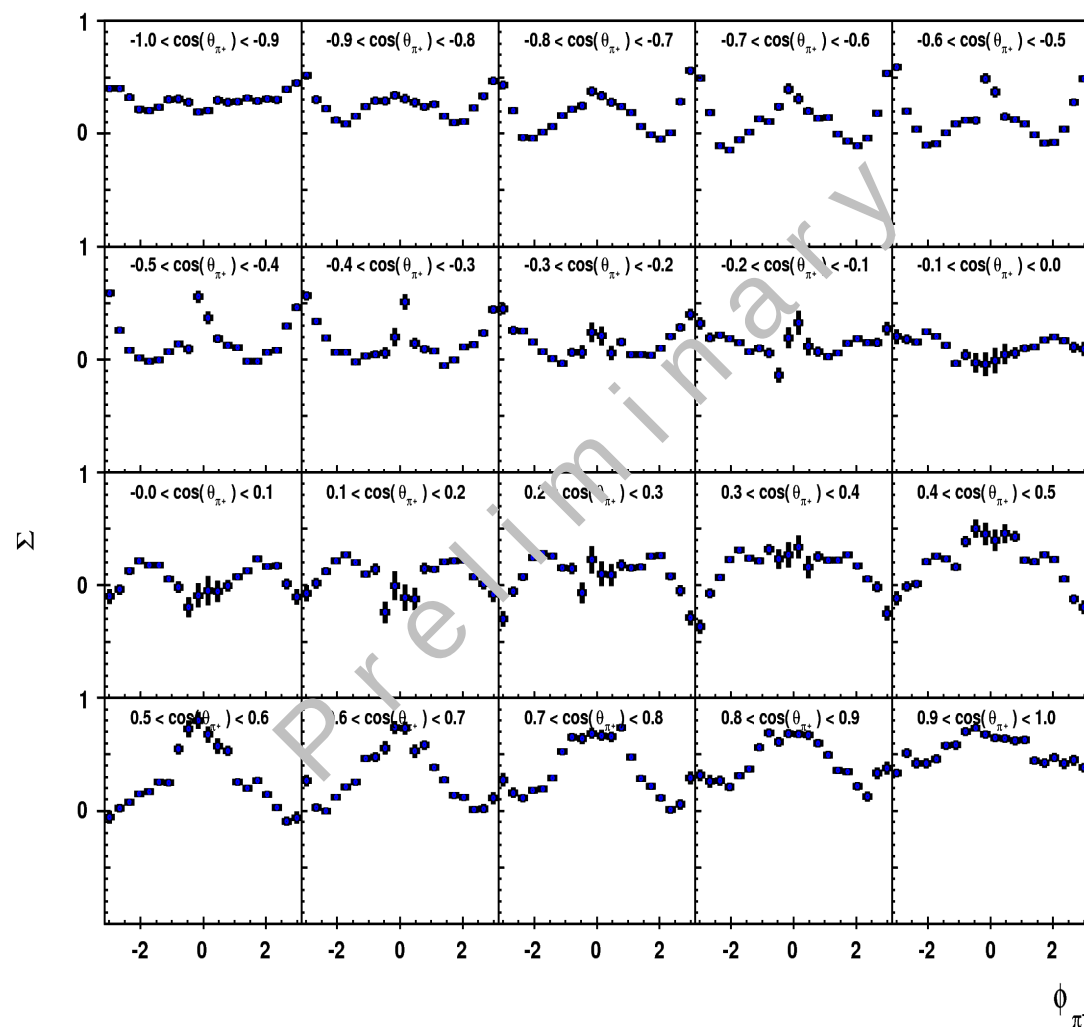


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$\gamma p \rightarrow p \pi^+ \pi^- ()$

- Photon energy of 1150 – 1200 MeV
- Each square is an bin in $\cos\theta^*$ of the π^+ .
- Y-axis is the value of the observable.
- X-axis is the ϕ of the π^+ .
- We see a non-zero value for I^c (Σ).
- I^c is symmetric around the origin.

