Baryon Spectroscopy using the CLAS Spectrometer and the Frozen Spin Target (FROST) at Jefferson Laboratory



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Nuclear Physics Seminar April 09, 2010

Outline

Introduction

- Baryon Spectroscopy
- Polarization Observable
- PROST Experiment
 - The CLAS at JLab
 - The FRozen-Spin Target (FROST)
 - The FROST-g9a run Period
- 3 Event Selection
 - The particle identification
 - The dilution factor
 - The beam and target polarization

4 The Preliminary Results

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Introduction

FROST Experiment Event Selection The Preliminary Results Baryon Spectroscopy Polarization Observable

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Baryon Spectroscopy Polarization Observable

What are hadrons?

Hadrons are composed of quarks bound by the strong interaction.



Quantum Chromodynamics (QCD)

- The theory of how quarks and gluons interact with themselves and each other
- →The study of the properties of baryon resonances
 →The N* Program

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Baryon Spectroscopy Polarization Observable

One of the Goals of the N* Program ...

Search for missing or yet unobserved resonances

Quark models predict many more baryons than have been observed

	****	***	**	*
N Spectrum	11	3	6	2
Δ Spectrum	7	3	6	6

Possible solutions:

1. Quark-diquark structure



one of the internal degrees of freedom is frozen

- →according to PDG
 (Phys. Lett. B 667, 1 (2008))
 →little known
 (many open questions left)
- 2. Have not been observed, yet

Nearly all existing data result from πN scattering experiments

 If the missing resonances did not couple to Nπ, they would not have been discovered!!

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Baryon Spectroscopy Polarization Observable

The excited states of the nucleon

Constituent quark models: Gluon-exchange model



Baryon Spectroscopy Polarization Observable

The excited states of the nucleon

Constituent quark models: N^* resonances (Isospin $\frac{1}{2}$)



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Baryon Spectroscopy Polarization Observable

The excited states of the nucleon

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Baryon Spectroscopy Polarization Observable

The excited states of the nucleon

Constituent quark models: N^* resonances (Isospin $\frac{1}{2}$)



Baryon Spectroscopy Polarization Observable

The Motivation for the $\pi^+\pi^-$ photoproduction



• The cross section of the $\pi^+\pi^-$ photoproduction dominates above W \approx 1.8GeV

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Baryon Spectroscopy Polarization Observable

The Motivation for the $\pi^+\pi^-$ photoproduction



• The cross section of the $\pi^+\pi^-$ photoproduction dominates above W \approx 1.8GeV

The excited states are found as broadly overlapping resonances

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Baryon Spectroscopy Polarization Observable

The Motivation for the $\pi^+\pi^-$ photoproduction



- The cross section of the $\pi^+\pi^-$ photoproduction dominates above W \approx 1.8GeV
- The excited states are found as broadly overlapping resonances
 - \rightarrow The polarization observables can isolate single resonances from other interference terms

Baryon Spectroscopy Polarization Observable

The differential cross section for $\gamma p \rightarrow p \pi^+ \pi^-$

The differential cross section for $\gamma p \rightarrow p \pi^+ \pi^-$

(without measuring the polarization of the recoiling nucleon)

 $\frac{\mathrm{d}\sigma}{\mathrm{d}x_{i}} = \sigma_{0}\left\{\left(\mathbf{1} + \vec{\Lambda}_{i} \cdot \vec{\mathbf{P}}\right) + \delta_{\odot}\left(\mathbf{I}^{\odot} + \vec{\Lambda}_{i} \cdot \vec{\mathbf{P}}^{\odot}\right)\right\}$

+ δ_{l} [sin 2 β (l^s + $\vec{\Lambda}_{i} \cdot \vec{P}^{s}$) + cos 2 β (l^c + $\vec{\Lambda}_{i} \cdot \vec{P}^{c}$)]}

