

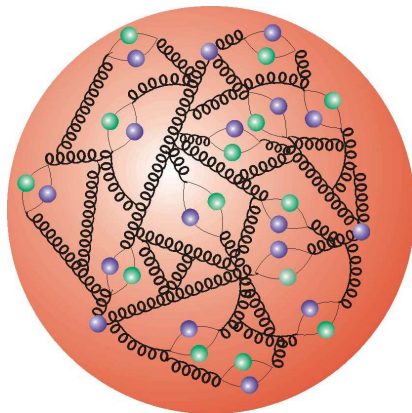
# Photoproduction of $\omega$ Mesons and $\pi^0\omega$ Meson Pairs off the Free Proton

Andrew Wilson

Florida State University  
Physics Department

December 5, 2012  
Dissertation Defense

# Nucleon (Protons / Neutrons)



- composite particle
- contains quarks and gluons
- bound by the strong nuclear force, described by Quantum Chromodynamics (QCD)
- contains 3 "valence" quarks
- contains a "sea" of quark anti-quark pairs and gluons.

What is the nature of the force that binds the proton?

# Strong Interaction

## quarks

u

up

s

strange

b

bottom

d

down

c

charm

t

top

## gauge boson

g

gluon

## Hadrons

**Baryons**



**Mesons**



- Quarks are only found bound inside hadrons.
- Gluons are the force carriers.
- QCD Lagrangian describes the interaction of quarks and gluons

The QCD Lagrangian has been shown to be correct at large interaction energies. ( $> 10$  GeV)

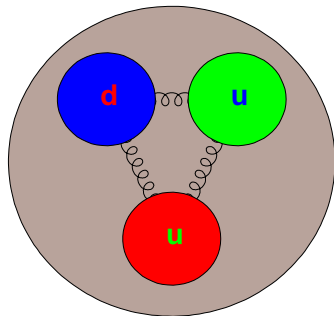
# Solving the QCD Lagrangian

- Perturbative QCD: Used to show the QCD Lagrangian is correct. Can not be used! → Non-perturbative at these energies.
- Numerical Solutions: Lattice QCD (not finished yet)

→ No current way to solve the QCD Lagrangian at these energies.

Can we construct effective models to understand?

# Constituent Quark Model



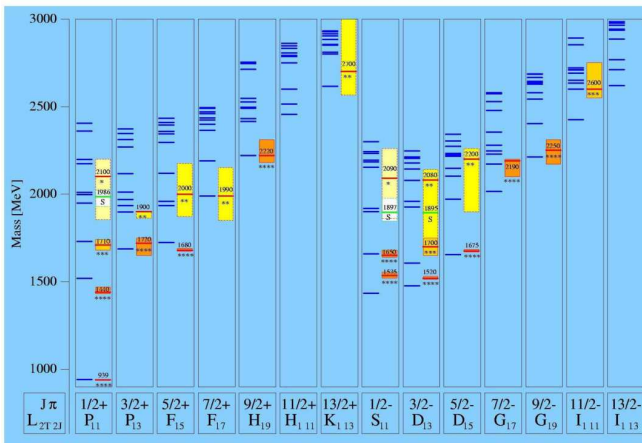
## Constituent Quark Model (CQM)

- Assume only 3 quarks with a fitted mass.
- Propose an interaction which respects the known properties of the strong force.
- Fit the mass of the quarks to recreate the nucleon (proton/neutron) mass.
- Generate the excited states.

# Constituent Quark Model Predictions

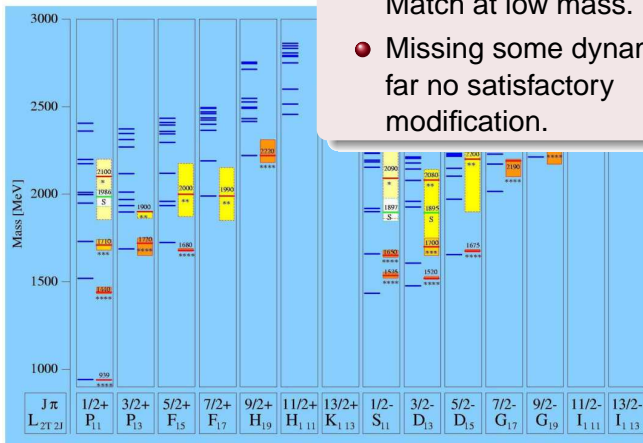
Isospin 1/2, Strangeness 0 Baryon Predictions  
Instanton-based, relativistic solution

(U. Loring *et al.* Eur. Phys. J. A **10**, 395 (2001))



# Constituent Quark Model Predictions

Isospin 1/2, Strangeness 0 Baryon Predictions  
 Instanton-based,  
 (U. Loring *et. al.* Eur. P



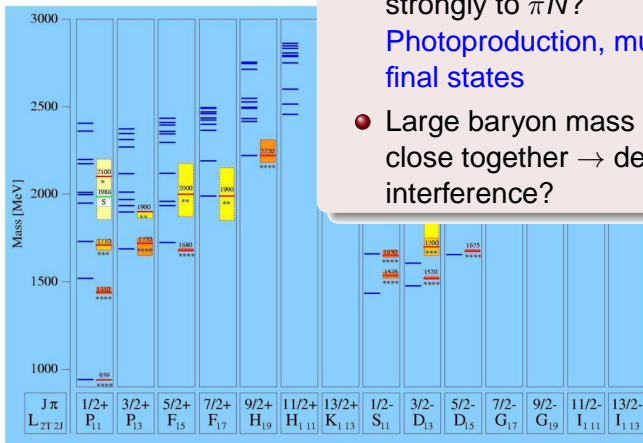
Are the CQMs wrong?

- Must be somewhat correct: Match at low mass.
- Missing some dynamics? So far no satisfactory modification.

# Constituent Quark Model Predictions

Experiment hasn't found them yet?

Isospin 1/2, Strangeness 0  
 Instanton-based,  
 (U. Loring *et. al.* Eur. P



- Most baryon states found using  $\pi N$  scattering. Missing resonances do not couple strongly to  $\pi N$ ?

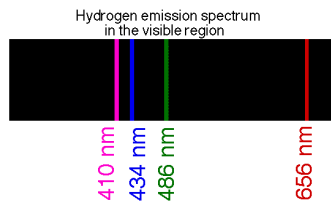
Photoproduction, multi-particle final states

- Large baryon mass widths close together  $\rightarrow$  destructive interference?



# Spectroscopy

The pattern of excited states of a bound state depends directly on the binding force.



Using a historical example:  
Atomic Spectroscopy

- Map out the excited states of the atom
- Create Models to fit the spectrum
- Lead to Quantum Physics and Quantum Electrodynamics

# Baryon Spectroscopy Observables

Scattering Experiment Observables used in this analysis

## Differential Cross Sections ( $\frac{d\sigma}{dX_i}$ )

Conceptually:

A measurement proportional to the probability for a reaction to happen scattering into some final state kinematics ( $X_i$ ).

**Total Cross Section:** Integrate over final state kinematics.

## Polarization Observables

Conceptually:

A relative measurement used to quantify how the differential cross section depends on the spin polarizations of the initial or final state particles.