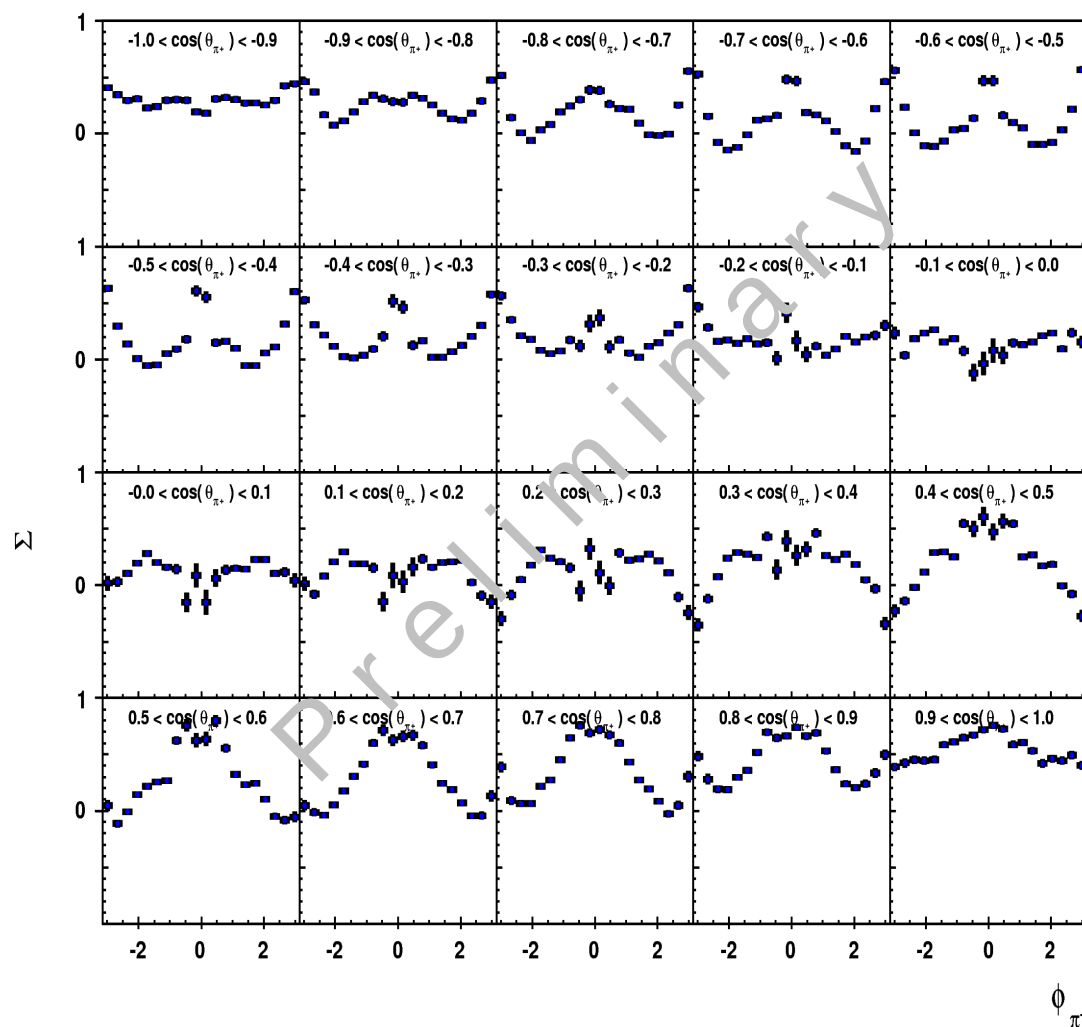


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- We see a non-zero value for I^c (Σ).
- I^c is symmetric around the origin.

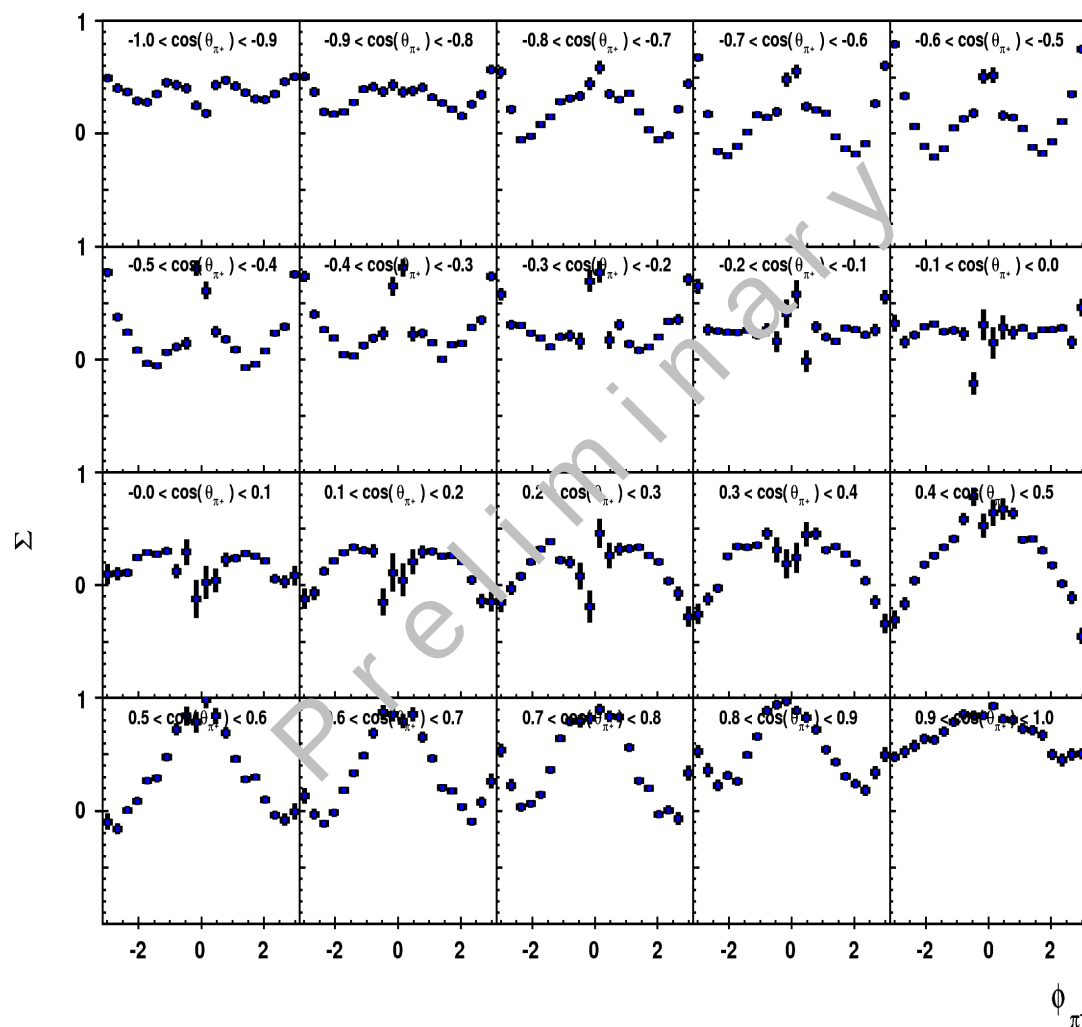


Preliminary Results : I^c

$$\frac{[(\delta_l^{para} + \delta_l^{perp}) I^s \sin(2\beta) + (\delta_l^{para} + \delta_l^{perp}) I^c \cos(2\beta)]}{2 + (\delta_l^{para} - \delta_l^{perp}) I^s \sin(2\beta) + (\delta_l^{para} - \delta_l^{perp}) I^c \cos(2\beta)}$$