- σ_0 : The unpolarized cross section
- β : The angle between the direction of polarization and the x-axis
- $\delta_{\odot,I}$: The degree of polarizaton of the photon beam $\Rightarrow \delta_{\odot}$, and δ_{I}
- $\vec{\Lambda}_i$: The polarization of the initial nucleon $\Rightarrow (\Lambda_x, \Lambda_y, \Lambda_z)$
- $I^{\odot, s, c}$: The observable arising from use of polarized photons $\Rightarrow I^{\odot}, I^{s}, I^{c}$
- \vec{P} : The polarization observable \Rightarrow (P_x , P_y , P_z) (P_x^{\odot} , P_y^{\odot} , P_z^{\odot}) (P_x^s , P_y^s , P_z^s) (P_x^c , P_y^c , P_z^c) 15 Observables

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Baryon Spectroscopy Polarization Observable

Polarization Observable

The differential cross section for $\gamma p \rightarrow p \pi^+ \pi^-$

(without measuring the polarization of the recoiling nucleon)

 $\frac{\mathrm{d}\sigma}{\mathrm{d}x_{i}} = \sigma_{0}\left\{\left(\mathbf{1} + \vec{\Lambda}_{i} \cdot \vec{\mathbf{P}}\right) + \delta_{\odot}\left(\mathbf{I}^{\odot} + \vec{\Lambda}_{i} \cdot \vec{\mathbf{P}}^{\odot}\right) \quad 15 \text{ Observables}\right\}$

+ δ_{l} [sin 2 β (l^s + $\vec{\Lambda}_{i} \cdot \vec{P}^{s}$) + cos 2 β (l^c + $\vec{\Lambda}_{i} \cdot \vec{P}^{c}$)]}

The circularly-polarized beam $\rightarrow \delta_I = 0$

The longitudinally-polarized target $\rightarrow \Lambda_{\textbf{x}} = \Lambda_{\textbf{y}} = 0$

$$\frac{\mathrm{d}\sigma}{\mathrm{d}x_{\mathrm{i}}} = \sigma_0 \left\{ \left(1 + \Lambda_z \cdot \mathbf{P}_z \right) + \delta_{\odot} \left(\mathbf{I}^{\odot} + \Lambda_z \left(\mathbf{P}_z^{\odot} \right) \right\} \right\}$$

3 Observables

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The CLAS at JLab The FRozen-Spin Target (FROST) The FROST-g9a run Period

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Jefferson Laboratory in Newport News, VA





The continuous electron beam accelerator facility (CEBAF) can deliver a continuous electron beam up to 6 GeV. a continuous electron beam up to 6 GeV.

The CLAS at JLab The FRozen-Spin Target (FROST) The FROST-g9a run Period

CEBAF Large Acceptance Spectrometer (CLAS)



The CLAS at JLab The FRozen-Spin Target (FROST) The FROST-g9a run Period

CEBAF Large Acceptance Spectrometer (CLAS)



The CLAS at JLab The FRozen-Spin Target (FROST) The FROST-g9a run Period

The tagging system at CLAS

JLAB Hall B bremsstrahlung photon tagger



The CLAS at JLab The FRozen-Spin Target (FROST) The FROST-g9a run Period

The tagging system at CLAS

JLAB Hall B bremsstrahlung photon tagger



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The CLAS at JLab The FRozen-Spin Target (FROST) The FROST-g9a run Period

The things we need for the FROST experiment

- The polarized photon beam The tagging system at CLAS
- The polarized proton The Frozen-Spin Target



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The Basics of Polarization



In absence of a magnetic field, a collection of spins is randomly oriented.



With the magnetic field,

the spins either parallel or anti-parallel to the field will be oriented

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Polarization = excess of one orientation over the other

- Oscillating EM fields, produced by atomic vibration, tends to randomize (de-polarize) the spins.
- Strength of vibrations decreases at low temperature.