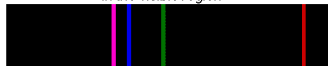
→ Spin-density Matrix Elements ( $\rho$ )

(Cross section dependence on the spin polarization of a single final state particle)

# Spectroscopy

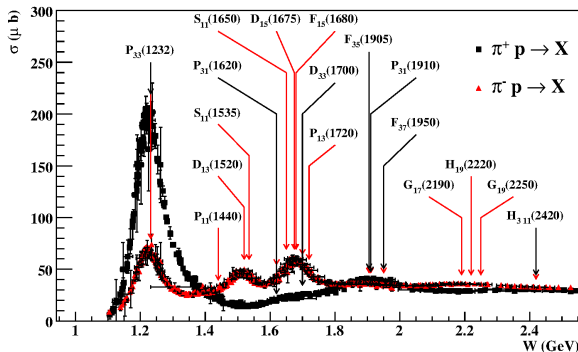
The pattern of excited states of a bound state depends directly on the binding force.

Hydrogen emission spectrum  
in the visible region



Using a historical example:  
Atomic Spectroscopy

- Map out the excited states of

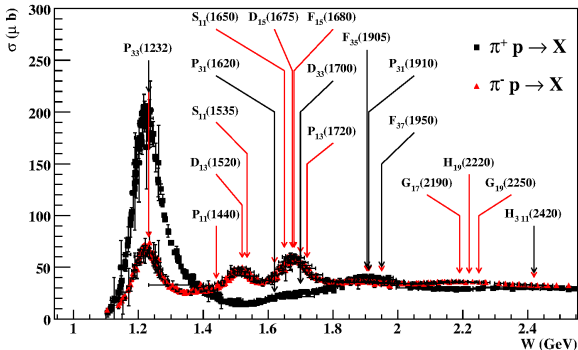


Is to fit the

Atom Physics and  
hydrodynamics

# Sp Baryon Spectroscopy

- Baryon Resonances have short lifetimes  $\rightarrow$  large mass widths  $\rightarrow$  destructive interference
  - Sophisticated interpretation analysis technique (Partial Wave Analysis)
  - Polarization Observables
- Baryon Resonances decay to a zoo of particles: photons, mesons, baryons, leptons.  $\rightarrow$  specialized detectors



# Why $\omega$ photoproduction?

- Quark model predictions suggest many of these "missing" baryon resonances couple significantly to  $p\omega$  and intermediate decays in  $p\pi^0\omega$ .
- The first measurement of observables over the full kinematic range with characteristics useable for isolating baryon resonances.

## Observables Measured

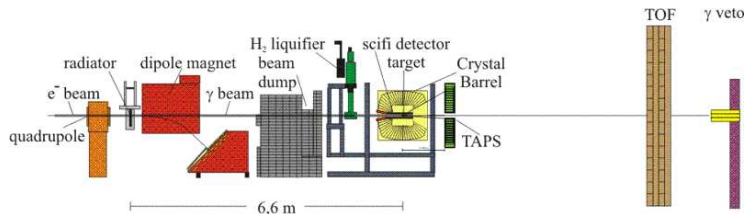
$$\gamma p \rightarrow p\omega$$

Differential cross sections & Spin-density matrix elements

$$\gamma p \rightarrow p\pi^0\omega$$

Differential cross sections

# CBELSA/TAPS Experiment



## Beam

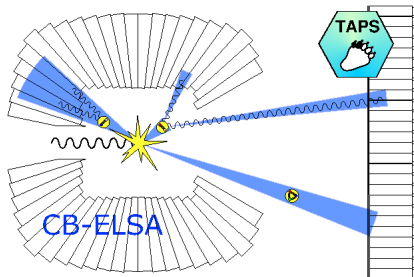
Unpolarized or linearly polarized photons tagged with energy and timing.

## Target

Stationary, unpolarized free protons. Liquid hydrogen.

# CBELSA/TAPS Detectors

- Crystal Barrel Detector: 1290 CsI (TI) crystals + photodiodes
- TAPS Detector: 528 BaF<sub>2</sub> crystals + photomultiplier tubes
- Excellent photon energy and position reconstruction characteristics.
- Only detect the presence of charged particles by scintillating materials between the target and crystals.



# Data

Data recorded : October - November 2002

Beam: Unpolarized tagged photons

Target: Unpolarized Liquid Hydrogen

## Decay channel selected for:

$$\gamma p \rightarrow p\omega \rightarrow p\pi^0\gamma \rightarrow p\gamma\gamma\gamma$$

$$\gamma p \rightarrow p\pi^0\omega \rightarrow p\pi^0\pi^0\gamma \rightarrow p\gamma\gamma\gamma\gamma\gamma$$

(Branching Ratios :  $\omega \rightarrow \pi^0\gamma = 8.9\%$  ,  $\pi^0 \rightarrow \gamma\gamma = 98\%$ )

## Selected events with:

$\gamma p \rightarrow p\omega$  3 uncharged reconstructed particles & 0-1 charged reconstructed particles

$\gamma p \rightarrow p\pi^0\omega$  5 uncharged reconstructed particles & 0-1 charged reconstructed particles

## Monte Carlo Simulated Data

Detector Acceptance, Cut tuning.



# Data Reduction Cuts

- Multiplicity Cuts - 3(5) photons and less than 2 protons
- Timing Cut - Imposes causality between the initial photon and final state particles. Reduces the number of initial photon candidates.
- Coplanarity Cut - cut away events where extra undetected particles could have been involved
- Trigger Cut - cut away events where proton identification simulation could be wrong.

# Kinematic Fitting

To evaluate the probability of an event being a desired final state.

## Definition

- Varies the measured values within the quoted errors of the experiment
- Matches to an "ideally" measured event using Energy-Momentum conservation and invariant masses
- Returns a confidence level value between 0 - 1 for each event.
- Used to select initial photon.

Natural Mass Width of the  $\omega$  meson  $\sim$  measurement error

Natural Mass Width of the  $\pi^0$  meson  $\sim 8$  eV  $\rightarrow$  Negligible

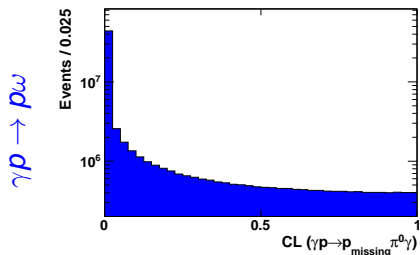
## Fit Hypothesis

$$\gamma p \rightarrow p \omega : \gamma p \rightarrow p_{\text{missing}} \pi^0 \gamma$$

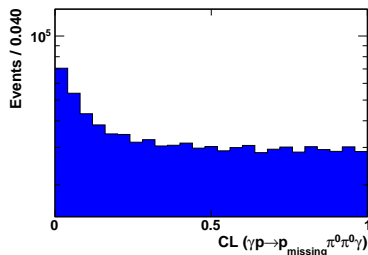
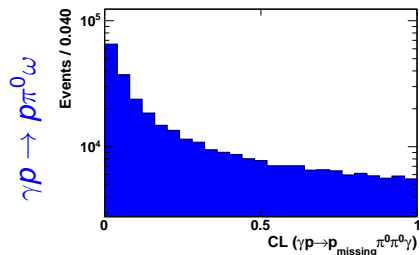
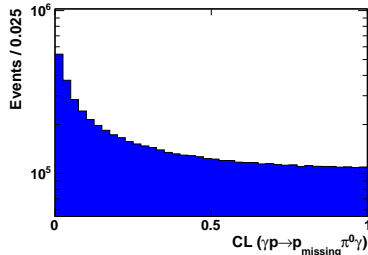
$$\gamma p \rightarrow p \pi^0 \omega : \gamma p \rightarrow p_{\text{missing}} \pi^0 \pi^0 \gamma$$