$\gamma p \rightarrow p \pi^+ \pi^- ()$

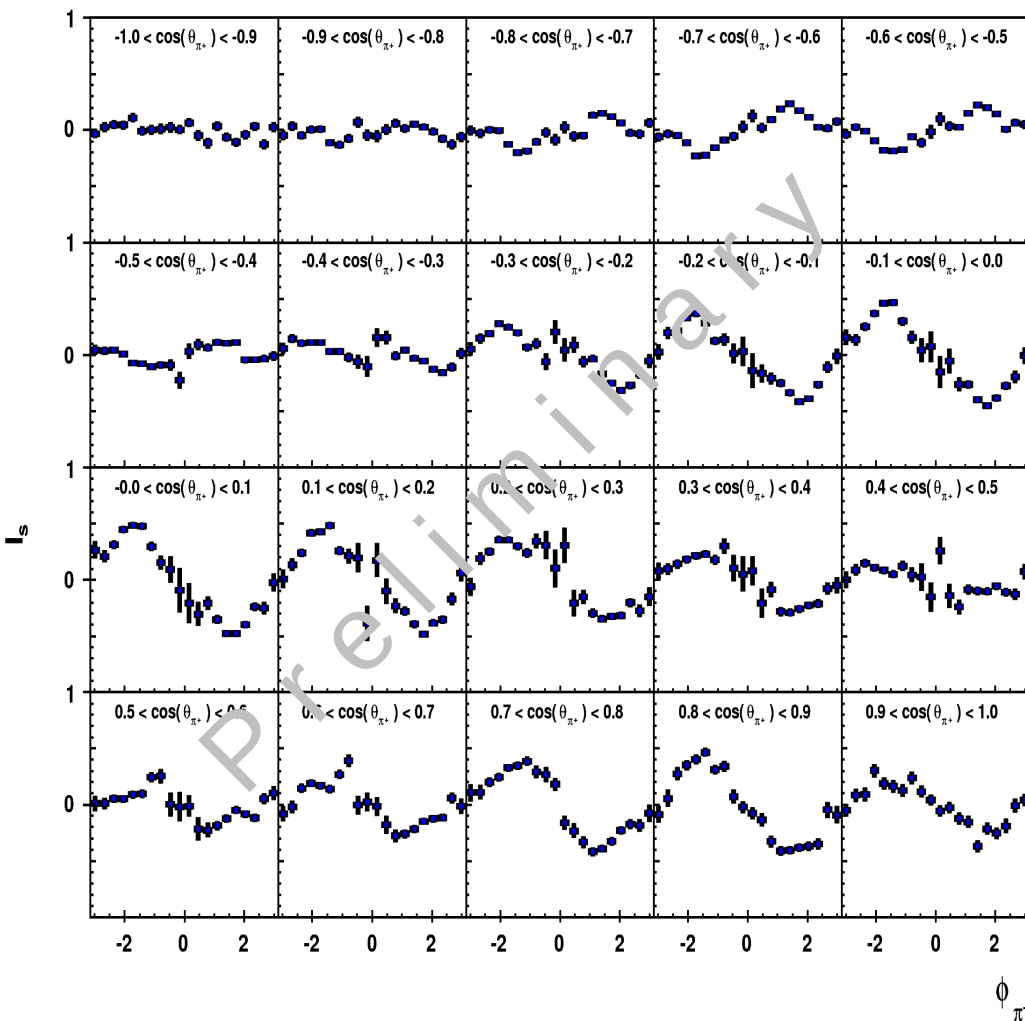
- Photon energy of 1250 – 1300 MeV
- Each square is an bin in $\cos\theta^*$ of the π^+ .
- Y-axis is the value of the observable.
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- We see a non-zero value for I^c (Σ).
- I^c is symmetric around the origin.



Preliminary Results : I^S

$$\frac{[(\delta_l^{para} + \delta_l^{perp}) I^s \sin(2\beta) + (\delta_l^{para} - \delta_l^{perp}) I^c \cos(2\beta)]}{2 + (\delta_l^{para} - \delta_l^{perp}) I^s \sin(2\beta) + (\delta_l^{para} + \delta_l^{perp}) I^c \cos(2\beta)}$$