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Polarization and Thermal Equilibrium

Any ensemble of atoms or nuclei with a magnetic moment can be polarized via the Zeeman interaction: $\vec{\mu} \cdot \vec{B}$

In general, the populations of the Zeeman levels (once equilibrium has been reached) will obey a Boltzmann distribution:

$$\frac{N(\uparrow)}{N(\downarrow)} = e^{\frac{-2\,\vec{\mu}\cdot\vec{B}}{kT}} \qquad P_{te} = \frac{N(\uparrow) - N(\downarrow)}{N(\uparrow) + N(\downarrow)} = \tanh\left(\frac{\vec{\mu}\cdot B}{kT}\right)$$
$$(T = \text{Temperature}, P_{te} = \text{Thermal Equilibrium Polarization})$$

The polarization will approach thermal equilibrium with a characterisite 1/e time constant t_1 :

 $P(t) = P_{te} (1 - e^{-t/t_1})$ "t₁: Spin-Lattice Relaxation Time"

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A Simple Way to Polarize

Brute Force Polarization $P_{\text{te}} = \tanh{\left(\frac{\vec{\mu} \cdot \vec{B}}{\iota \tau}\right)}$

To get high polarization

minimize T

Thermal Equilibrium Polarization at 5 Tesla



Disadvantages:



Requires very large magnet

- Low temperatures require low luminosity
- Polarization can take a very long time (protons slow, electrons fast)

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A Better Way – Dynamic Nuclear Polarization

- (1) Use brute force to polarize free electrons in the target material.
- (2) Use microwaves to "transfer" this polarization to nuclei.

Mutual electron-nucleus spin flips re-arranges the nuclear Zeeman populations to favor one spin state over the other.

For best results:

DNP is performed at B/T conditions where electrons t_1 is short (ms) and nuclear t_1 is long (minutes):

JLab:
$$B = 5 T$$

 $T = 1 K$

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Materials for DNP Targets

Choice of material dictated by:

Maximum polarization

Resistance to ionizing radiation

Presence of unpolarized nuclei

Presence of unwanted, polarized nuclei

Compromise: Butanol (C₄H₉OH)

• Quality (dilution) factor:

$$f = \vec{N}/N_{\text{total}} = 10/74 \approx 0.13$$

The holding magnet for FROST : 0.5 T



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The Frozen-Spin Target (FROST)

Operation is more complicated:

- (1) Polarize target material via DNP at 5 T and 0.5 K (Polarizing Mode)
- (2) After optimum polarization is obtained, turn off microwaves and 5 T magnet
- (3) Use a 2nd magnet (~0.5 T) and very low temperatures to "freeze" the polarization (Frozen Spin Mode)
- (4) Polarization will decay very slowly with a time constant of several days
- (5) After polarization decays to about 50 % of its initial value, go back to step 1



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The Frozen-Spin Target (FROST)



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The Frozen-Spin Target (FROST)





How to polarize the FROST?



The magnets in the FROST experiment

- (a) The longitudinal holding magnet. (About 0.5 T)
- (b) The transvere holding magnet. (g9b) (Charles Hanretty)
- (c) The polarizing magnet. (5 Tesla internal solenoid)

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Polarizing Mode

- Temperature 0.5 K
- 5T magnet ON (polarize the electrons)
- * Microwave ON (transfer this polarization to nuclei)

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* Photon beam OFF

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The Frozen-Spin Target (FROST)





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The Frozen-Spin Target (FROST)





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The Frozen-Spin Target (FROST)





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How to polarize the FROST?



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- (c) The polarizing magnet. (5 Tesla internal solenoid)

Frozen Spin Mode

- * 5T magnet OFF
- Microwave OFF
- * Temperature \sim 0.05 K
- 0.5T magnet ON (holding magnet)
- * Photon beam ON

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The Frozen-Spin Target (FROST) - polarizing mode



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The Frozen-Spin Target - Summary of Results

	Expectation	Result
Base temperature:	50 mK	28 mK (w/o beam) 30 mK (w/ beam)
Cooling Power:	10 μ W (Frozen) 20 mW (Polarizing)	800 μW @ 50mK 60mW @ 300 mK
Polarization:	80 %	+ 82 % - 85 %
1/e Relaxation Time:	500 hours	2700 hours (+ Pol.) 1600 hours (-Pol.)