# Kinematic Fitting Results

## Experimental Data

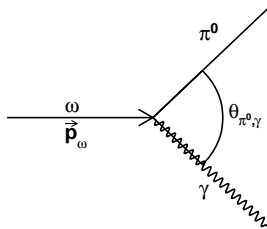


## Monte Carlo Simulated

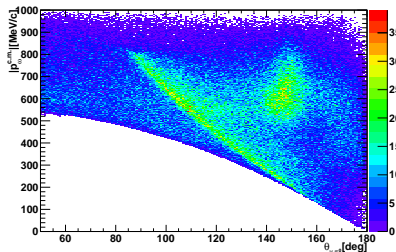


Kept all events with  $CL > 0.005$

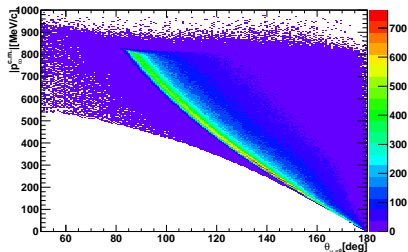
# $\gamma p \rightarrow p\omega$ Opening Angle Cut



## Experimental Data

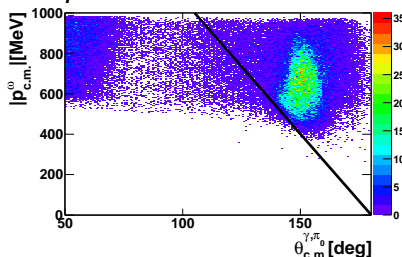


## Monte Carlo Simulated

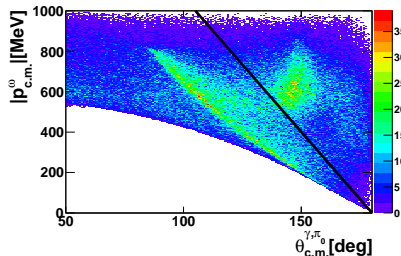


# $\gamma p \rightarrow p\omega$ Opening Angle Cut

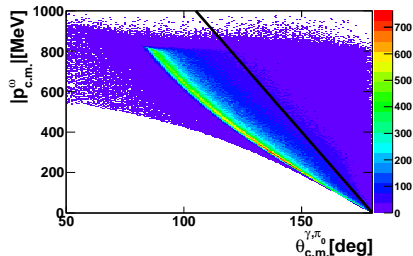
$p\pi^0$  Monte Carlo Simulated



Experimental Data



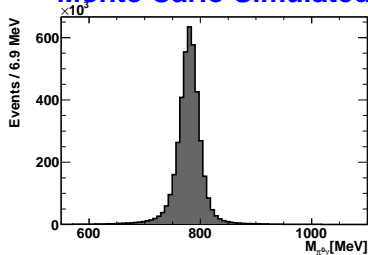
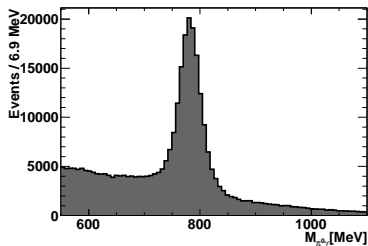
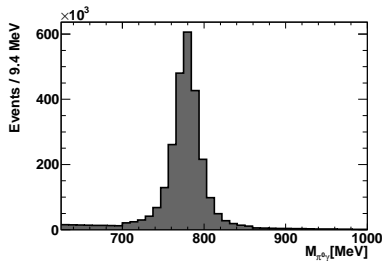
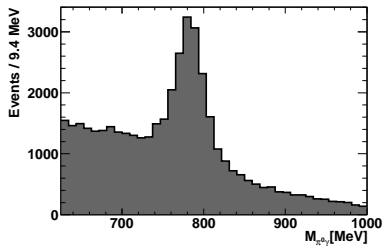
Monte Carlo Simulated



# FINALLY!... Invariant Mass Distributions

## Experimental Data

## Monte Carlo Simulated

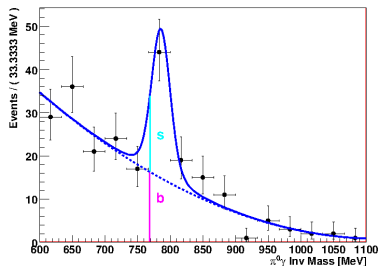
 $\gamma p \rightarrow p\omega$ 

 $\gamma p \rightarrow p\pi^0\omega$ 


Two entries per event in  $\gamma p \rightarrow p\pi^0\omega$  plots.

# Q-factor Background Subtraction

For each event left in the analysis.

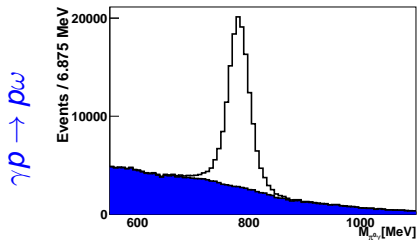
- Find the nearest neighbors in the final state's kinematic phase space.
- Fit the invariant mass spectrum of a particle in the desired final state. (background & signal functions)
- Define a Q-factor from the fit. (Probability the event is the desired final state)
- Weight each event with the Q-factor.



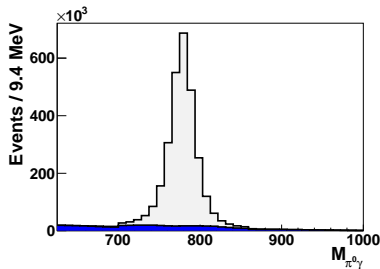
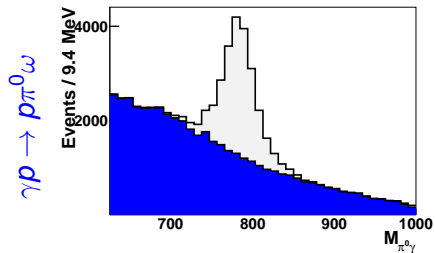
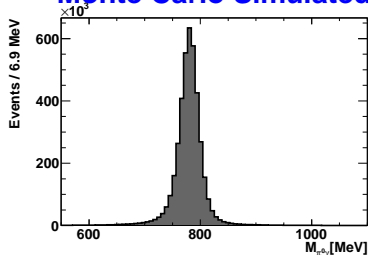
$$Q = \frac{s}{s+b}$$

# Background Subtracted

## Experimental Data



## Monte Carlo Simulated



Two entries per event in  $\gamma p \rightarrow p\pi^0\omega$  plots.



# Differential Cross Section Methodology

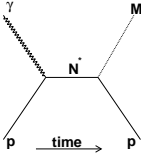
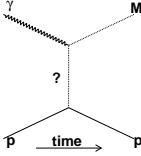
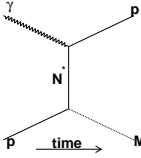
$$\frac{d\sigma}{dX_i} = \frac{N_{data}}{\mathcal{A} \mathcal{F} \rho_t^A B_r \Delta X_i}$$

$N_{data}$	# of observed events (timing background subtracted)
$\mathcal{A}$	detector acceptance
$\mathcal{F}$	Photon Flux
$\rho_t^A$	Target Area Density
$B_r$	Branching Ratio
$X_i$	Kinematic Variable
$\Delta X_i$	Kinematic Bin Width