$\gamma p \rightarrow p \pi^+ \pi^- ()$

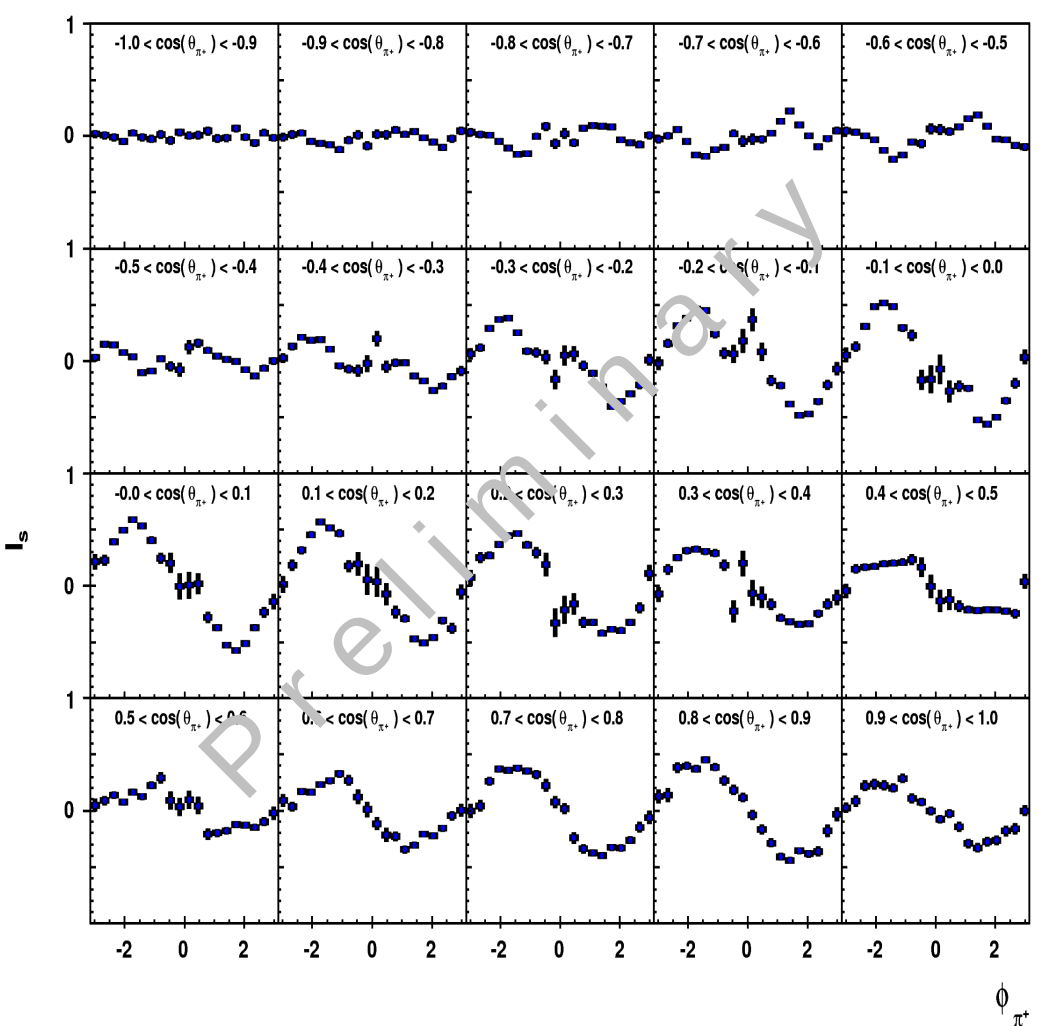


- Photon energy of 1100 – 1150 MeV
- Each square is an bin in $\cos\theta^*$ of the π^+ .
- Y-axis is the value of the observable.
- X-axis is the ϕ of the π^+ .
- I^S is antisymmetric around the origin.

Preliminary Results : I^S

$$\frac{[(\delta_l^{para} + \delta_l^{perp}) I^s \sin(2\beta) + (\delta_l^{para} - \delta_l^{perp}) I^c \cos(2\beta)]}{2 + (\delta_l^{para} - \delta_l^{perp}) I^s \sin(2\beta) + (\delta_l^{para} + \delta_l^{perp}) I^c \cos(2\beta)}$$