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The CLAS at JLab The FRozen-Spin Target (FROST) The FROST-g9a run Period

The FROST-g9a run Data

The FROST run period: Nov. 3, 2007 - Feb. 12, 2008 Data set: 35 TBytes

Production Data

Target:

Longitudinal polarized target

Average target polarization

 \sim 82 % (+Pol) and 85 % (-Pol)

Photon beam:

- Circularly and linearly polarized photon beam 0.5 - 4.5 GeV
- Electron beam polarization $\sim 85\%$





10.5 Billion events

The particle identification The dilution factor The beam and target polarization

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The particle identification The dilution factor The beam and target polarization

The particle identification - The beta cut



The beta cut

 $\begin{array}{c} \mbox{Calculated beta - Measured beta} \\ (\frac{momentum}{energy}) & (TOF) \end{array}$



The particle identification The dilution factor The beam and target polarization

The 4 different topologies of $\gamma p \rightarrow p \pi^+ \pi^-$

- ♦ The topology : $\gamma p \rightarrow p \pi^+(\pi^-)$
- ♦ The topology : $\gamma p \rightarrow p \pi^-(\pi^+)$
- ♦ The topology : $\gamma p \rightarrow \pi^+ \pi^-(p)$
- \diamond The topology : $\gamma p \rightarrow p \pi^+ \pi^-$



The particle identification The dilution factor The beam and target polarization

selecting the target



The particle identification The dilution factor The beam and target polarization

What is the dilution factor?





- The hydrongen atoms are polarized longitudinally in FROST experiment
- ♦ The butanol (C_4H_9OH) target has the unpolarized atoms like the carbon (C) or the oxyzen (O).



The particle identification The dilution factor The beam and target polarization

What is the dilution factor?





 The dilution factor is defined as the ratio between the polarized hydrogen and the full butanol contribution to the cross section

The dilution factor =
$$\frac{\sigma_H}{\sigma_{C_4H90H}} = \frac{N_{butanol} - S \cdot N_{carbon}}{N_{butanol}}$$
 (where S : The scaling factor)

The particle identification The dilution factor The beam and target polarization

The scaling factor



The scaling factor normalizes the distribution of the two targets.

◊ comparing [0.6,0.8] of the two targets; the butanol and carbon.

The scaling factor = $\frac{The \ blue \ part}{The \ pink \ part}$

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The particle identification The dilution factor The beam and target polarization

The scaling factor



- The scaling factor normalizes the distribution of the two targets.
- ◊ comparing [0.6,0.8] of the two targets; the butanol and carbon.
- (The red plot) = (The pink plot) X (The scaling factor)

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The particle identification The dilution factor The beam and target polarization

The dilution factor



The particle identification The dilution factor The beam and target polarization

The dilution factor



The average dilution factor is 0.578

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The particle identification The dilution factor The beam and target polarization

the beam and target polarization

- $\diamond~$ Target polarization, $\Lambda_z \sim 0.8$
- $\diamond~$ Electron beam polarization, $P_e \sim 0.85$

$$P_{\gamma} = P_{e} \cdot \frac{(\frac{4}{E_{e}})E_{\gamma} - (\frac{4}{E_{e}})^{2}E_{\gamma}^{2}}{4 - (\frac{4}{E_{e}})E_{\gamma} + 3(\frac{4}{E_{e}})^{2}E_{\gamma}^{2}}$$



The photon energy [GeV]	The photon polarization	
[0.3,0.4]	0.209	
[0.4,0.5]	0.277	
[0.5,0.6]	0.348	
[0.6,0.7]	0.419	
[0.7,0.8]	0.490	
[0.8,0.9]	0.559	
[0.9,1.0]	0.624	
[1.0,1.1]	0.683	
[1.1,1.2]	0.734	
[1.2,1.3]	0.777	
[1.3,1.4]	0.810	
[1.4,1.5]	0.833	
[1.5,1.6]	0.846	