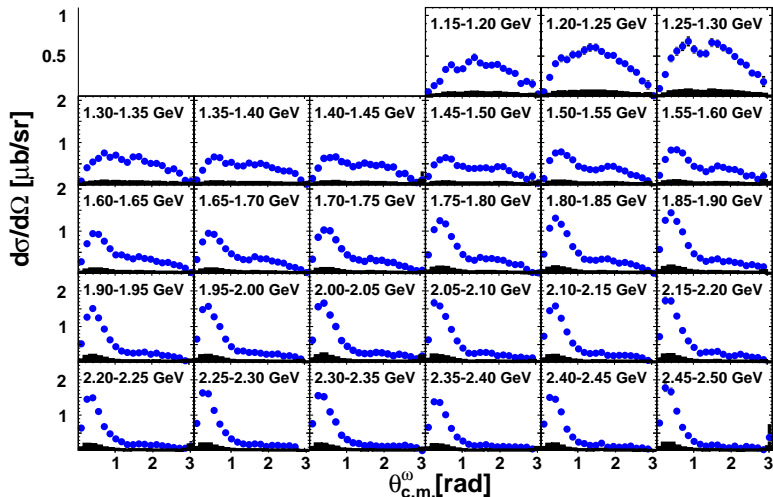
# Differential Cross Section Uncertainties

- Statistical Uncertainties - Measurement Precision
- Systematic Uncertainties - Measurement Accuracy
  - Possible Target Shift
  - Acceptance Correction ( $\gamma p \rightarrow p\pi^0\omega$  Only)
  - Q-factor Fit Uncertainties
  - Confidence Level Systematic Uncertainty
  - Monte Carlo Simulation Uncertainty
  - Background Contributions from other reactions

# Differential Cross Section Interpretation

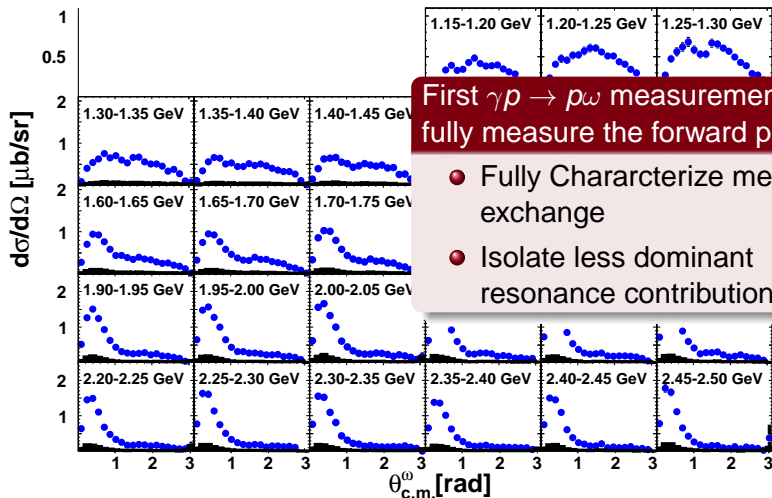
Name	Process	Meson Angular Distribution
Resonance Production		symmetric
Meson Exchange		forward angle peaked
Baryon Exchange		backward angle peaked

# $\gamma p \rightarrow p\omega$ Differential Cross Sections



Labeled with incoming photon energy.

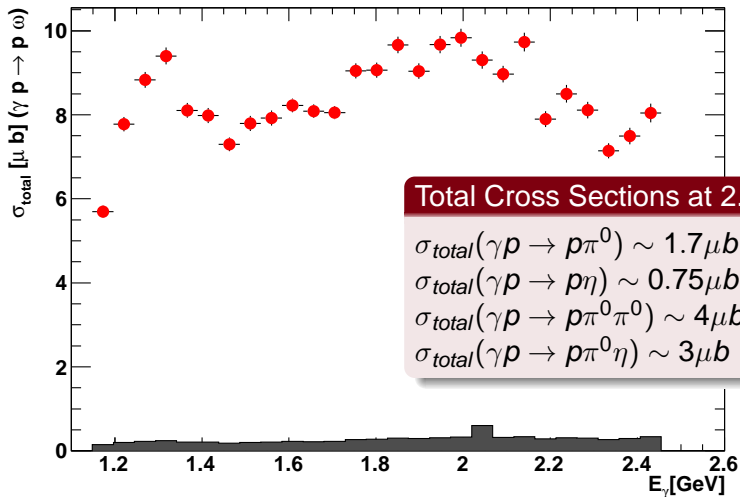
# $\gamma p \rightarrow p\omega$ Differential Cross Sections



First  $\gamma p \rightarrow p\omega$  measurement to fully measure the forward peak.

- Fully Characterize meson exchange
- Isolate less dominant resonance contributions

# $\gamma p \rightarrow p\omega$ Total Cross Section



# Unpolarized $\gamma p \rightarrow p \omega$ Spin Density Matrix Elements

Information in the spin of the  $\omega$  meson can be extracted by analyzing the decay angular distribution.

$$\rho_{ij} \sim M_i M_j^*$$

$i$  and  $j$  is the spin polarization of the  $\omega$   $(-1, 0, 1)$

## Unpolarized Spin-density Matrix Element Fitting

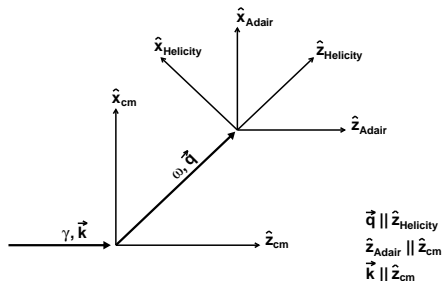
$$W^0(\theta_d, \phi_d, \rho^0) = \frac{3}{8\pi} (\sin^2 \theta_d \rho_{00}^0 + (1 + \cos^2 \theta_d) \rho_{11}^0) \\ + \sin^2 \theta_d \cos 2\phi_d \rho_{1-1}^0 + \sqrt{2} \sin 2\theta_d \cos \phi_d \operatorname{Re} \rho_{10}^0$$

$\theta_d$  and  $\phi_d$  angles of the  $\gamma$  in  $\omega \rightarrow \pi^0 \gamma$  in  $\omega$  rest frame.

Binning used : 6  $\cos \theta_d$  bins and 8  $\phi_d$  bins

# Spin Density Matrix Element Reference Systems

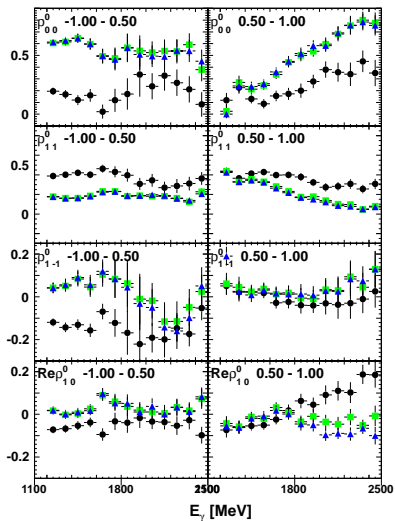
- Helicity system - z-axis  $\parallel \vec{q}_\omega^{c.m.}$
- Adair system - z-axis  $\parallel \vec{k}_{c.m.}$
- Gottfreid-Jackson system - z-axis  $\parallel \vec{k}_\omega$  frame .



y-axis is defined to be  $\hat{k} \times \hat{q}$   
 Extracted SDMEs in all three systems.