$\gamma p \rightarrow p \pi^+ \pi^- ()$ ↓

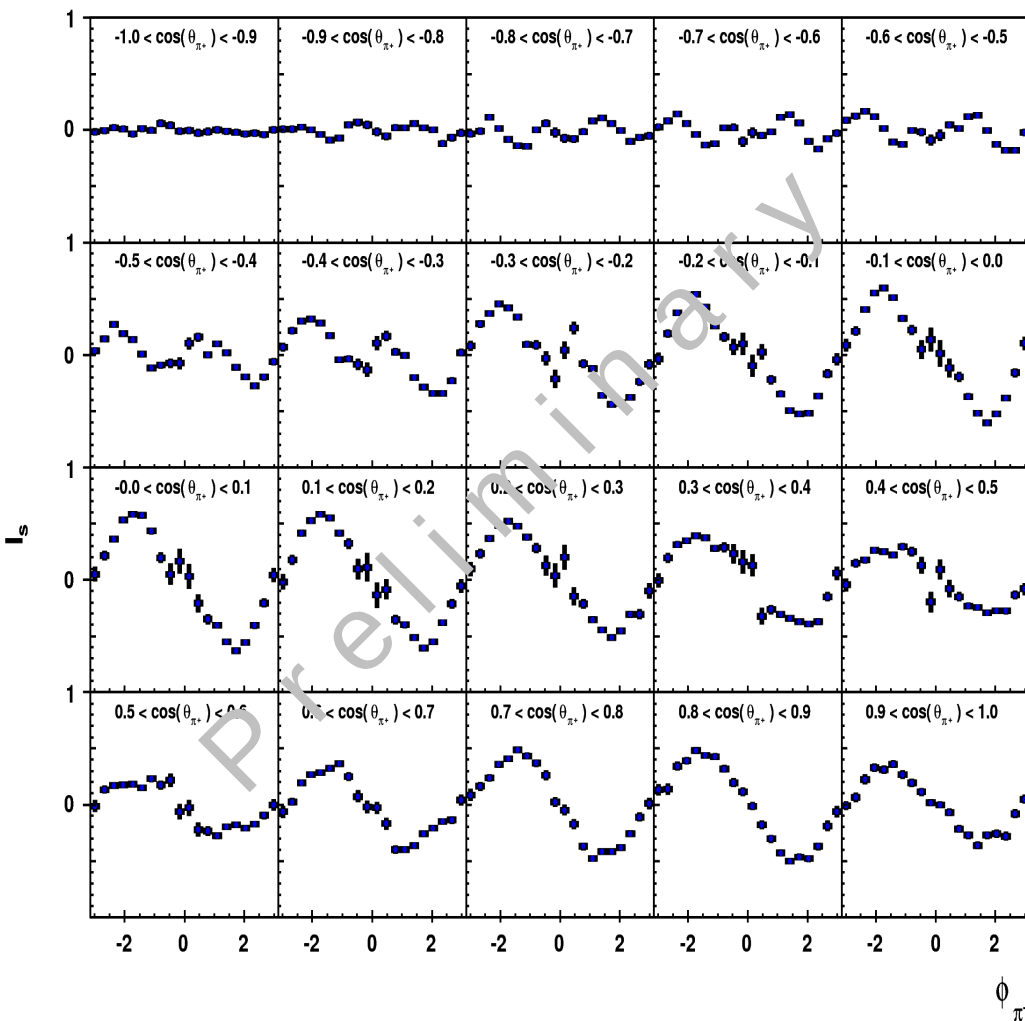


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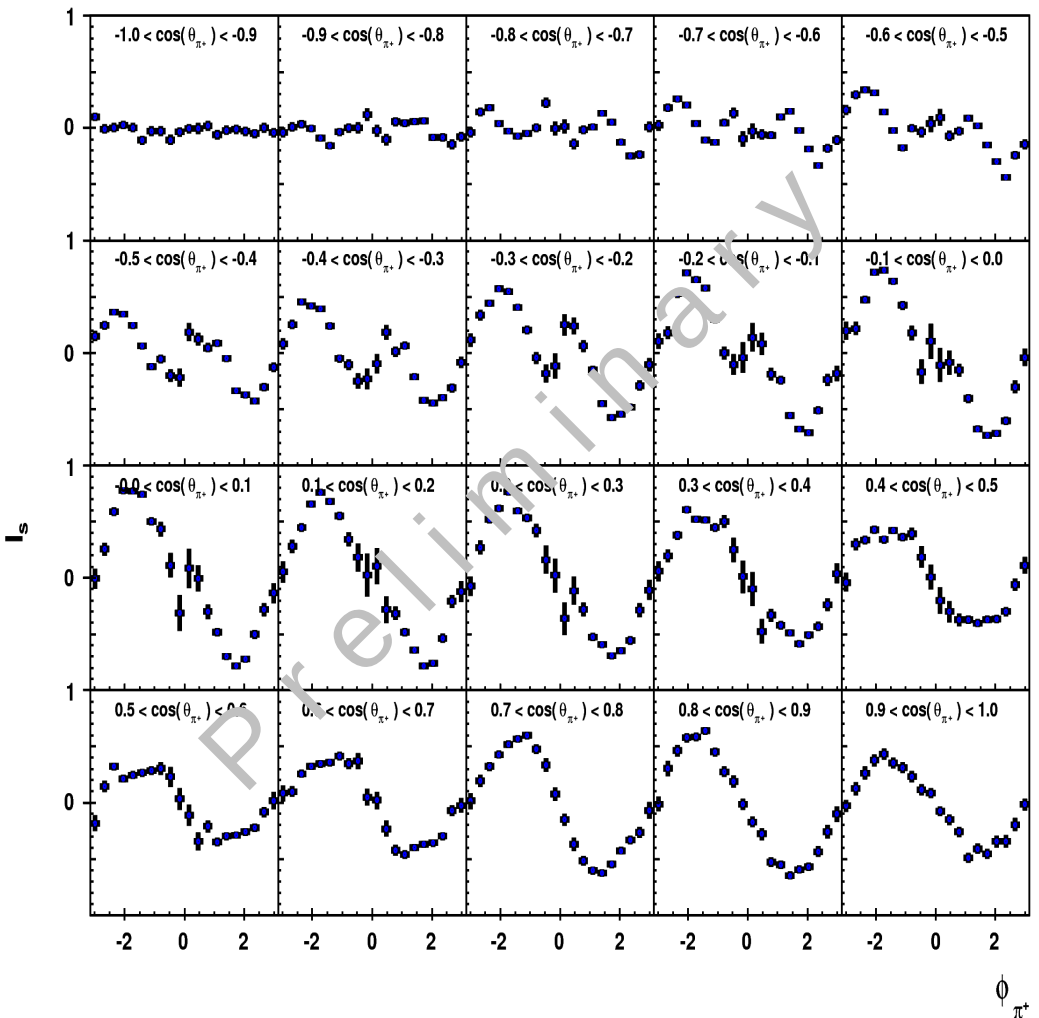