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Polarization Observable P_z^o

$$\mathbf{P}_{\mathbf{z}}^{\odot} = \frac{1}{f \cdot \delta_{\odot} \cdot \Lambda_{\mathbf{z}}} \bigg\{ \frac{\bigg(N(\to \Rightarrow) + N(\to \Rightarrow) \bigg) - \bigg(N(\to \Leftarrow) + N(\leftarrow \Leftarrow) \bigg)}{\bigg(N(\to \Rightarrow) + N(\to \Rightarrow) \bigg) + \bigg(N(\to \Leftarrow) + N(\leftarrow \Leftarrow) \bigg)} \bigg\}$$

- f dilution factor
- δ_{\odot} beam polarization
- Az target polarization
- $N(\rightarrow \Rightarrow)$ the number of events

with the circular beam polarization and longitudinal target polarization

- \rightarrow and $\leftarrow:$ circular polarization of the beam in its two possible settings
- \Rightarrow and \Leftarrow : longitudinal target polarization parallel or anti-parallel to the beam

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Photoproduction of $\pi^+\pi^-$ off the Proton: Kinematics

The π^+ π^- in the final state require 5 independent variables!



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The asymmetry plot for P_z^o

The topology $\gamma p \rightarrow \pi^+ \pi^-(p)$ (Energy Bin 1100 MeV - 1200 MeV)



The asymmetry plot for P_z^o

The topology $\gamma p \rightarrow \pi^+ \pi^-(p)$ (Energy Bin 1200 MeV - 1300 MeV)



The asymmetry plot for P_z^o

The topology $\gamma p \rightarrow \pi^+ \pi^-(p)$ (Energy Bin 1300 MeV - 1400 MeV)



Summary

- The first part of FROST with a longitudinally polarized target has been completed
- $\diamond~$ Preliminary results for $\mathbf{P_z}^{\odot}$ in π^+ π^- photoproduction
- The second part of FROST with a transversely polarized target already start from March 2010 to July 2010

To do:

- Energy Loss Correction
- Momentum Correction
- Kinematic Fitting
- Normalization

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Thank you

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The excited states of the nucleon

Constituent quark models: N^* resonances (Isospin $\frac{1}{2}$)



Side view of the CLAS spectrometer



Creating a Circular Polarized Photon Beam



- Laser → Unpolarized Photons
- Pockels Cell \rightarrow Circularly Polarized Photons
- Ilectron Gun → Low Energy Linearly Polarized Electrons
- Møller Detector \rightarrow Measures Degree of Polarization of Beamline at Radiator

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Creating a Circular Polarized Photon Beam

Bremsstrahlung Radiation is described by QED exactly



Crede, Volker. $\pi^0\eta$ Helicity Difference CB-ELSA Proposal to PAC, 2005

Kammer, Susanne. Strahlpolarimetrie am CBELSA/TAPS Experiment, DPG Meeting 2008

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Particle identification



• Particle identification in CLAS relies on the combination of measured charged-particle momenta (from DC) and the flight time from the target to the respective TOF counters.

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The contamination



How do we make the low temperature?

Refrigeration below 4.2 K - Evaporative cooling

Liquid 4He boils at 4.2K under atmospheric conditions (3He at 3.1K).

Liquid Temperature can be lowered by reducing the vapor Pressure



How do we make the low temperature?

3He/4He Dilution Refrigeration: The Basics

- below 0.8 K, a 3He/4He mixture will separate into two phases



-The specific heat of a 3He atom is higher in the lower, dilute phase than in the upper, concentrated phase.

 $C_d>C_c$

-Therefore, 3He will absorb energy when it dissolves into the dilution phase.

How do we make the low temperature?

Dilution Refrigeration



The cooling precess of 3He

- (1) 1.5 K : in Condenser
- (2) 0.7 K : in Distillation chamber
- (3) 0.005 K : in mixing chamber

- Things near mixing chamber are cooled to around 0.005 K

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How do we make the low temperature?

Horizontal Dilution Refrigerator: T.O. Niinikoski, CERN 1971



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The Scaling factor



The average scaling factor is 5.636

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The dilution factor



The average dilution factor is 0.578

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