# Unpolarized $\gamma p \rightarrow p\omega$ Spin Density Matrix Elements

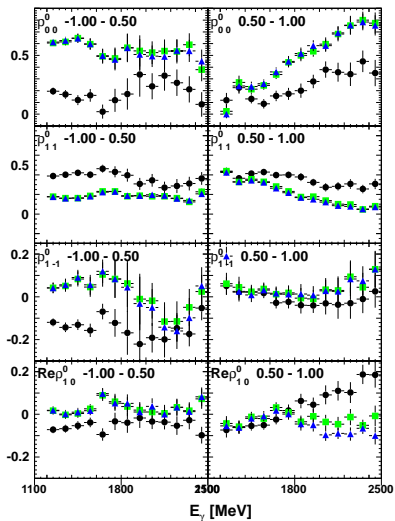


Labeled with SDME and  $\cos \theta_{c.m.}^\omega$ .

- Same Data used in differential cross sections
- Q-factors used to fill angular distributions
- Angular distributions were acceptance corrected
- Monte Carlo simulation uncertainty, Q-factor fit errors, and statistical errors propagated
- Extracted using a  $\chi^2$  minimization fit

● Helicity System    ■ Gottfried-Jackson system    ▲ Adair system

# Unpolarized $\gamma p \rightarrow p\omega$ Spin Density Matrix Elements



Labeled with SDME and  $\cos \theta_{C.m.}^\omega$

First  $\gamma p \rightarrow p\omega$  Unpolarized SDMEs measured over the full kinematic range

- Helicity System
- Gottfried-Jackson system
- ▲ Adair system

# Polarized $\gamma p \rightarrow p \omega$ Spin Density Matrix Elements

transversely polarized beam photons

## Polarized Spin-density Matrix Element Fitting

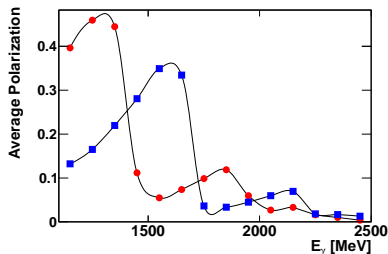
$$\begin{aligned}
 W^L(\cos \theta_d, \phi_d, \Phi_{pol}, \rho) = & \frac{3}{8\pi} (\sin^2 \theta_d \rho_{00}^0 + (1 + \cos^2 \theta_d) \rho_{11}^0) \\
 & - \frac{3}{8\pi} P_\gamma \cos 2\Phi_{pol} (\sin^2 \theta_d \rho_{00}^1 \\
 & + (1 + \cos^2 \theta_d) \rho_{11}^1)
 \end{aligned}$$

$\theta_d$  and  $\phi_d$  angles of the  $\gamma$  in  $\omega \rightarrow \pi^0 \gamma$  in  $\omega$  rest frame.

$\Phi_{pol}$  is the angle between the production plane ( $\omega + p_{final}$ ) and the polarization direction.

# Polarized $\gamma p \rightarrow p\omega$ Data

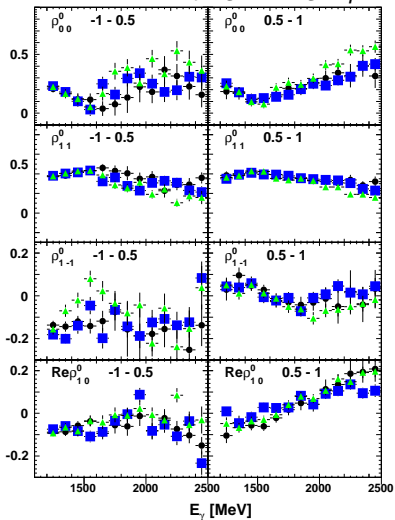
- March 2003 (1350 MeV Coherent Edge) (●)
- May 2003 (1600 MeV Coherent Edge) (■)



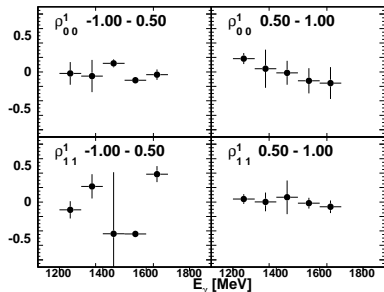
Nearly the same reconstruction and reaction selection.  
No TAPS trigger cut due to a slightly different trigger.

# Polarized $\gamma p \rightarrow p \omega$ Spin Density Matrix Elements

Extracted by ignoring  $\Phi_{pol}$



Integrated over  $\phi_d$  before fitting.



First Measurement

Helicity system

Labeled with SDME and  $\cos \theta_{c.m.}^\omega$

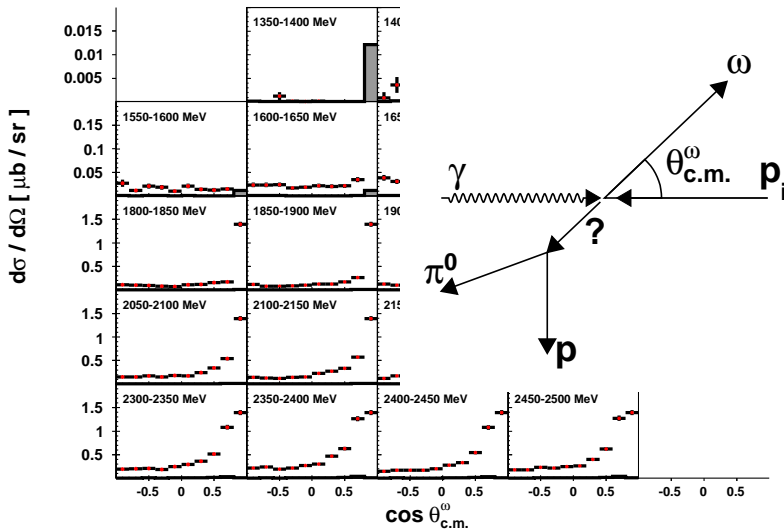
- unpolarized data
- March 2003 data
- ▲ May 2003 data

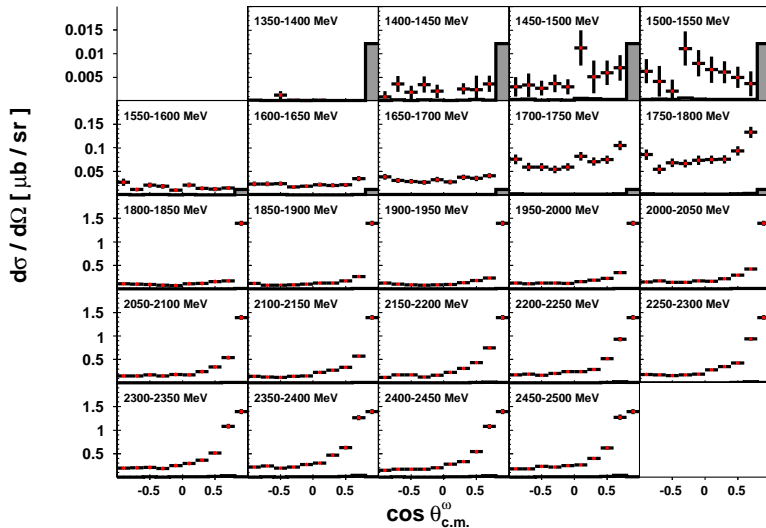
# $\gamma p \rightarrow p \pi^0 \omega$ Differential Cross Sections Overveiw

- 5 Independent Kinematic Variables
- 3 different Kinematic Variable binning
- 12,500 Differential Cross Section data points per binning