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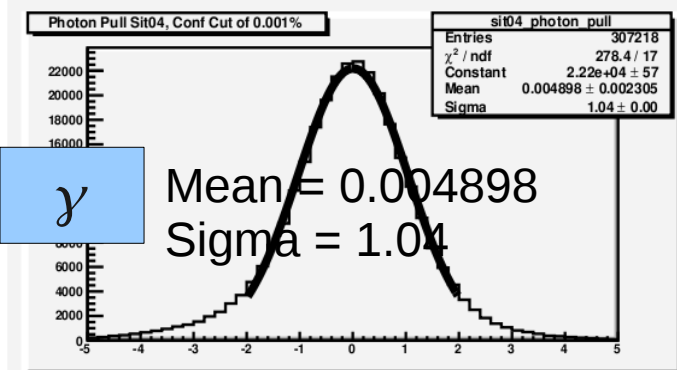
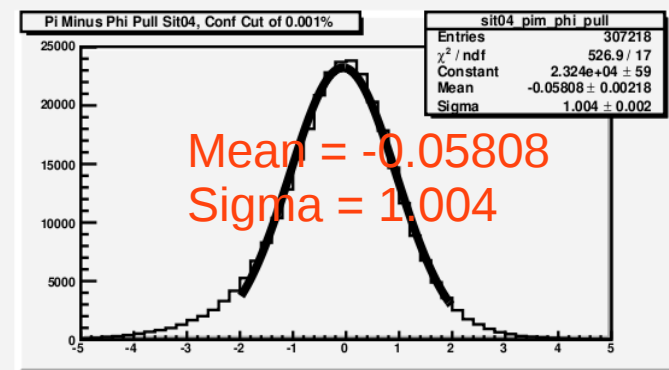
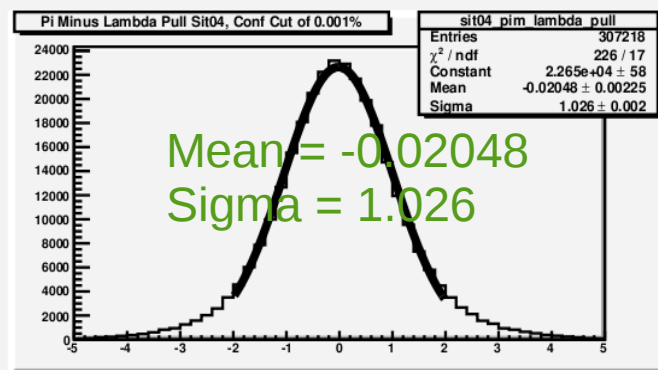
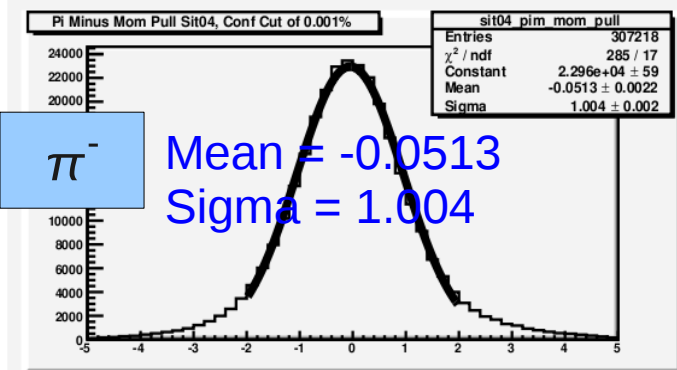
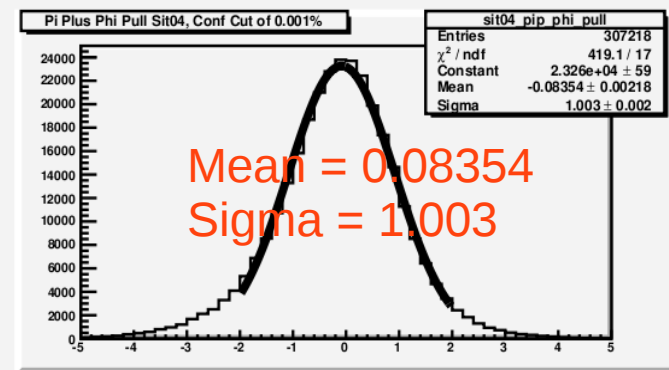
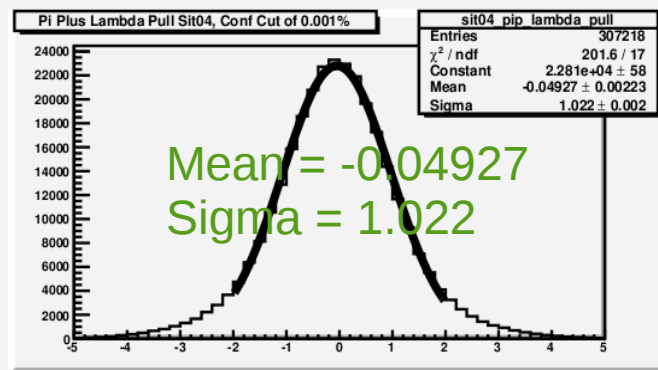
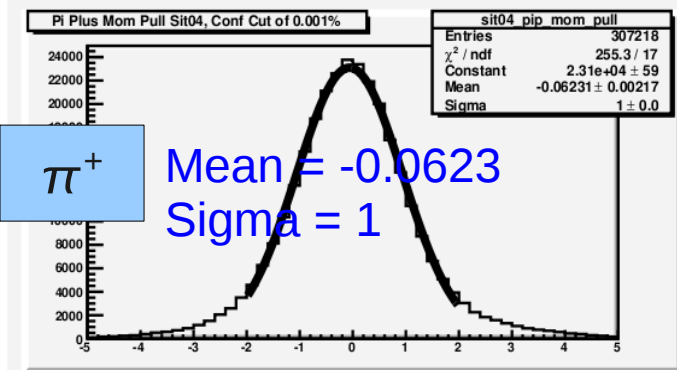
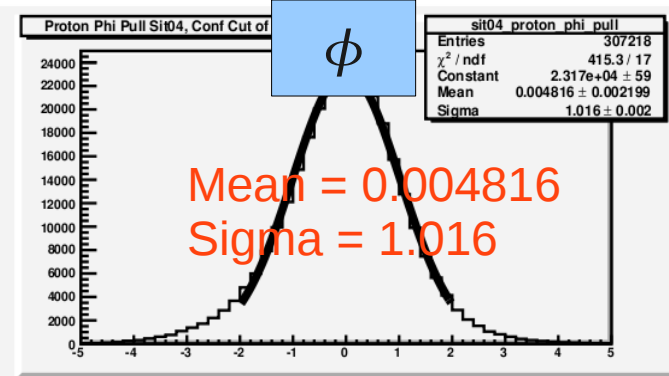
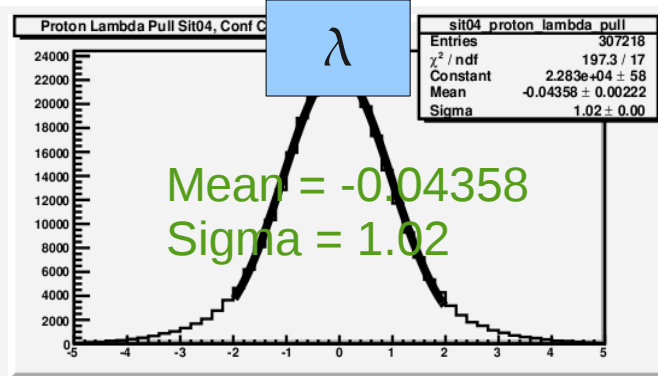
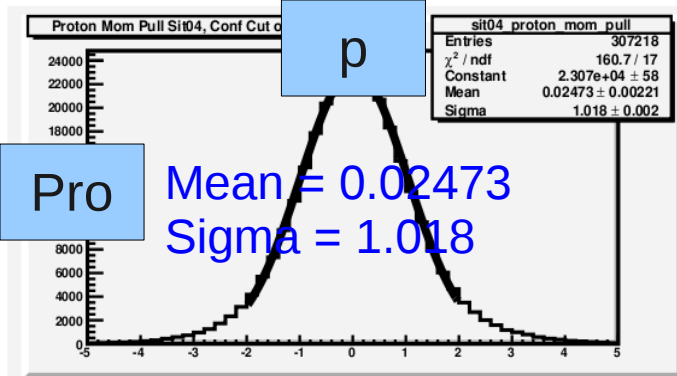
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Summary