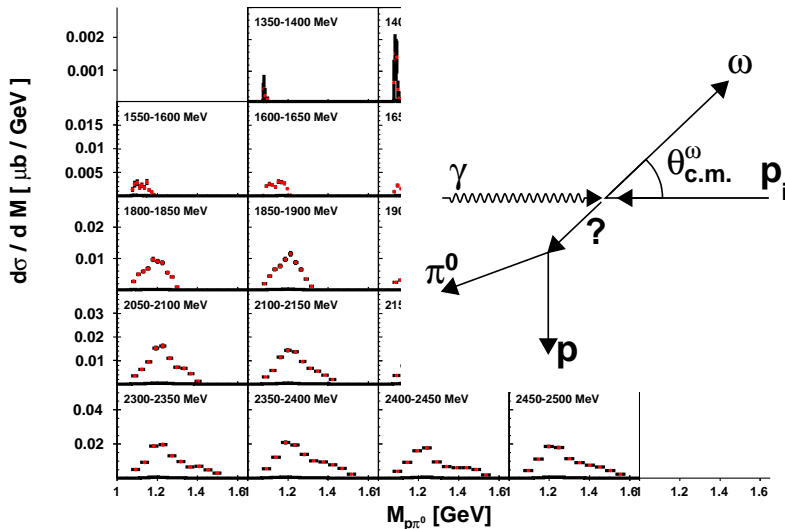
Binning Name	$a$	$b_1$	$b_2$
$\omega_{cms}$	$\omega$	$\pi^0$	$p$
$\pi^0_{cms}$	$\pi^0$	$p$	$\omega$
$p_{cms}$	$p$	$\omega$	$\pi^0$

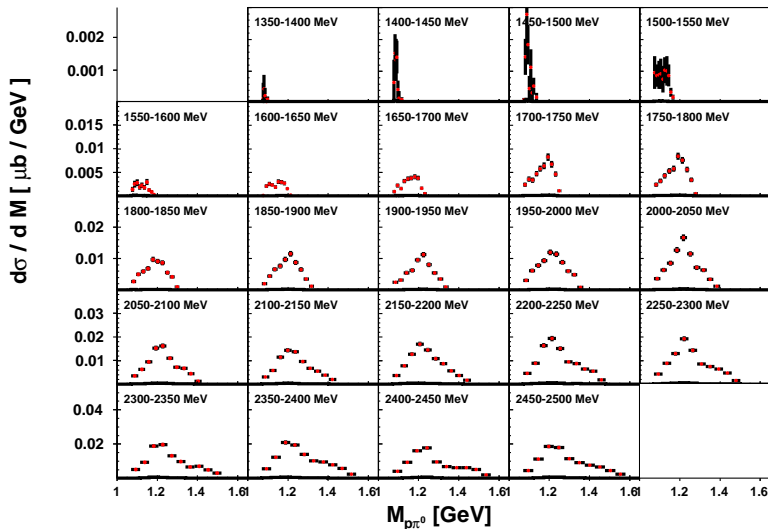
# $\gamma p \rightarrow p\pi^0\omega$ $\omega_{cms}$ Differential Cross Section



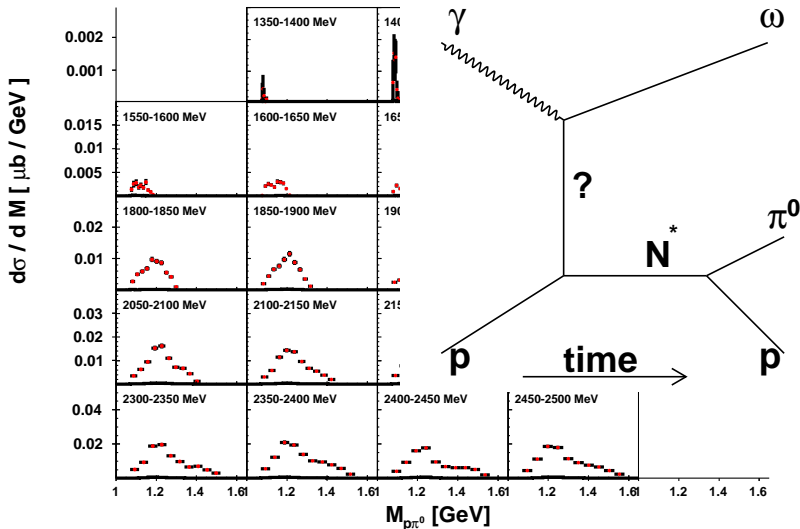
$\gamma p \rightarrow p\pi^0\omega$   $\omega_{CMS}$  Differential Cross Section




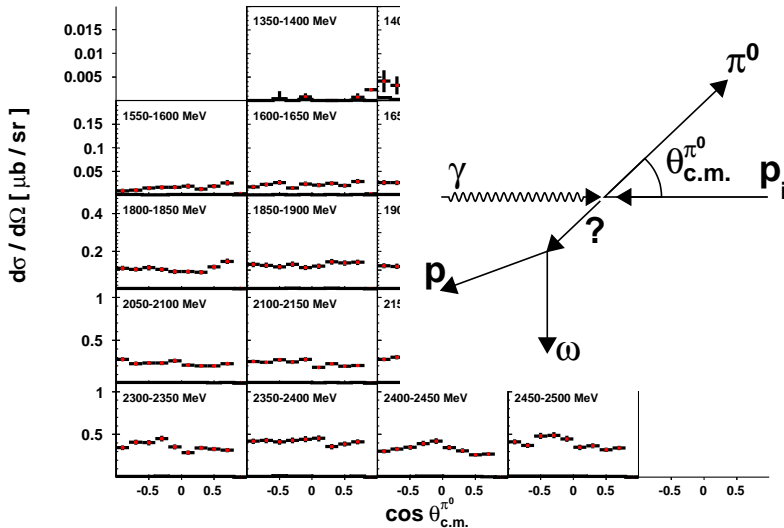
$\gamma p \rightarrow p \pi^0 \omega$   $\omega_{cms}$  Differential Cross Section


$\gamma p \rightarrow p \pi^0 \omega$   $\omega_{CMS}$  Differential Cross Section


# $\gamma p \rightarrow p \pi^0 \omega$ $\omega_{cms}$ Differential Cross Section

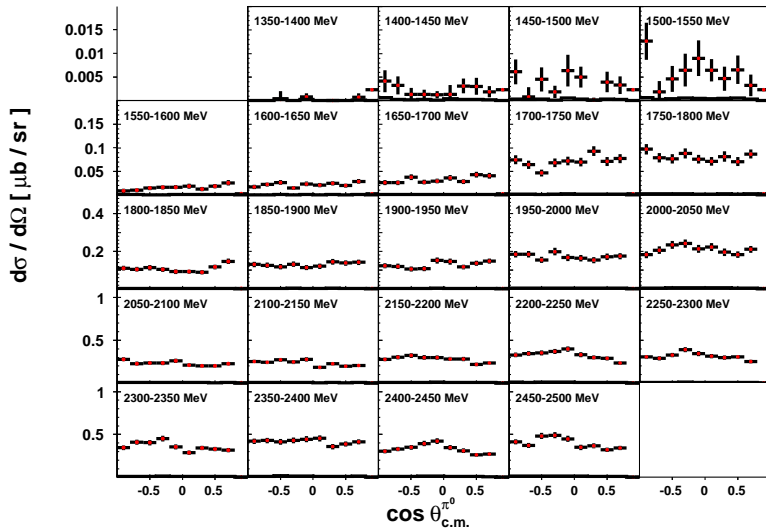


# $\gamma p \rightarrow p \pi^0 \omega \pi^0_{cms}$ Differential Cross Section



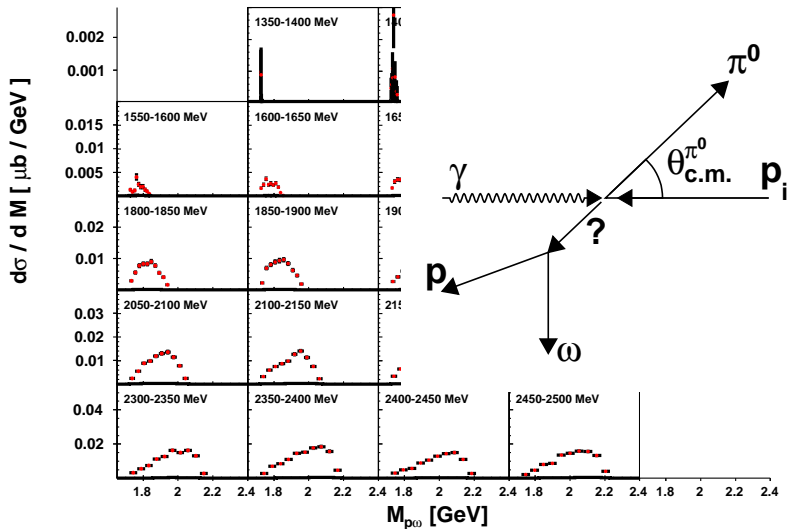
$$\gamma p \rightarrow p \pi^0 \omega \pi^0_{cms}$$

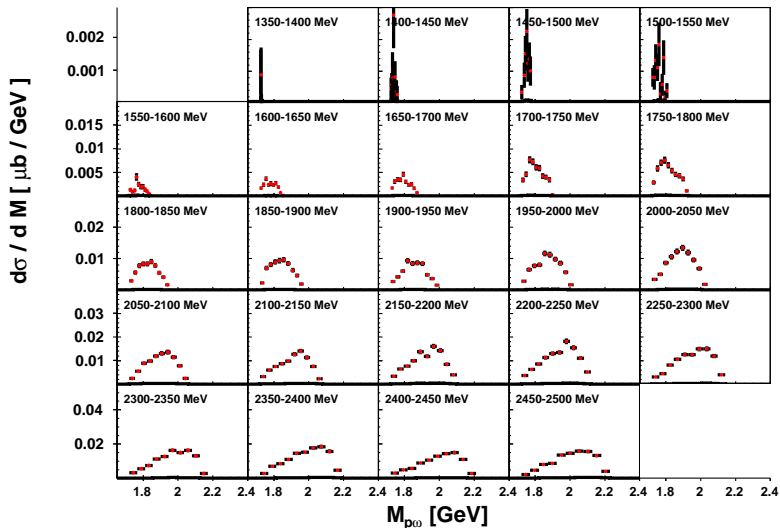
# Differential Cross Section



$$\gamma p \rightarrow p \pi^0 \omega \pi^0_{cms}$$

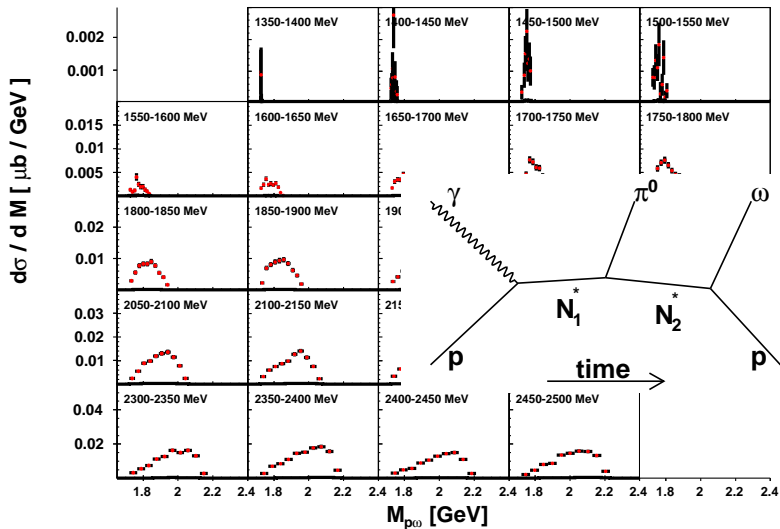
# Differential Cross Section



$\gamma p \rightarrow p \pi^0 \omega \pi^0_{cms}$  Differential Cross Section


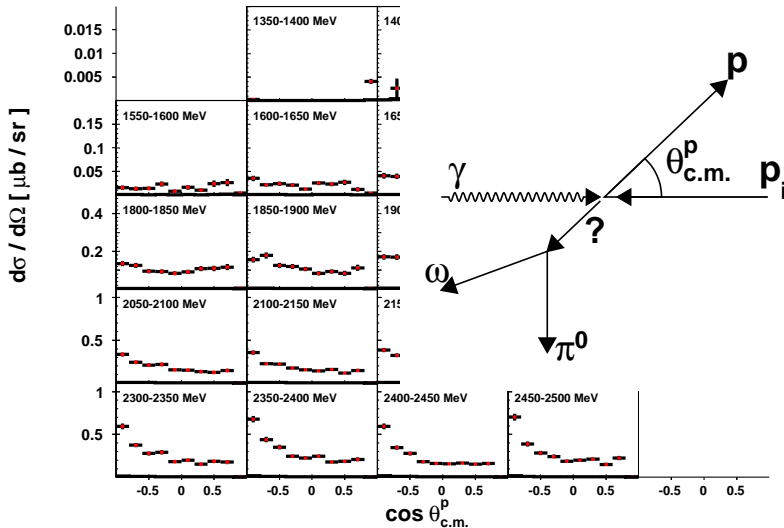
$$\gamma p \rightarrow p \pi^0 \omega \pi^0_{cms}$$

# Differential Cross Section

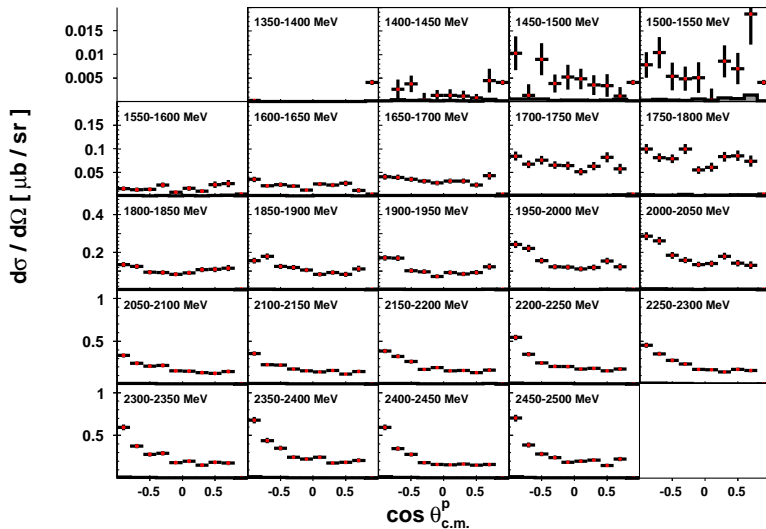




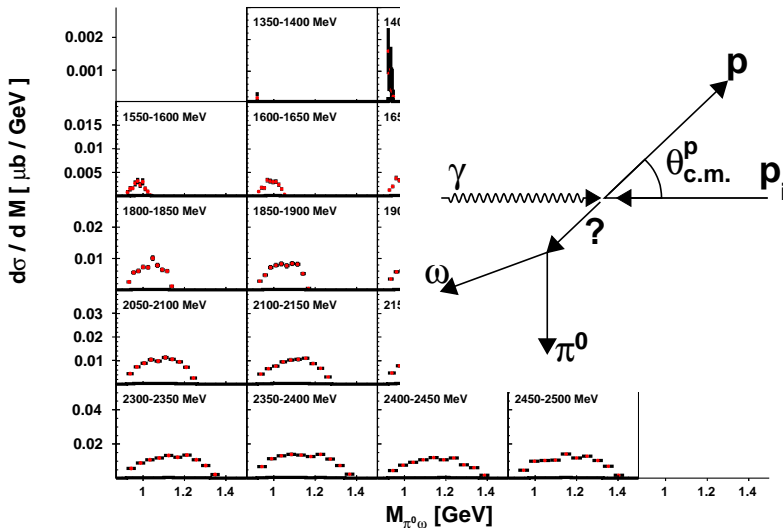
# $\gamma p \rightarrow p \pi^0 \omega$ $p_{cms}$ Differential Cross Section