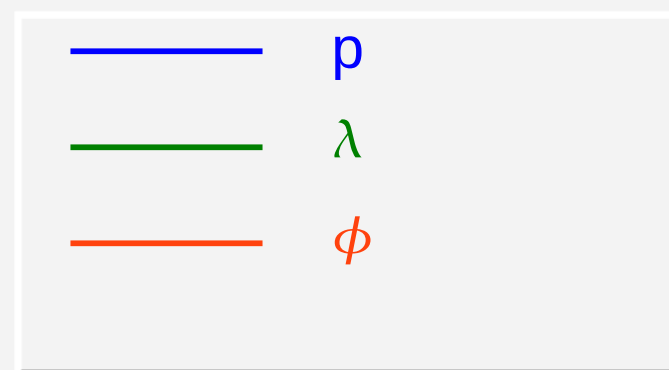
- The first (preliminary) measurements of the polarization observables I^s and I^c for $\gamma p \rightarrow p \pi^+ \pi^-$ have been made using the high amount of statistics available in the g8b data set (possible binning in third angle).
- The measurements of polarization observables are key to understanding the missing resonances in the CQM.
- Other polarized photoproduction experiments (such as FROST) will allow access to more polarization observable measurements.
- To Do: Background subtraction (via Q-value method)₂₈