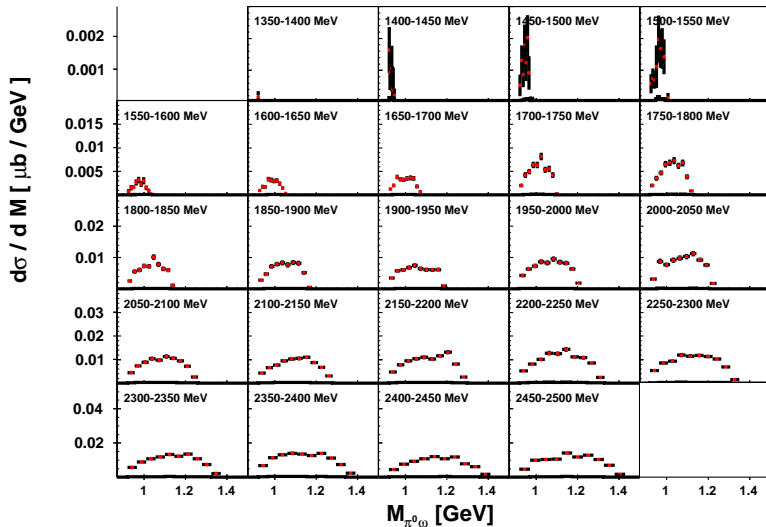
# $\gamma p \rightarrow p \pi^0 \omega$ $p_{cms}$ Differential Cross Section



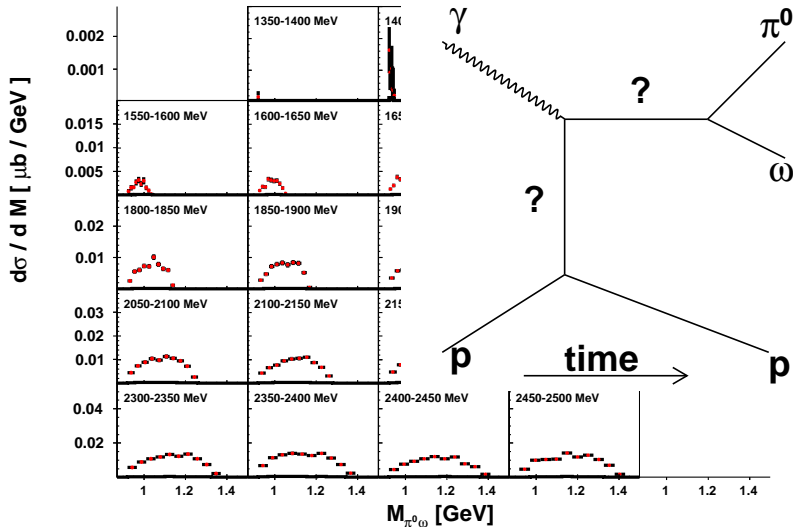
# $\gamma p \rightarrow p \pi^0 \omega$ $p_{cms}$ Differential Cross Section



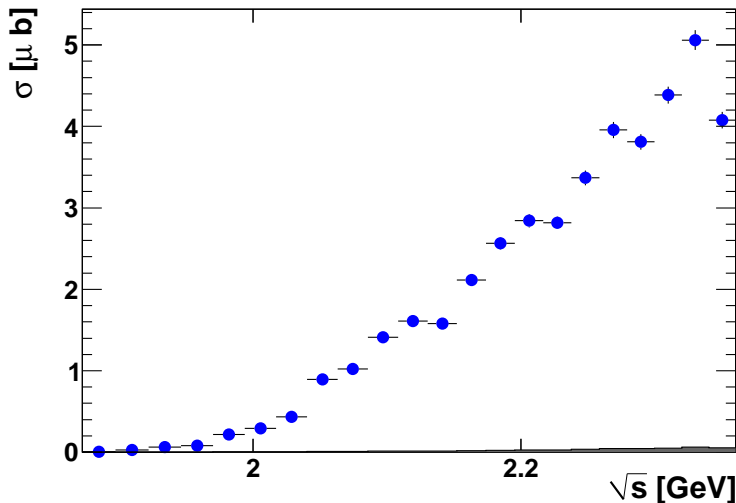
# $\gamma p \rightarrow p \pi^0 \omega$ $p_{cms}$ Differential Cross Section



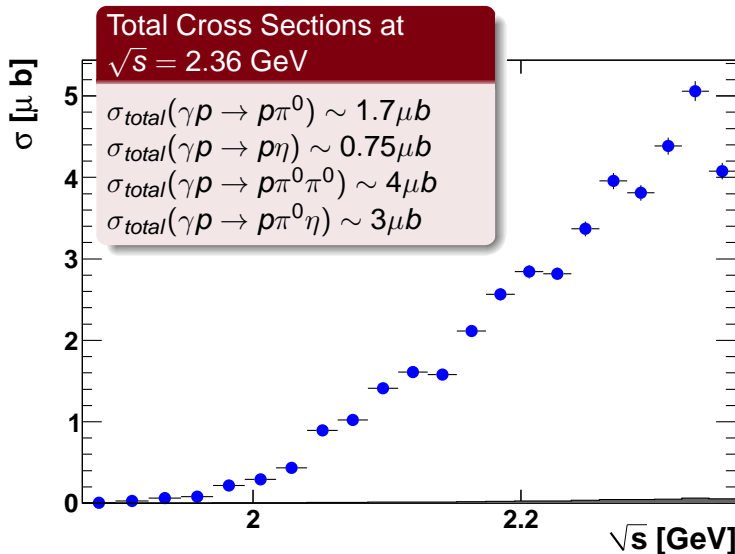
# $\gamma p \rightarrow p \pi^0 \omega$ $p_{cms}$ Differential Cross Section



# $\gamma p \rightarrow p\pi^0\omega$ Total Cross Section



# $\gamma p \rightarrow p\pi^0\omega$ Total Cross Section



# Summary and Future Work

## First Time Measurements

- $\gamma p \rightarrow p\omega$  and  $\gamma p \rightarrow p\pi^0\omega$  differential cross sections,  $\gamma p \rightarrow p\omega$  Spin-density Matrix Elements
  - have been measured over the full reaction kinematics
  - characteristics usable for isolating "missing" baryon resonances.
- $\gamma p \rightarrow p\omega$  polarized spin-density matrix elements
- $b_1(1235)$  meson photoproduction measurements at these energies

## Future Work

- $\gamma p \rightarrow p\omega$  and  $\gamma p \rightarrow p\pi^0\omega$  interpretation analysis
- $\gamma p \rightarrow p\pi^0\omega$  spin-density matrix elements



# Acknowledgements

## Thank You!

- Committee Members
- Volker Crede
- My Wife
- My Parents