END

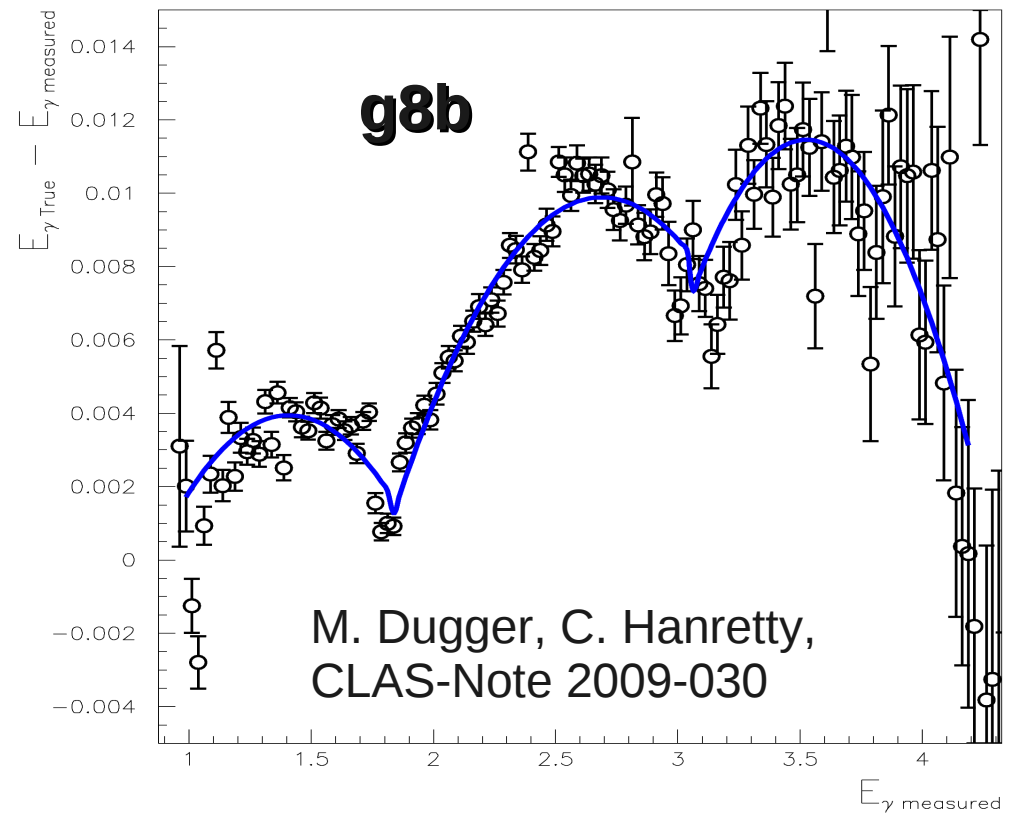
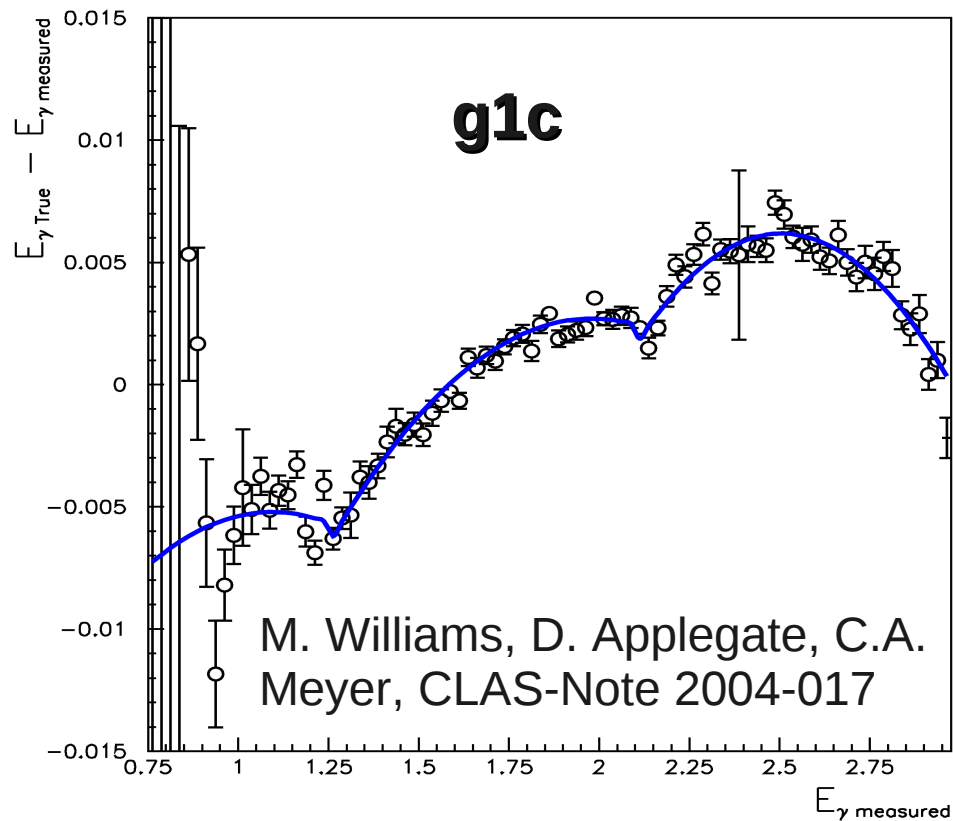


$\gamma p \rightarrow p \pi^+ \pi^- ()$
(Run #048326)



Tagger Sag

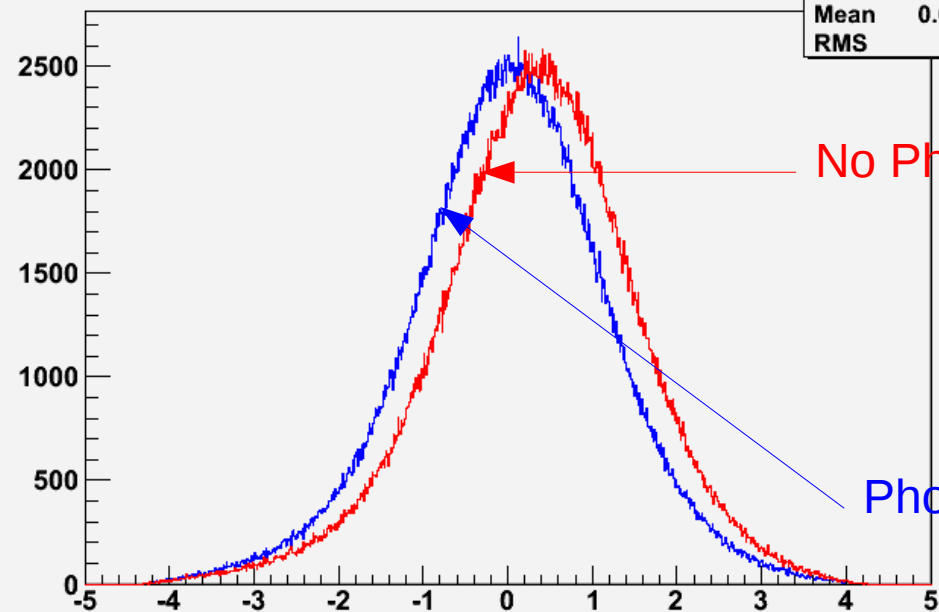
- Results from a physical “sagging” of the E-counter scintillator bars.
- Has been seen in several runs.



Tagger Sag

- Results from a physical “sagging” of the E-counter scintillator bars.
- Has been seen in several runs.
- Requires and energy-dependent photon energy correction.

Photon Pull Values v. Photon Energy, ConfCut > 0.1 Percent



Fitted value of par[1]=Mean

